

Moving to 2D and 3D world

So far we considered 1D motion

1 spatial coordinate $\hat{x} \leftrightarrow \hat{p}_x$

Schrodinger eqn $-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + U(x)\psi(x) = E\psi(x)$

Two-dimensional motion: 2 spatial coordinates

$\hat{x} \rightarrow \hat{p}_x$

$\hat{p}_x = i\hbar \frac{\partial}{\partial x}$

$\hat{y} \rightarrow \hat{p}_y$

$\hat{p}_y = i\hbar \frac{\partial}{\partial y}$

} cartesian coordinates

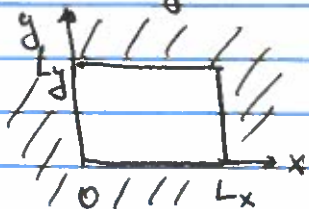
$\hat{H} = \frac{\hat{p}^2}{2m} + U(\hat{r}) = \frac{\hat{p}_x^2 + \hat{p}_y^2}{2m} + U(\hat{x}, \hat{y})$

Schrodinger equation $\hat{H}\psi(x,y) = E\psi(x,y)$

$-\frac{\hbar^2}{2m} \left[\frac{\partial^2\psi(x,y)}{\partial x^2} + \frac{\partial^2\psi(x,y)}{\partial y^2} \right] + U(x,y)\psi(x,y) = E\psi$

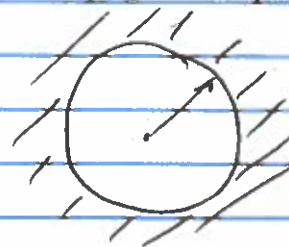
Two common geometries for a 2D potential wells

Rectangular well



Natural to use Cartesian coordinates

Radially/spherically symmetric well



Will use polar (2D) or spherical (3D) coordinates.

Rectangular ^{or square} potential well

$$U(x,y) = \begin{cases} 0 & 0 < x,y < L \\ \infty & \text{otherwise} \end{cases}$$

A particle will bounce freely b/w four walls. Note, that the motion along x and y are independent degrees of freedom.

That is why each of these motion is described by an independent wave function

$$\Psi(x,y) = \underbrace{X(x)}_{\text{a wave function for x-motion}} \cdot \underbrace{Y(y)}_{\text{a wave function for y-motion}}$$

Each of these wave functions describes the particle behavior in two corresponding potential wells

along x: $U(x) = \begin{cases} 0 & 0 < x < L \\ \infty & \text{otherwise} \end{cases}$

$$U(y) = \begin{cases} 0 & 0 < y < L \\ \infty & \text{otherwise} \end{cases}$$

~~$\frac{d^2 \Psi}{dx^2} + \frac{d^2 \Psi}{dy^2} = -E \Psi$~~

We know solutions for these problems!

$$X_n(x) = \sqrt{\frac{2}{L}} \sin \frac{\pi n x}{L}$$

$n = 1, 2, \dots$

$$Y_s(y) = \sqrt{\frac{2}{L}} \sin \frac{\pi s y}{L}$$

$s = 1, 2, \dots$

$n, s = 1$

Two dimensional wave function

$$\Psi_{ns}(x,y) = \frac{2}{L} \sin \frac{\pi n x}{L} \sin \frac{\pi s y}{L}$$

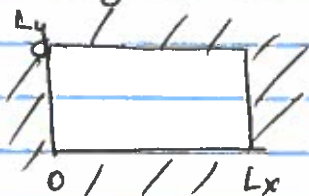
This is the eigenstate of the 2D Schrodinger eq corresponding eigenvalue:

$$E_{ns} = \frac{\pi^2 \hbar^2}{2mL^2} (n^2 + s^2)$$

$$\hat{H} \Psi_{ns}(x,y) = E_{ns} \Psi_{ns}(x,y)$$

Aside:

Easy to show that for (rectangular) non-square well



$$\Psi_{ns}(x,y) = \frac{2}{\sqrt{L_x L_y}} \sin \frac{\pi n x}{L_x} \sin \frac{\pi n y}{L_y}$$

$$E_{ns} = \frac{\pi^2 \hbar^2}{2m} \left[\frac{n^2}{L_x^2} + \frac{s^2}{L_y^2} \right]$$

