

Maxwell's equation - crown jewel
of electromagnetism

Most of the equations were known
before, Maxwell put things together
and analysed the results

In ~~vacuum~~ ^{electric} ~~in the absence of~~
~~any charges~~

General form of Maxwell's eqn

$$\left\{ \begin{array}{l} \nabla \cdot \vec{D} = \rho \\ \nabla \cdot \vec{B} = 0 \\ \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \\ \nabla \times \vec{H} = \vec{j} + \frac{\partial \vec{D}}{\partial t} \end{array} \right. \quad \left\{ \begin{array}{l} \rho - \text{electric charge} \\ \quad \text{density} \\ \vec{j} - \text{electric current} \\ \quad \text{density} \end{array} \right.$$

\vec{E} - electric field } external
 \vec{B} - magnetic field } parameters
 $\vec{D} = \vec{E} + \vec{P}$

polarization, electrical response of
the material $\vec{P} = \epsilon_0 \chi \vec{E}$, $\vec{D} = \epsilon_0 \epsilon \vec{E}$

χ - susceptibility, ϵ - dielectric constant

$\vec{H} = \vec{B} / \mu \mu_0$ μ - permeability

For all materials we care about $\mu = 1$

$$\vec{H} = \vec{B} / \mu_0$$

$$\vec{\nabla} = \vec{e}_x \frac{\partial}{\partial x} + \vec{e}_y \frac{\partial}{\partial y} + \vec{e}_z \frac{\partial}{\partial z} \quad \text{gradient vector}$$

Looking at Maxwell's equations we see that electric field can originate from electrical charges (first eqn) or time-varying magnetic field (3rd eqn)

Magnetic field can originate from electrical currents or time-varying electric field. (4th eqn)

Electric field varying in time and space \rightleftharpoons Magnetic field varying in time and space

electro-magnetic wave

In vacuum, with no charges around

$$\begin{cases} \nabla \cdot \vec{E} = 0 & \nabla \times \vec{E} = -\partial \vec{B} / \partial t \\ \nabla \cdot \vec{B} = 0 & \nabla \times \vec{B} = \mu_0 \epsilon_0 \partial \vec{E} / \partial t \end{cases}$$

~~or~~
$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

or
$$\frac{\partial^2 \vec{E}}{\partial x^2} + \frac{\partial^2 \vec{E}}{\partial y^2} + \frac{\partial^2 \vec{E}}{\partial z^2} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

same for \vec{B}

Solution: $\vec{E} = \vec{E}_0 \cos(\vec{k} \cdot \vec{r} - \omega t + \varphi)$

\vec{k} - wave propagation direction

Previously we chose \vec{k} to point along z direction

$$\vec{E} = \vec{E}_0 \cos(kz - \omega t + \varphi)$$

$$\vec{B} = \vec{B}_0 \cos(kz - \omega t + \varphi)$$

We can now say more about the waves!

$$\frac{\partial^2 \vec{E}}{\partial z^2} = -k^2 \vec{E} \quad \frac{\partial^2 \vec{E}}{\partial t^2} = -\omega^2 \vec{E}$$

$$k^2 = \mu_0 \epsilon_0 \omega^2$$

$$c = \frac{\omega}{k} = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

Speed of light in vacuum

$$\nabla \cdot \vec{E} = 0 \Rightarrow \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0$$

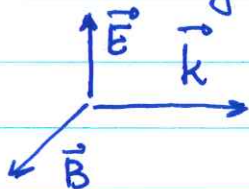
= 0 no dependence on x, y

E_z must be 0

$$\nabla \cdot \vec{B} = 0$$

B_z must be 0

Electromagnetic wave is a transverse wave



$$\vec{E} \perp \vec{B} \perp \vec{k}$$

\vec{E} - direction of the electric field determines the polarization of the light
For a linearly-polarized light we can define

$$E_x(z, t) = E_0 \cos(kz - \omega t + \varphi)$$

$$B_y(z, t) = B_0 \cos(kz - \omega t + \varphi)$$

$$B_0 = \frac{k E_0}{\omega} = \frac{E_0}{c}$$

What about materials?

If there are free charges around, electric field can move them, producing real work. Because of the energy conservation, that means that the energy of the electromagnetic field is reduced, and light is attenuated - absorbed.

That is why light cannot propagate inside metals (too many free electrons!)

In dielectrics - no free charges
 $\rho = 0$ $\vec{j} = 0$ $\vec{D} = \epsilon_0 \epsilon \vec{E}$
all response of bound charges in the atoms / molecules is hidden inside ϵ - dielectric constant.

Maxwell's eqns in a dielectric ($\mu=1$)
$$\left\{ \begin{array}{l} \nabla \cdot \vec{E} = 0 \\ \nabla \cdot \vec{B} = 0 \end{array} \right. \quad \left\{ \begin{array}{l} \nabla \times \vec{E} = -\partial \vec{B} / \partial t \\ \nabla \times \vec{B} = \mu_0 \epsilon_0 \epsilon \partial \vec{E} / \partial t \end{array} \right.$$

same as in vacuum, but

$$k^2 = \mu_0 \epsilon_0 \epsilon$$

$$v = \frac{1}{\sqrt{\mu_0 \epsilon_0 \epsilon}} = \frac{c}{\sqrt{\epsilon}} = \frac{c}{n}$$

Refractive index $n = \sqrt{\epsilon}$

Inside the dielectric the e-m wave propagates just as in vacuum, but slower, since now the electric field consists of not only the original wave, but also of re-radiated wave from excited bound electrons