

Raising & lowering operator

Brief reminder of a new notation for spin states $|S, m\rangle$ for $|j, m\rangle$
 In this notation we are "locked" to

$$z\text{-basis} \quad \hat{S}^2 = \hat{S}_x^2 + \hat{S}_y^2 + \hat{S}_z^2 \quad \text{or} \quad \hat{J}^2 = \hat{J}_x^2 + \hat{J}_y^2 + \hat{J}_z^2$$

$$\hat{J}^2 |j, m\rangle = \hbar^2 j(j+1) |j, m\rangle$$

$$\hat{J}_z |j, m\rangle = \hbar m |j, m\rangle$$

Spin- $1/2$ particle $j = S = 1/2, m = \pm 1/2$

Spin-1 particle $j = S = 1, m = 0, \pm 1$

Eigenstates of \hat{J}^2 and \hat{J}_z are:

$$|1, 1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad |1, 0\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad |1, -1\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\hat{S}_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad \hat{S}_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\hat{S}_y = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}$$

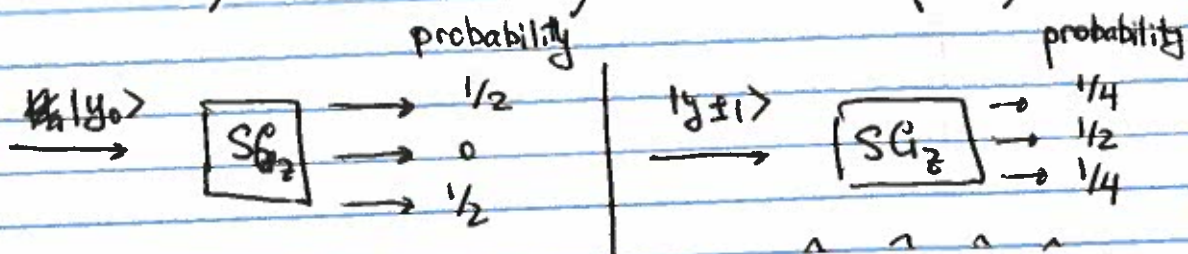
If $|y_{0, \pm 1}\rangle$ are the eigenstates of \hat{S}_y

$$\hat{S}_y |y_0\rangle = 0$$

$$\hat{S}_y |y_{\pm 1}\rangle = \pm \hbar |y_{\pm 1}\rangle \leftarrow \text{true in any basis}$$

$$\frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} |y_{\pm 1}\rangle = \pm \hbar |y_{\pm 1}\rangle \rightarrow \text{find } |y_{\pm 1}\rangle \text{ in } z\text{-basis}$$

$$|y_0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ +1 \end{pmatrix} \quad |y_{+1}\rangle = \frac{1}{2} \begin{pmatrix} 1 \\ i\sqrt{2} \\ -1 \end{pmatrix} \quad |y_{-1}\rangle = \frac{1}{2} \begin{pmatrix} 1 \\ -i\sqrt{2} \\ -1 \end{pmatrix}$$



Extra family members in $\hat{J}_x, \hat{J}_y, \hat{J}_z, \hat{J}^2$ family

$$\hat{J}_+ = \hat{J}_x + i\hat{J}_y$$

raising operator

$$\hat{J}_- = \hat{J}_x - i\hat{J}_y$$

lowering operator

In general $\hat{J}_+ |j, m\rangle \propto |j, m+1\rangle$

$$\hat{J}_- |j, m\rangle \propto |j, m-1\rangle$$

$$j \quad \begin{matrix} \text{--- } m=j \text{ ---} \\ \text{--- } m=j-1 \text{ ---} \end{matrix}$$

$$\begin{matrix} \text{--- } m=-j+1 \text{ ---} \\ \text{--- } m=-j \text{ ---} \end{matrix} \quad \begin{matrix} \uparrow \hat{J}_+ \\ \downarrow \hat{J}_- \end{matrix}$$

also called
ladder
operators

Spin $1/2$

$$\hat{J}_+ |1/2, 1/2\rangle = 0$$

$$\hat{J}_- |1/2, -1/2\rangle = \text{const} \cdot |1/2, 1/2\rangle$$

$$\hat{J}_+ |1/2, -1/2\rangle = \text{const} |1/2, 1/2\rangle$$

$$\hat{J}_- |1/2, -1/2\rangle = 0$$

$$\hat{J}_{\pm} = \hat{J}_x \pm i\hat{J}_y = \frac{\hbar}{2} \left[\begin{pmatrix} 0 & \pm 1 \\ 1 & 0 \end{pmatrix} \pm i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \right]$$

$$\hat{J}_+ = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$\hat{J}_- = \frac{\hbar}{2} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

Raising & lowering operators

$$\hat{J}_+ = \hat{J}_x + i\hat{J}_y = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} + \frac{i\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ +i & 0 & -i \\ 0 & i & 0 \end{pmatrix} = \sqrt{2}\hbar \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\hat{J}_- = \hat{J}_x - i\hat{J}_y = \sqrt{2}\hbar \begin{pmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

As in the case of spin $-1/2$ particle,

\hat{J}_+ action will convert a state with m into a state $m+1$, and \hat{J}_- from m to $m-1$

$$J_+ |1, -1\rangle = \sqrt{2}\hbar \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \sqrt{2}\hbar \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \sqrt{2}\hbar |1, 0\rangle$$

$$J_+ |1, 0\rangle = \sqrt{2}\hbar |1, 1\rangle$$

$$J_+ |1, 1\rangle = \sqrt{2}\hbar \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = 0 \quad \hat{J}_+ |1, 1\rangle = 0$$

since $m=1$ is max.