Let's get back to a symmetric well

$$2D \quad E_{ns} = \frac{\pi^2 \hbar^2}{2mL^2} (n^2 + s^2)$$

$$3D \quad E_{nsj} = \frac{\pi^2 \hbar^2}{2mL^2} (n^2 + s^2 + j^2)$$

Energy spectrum	(possible particle energies)		
	1D	2D	3D
Ground state	$n=1 \quad E_1 = \frac{\hbar^2 \pi^2}{2mL^2}$	$n=1, s=1 \quad 2E_1$	$n=1, s=1, j=1 \quad 3E_1$
First excited state	$n=2 \quad 4E_1$	$n=1, s=2$ $s=2, s=1$ twice degenerate $5E_1$	$n=1, s=1, j=2$ and \odot 3x degenerate $6E_1$
Second excited state	$n=3 \quad 9E_1$		

After the first excited state we have to carefully analyze which combination gives lower energy

$s^2 \setminus n^2$	1	4	9	16
1	2	5	10	17
4	5	8	13	20
9	10	13	18	23
16	17	20	23	32

2D square well

(1,1) $E_1 = 2E$	(ground) singlet
(1,2) $E_2 = 5E$	2x degenerate
(2,2) $E_3 = 8E$	singlet
(1,3) $E_4 = 10E$	2-x (doublet)
(2,3) $E_5 = 13E$	— 4 —
(1,4) $E_6 = 17E$	— 4 —
(3,3) $E_7 = 18E$	singlet
(2,4) $E_8 = 20E$	doublet

Proper mathematical solution of the 2D Schrodinger eqn \rightarrow separation of variables

Inside the well $U(x,y) = 0$

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi(x,y)}{\partial x^2} + \frac{\partial^2 \psi(x,y)}{\partial y^2} \right) = E \psi(x,y)$$

As we discussed $\psi(x,y) = X(x) Y(y)$

$$-\frac{\hbar^2}{2m} \left(Y(y) \frac{d^2 X}{dx^2} + X(x) \frac{d^2 Y}{dy^2} \right) = E \cdot X \cdot Y$$

$$Y \cdot X'' + X \cdot Y'' + \frac{2mE}{\hbar^2} XY = 0 \quad \times \frac{1}{XY}$$

$$\left(\frac{X''}{X} + \frac{Y''}{Y} + \frac{2mE}{\hbar^2} \right) = 0$$

depends only on x depends only on y

must be constant

$$C_1 = -\frac{\hbar^2 k_x^2}{2m} \quad C_2 = -\frac{\hbar^2 k_y^2}{2m}$$

must be true for any x, y inside the well

$$\frac{X''}{X} = C_1 = -k_x^2$$

$$\frac{Y''}{Y} = C_2 = -k_y^2$$

$$X'' + k_x^2 X = 0$$

$$Y'' + k_y^2 Y = 0$$

general solution:

$$X(x) = A \cos k_x x + B \sin k_x x$$

$$Y(y) = \tilde{A} \cos k_y y + \tilde{B} \sin k_y y$$

Accounting for the boundary conditions (wavefunction vanishes at $x, y = 0$ & $x, y = L$)

$$X_n(x) = \sqrt{\frac{2}{L}} \sin \frac{\pi n x}{L}$$

$$Y_s(y) = \sqrt{\frac{2}{L}} \sin \frac{\pi s y}{L}$$

$$k_{x_n} = \frac{\pi n}{L}$$

$$k_{y_s} = \frac{\pi s}{L}$$

Substituting found values of k_x, s into
the Schrodinger equ

$$-k_x^2 - k_y^2 + \frac{2mE}{\hbar^2} = 0 \quad E = \frac{\hbar^2(k_x^2 + k_y^2)}{2m}$$

$$E_{n,s} = \frac{\pi^2 \hbar^2}{2mL^2} (n^2 + s^2)$$