Similarly $\hat{J}_- |1, 1\rangle = \sqrt{2}\hbar |1, 0\rangle$, $\hat{J}_- |1, 0\rangle = \sqrt{2}\hbar |1, -1\rangle$
and $\hat{J}_- |1, -1\rangle = 0$

These behavior is universal, and occurs for any ~~states~~ triplets of operators that follow the same commutation relationships

How can we prove that $\hat{J}_+ |j, m\rangle \propto |j, m+1\rangle$?

$$[\hat{J}_z, \hat{J}_\pm] = [\hat{J}_z, \hat{J}_x \pm i\hat{J}_y] = [\hat{J}_z, \hat{J}_x] \pm i[\hat{J}_z, \hat{J}_y] = \\ = i\hbar \hat{J}_y \pm i(-i\hbar \hat{J}_x) = i\hbar \hat{J}_y \pm \hbar \hat{J}_x = \pm \hbar (\hat{J}_x \pm i\hat{J}_y) = \pm \hbar \hat{J}_\pm$$

Thus, if $\hat{J}_z |j, m\rangle = \hbar m |j, m\rangle$, $\hat{J}_z \hat{J}_\pm - \hat{J}_\pm \hat{J}_z = \pm \hbar \hat{J}_\pm$

$$\text{Then } \hat{J}_z \hat{J}_+ |j, m\rangle = (\hat{J}_+ \hat{J}_z + \hbar \hat{J}_+) |j, m\rangle = \\ = \hat{J}_+ \underbrace{[\hat{J}_z |j, m\rangle]}_{\hbar m |j, m\rangle} + \hbar \hat{J}_+ |j, m\rangle = \hat{J}_+ \hbar (m+1) |j, m\rangle$$

$$\text{So } \hat{J}_z [\hat{J}_+ |j, m\rangle] = \hbar (m+1) [\hat{J}_+ |j, m\rangle]$$

must be identical to

$$\hat{J}_z |j, m+1\rangle = \hbar (m+1) |j, m+1\rangle$$

Thus $\hat{J}_+ |j, m\rangle \propto |j, m+1\rangle$

Now let's assume that

$$\hat{J}_+ |j, m\rangle = c_{jm} \hbar |j, m+1\rangle \text{ and find } c_{jm}$$

$$\langle j, m | \hat{J}_- \hat{J}_+ |j, m\rangle = \langle j, m | \hat{J}_+^\dagger \hat{J}_+ |j, m\rangle =$$

$$= |c_{jm}|^2 \hbar^2 \langle j, m+1 | j, m+1\rangle = |c_{jm}|^2 \hbar^2$$

$$\text{At the same time } \hat{J}_- \hat{J}_+ = (\hat{J}_x - i\hat{J}_y)(\hat{J}_x + i\hat{J}_y) =$$

$$= \frac{1}{2} \hat{J}_x^2 + \hat{J}_y^2 + i(\hat{J}_y \hat{J}_y - \hat{J}_y \hat{J}_x) = \hat{J}^2 - \hat{J}_z^2 + i \cdot i \hbar \hat{J}_z =$$

$$= \hat{J}^2 - \hat{J}_z^2 - \hbar \hat{J}_z$$

$$\langle j, m | \hat{J}^2 - \hat{J}_z^2 - \hbar \hat{J}_z | j, m \rangle = \hbar^2 j(j+1) - \hbar^2 m^2 - \hbar^2 m$$

$$|c_{jm}|^2 \cdot \hbar^2 = \hbar^2 (j(j+1) - m(m+1))$$

since $|c_{jm}|^2 \geq 0$ ~~thus we get~~ $-j-1 \leq m \leq j$

$$c_{jm} = \sqrt{j(j+1) - m(m+1)}$$

$$\hat{J}_+ | j, m \rangle = \hbar \sqrt{j(j+1) - m(m+1)} | j, m+1 \rangle$$

automatically get $\hat{J}_+ | j, j \rangle = 0$ $m_{\max} = j$

We can repeat the steps for

$$\hat{J}_- | j, m \rangle = d_{jm} | j, m-1 \rangle$$

and $\langle j, m | \hat{J}_+ \hat{J}_- | j, m \rangle = \langle j, m | \hat{J}^2 - \hat{J}_z^2 + \hbar \hat{J}_z | j, m \rangle$
 $= \hbar^2 (j(j+1) - m(m-1))$

since $|d_{jm}|^2 \geq 0$ $-j \leq m \leq j+1$

$$\hat{J}_- | j, m \rangle = \hbar \sqrt{j(j+1) - m(m-1)} | j, m-1 \rangle$$

and $-m \leq j \leq m$

Spin $-1/2$ $j = 1/2$ $m = \pm 1/2$

Spin -1 $j = 1$ $m = 0, \pm 1$

Also, since we must be able to "step" from minimum $m = -j$ to max $m = j$ in integer number of steps $= 2j$, then j must be either integer, or half-integer