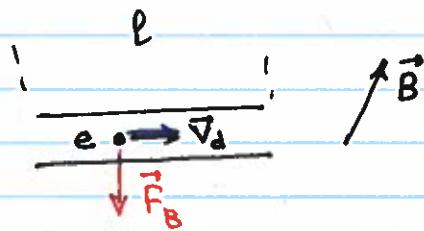


Magnetic force on a current-carrying conductor

electric

Since current is a stream of moving charges, and magnetic field exerts a force on each, there will be a total force on the ~~some~~ whole wire



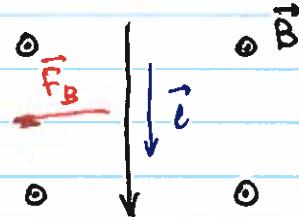
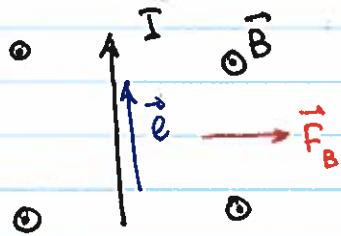
$$\vec{F}_{\text{total}} = N_{\text{charges}} \cdot e \vec{v}_d \times \vec{B}$$

$$N_{\text{charges}} = n \cdot A \cdot l$$

↑ length of the wire
 ↑ area of the wire
 charge density

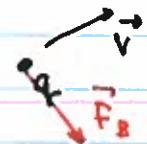
$$\vec{F} = n A l \cdot e \vec{v}_d \times \vec{B} = I \vec{l} \times \vec{B}$$

\vec{l} - vector length of the current-carrying wire, direction is along the current



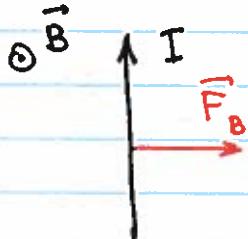
Magnetic force of a current-carrying conductors

\vec{B}



Moving charge circular / spiral motion

$$\vec{F} = q \vec{v} \times \vec{B}$$

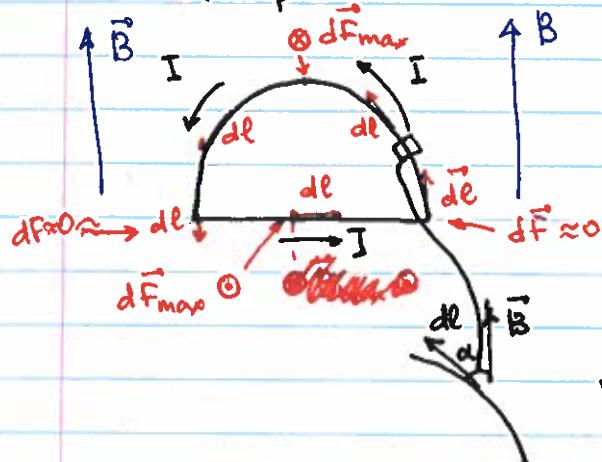


$$\vec{F}_B = I \cdot \vec{l} \times \vec{B}$$

l - the length of the wire in the direction of the current

If the wire is curved $d\vec{F}_B = I \cdot d\vec{l} \times \vec{B}$

Example



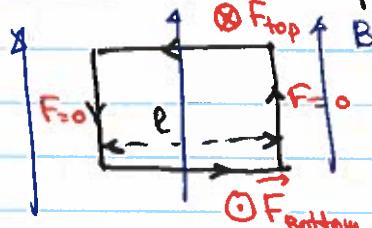
The magnetic force per unit length is strongest in the straight part and on top

$$|dF_{max}| = I dl \cdot B$$

It is zero near the bends
 $d\vec{l} \parallel \vec{B}$

$$|d\vec{F}| = I dl B \sin \theta < |d\vec{F}_{max}|$$

A closed loop in



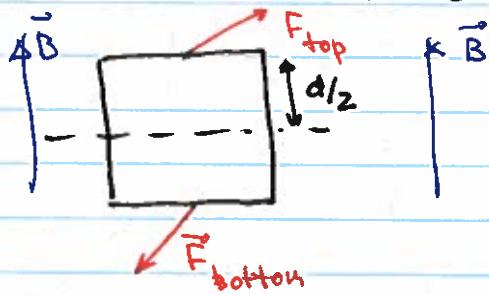
a magnetic field

$$|\vec{F}_{top}| = |\vec{F}_{bottom}| = l \cdot I \cdot B$$

but in opposite directions

$$F_{tot} = 0 \text{, no net force}$$

But! there is torque!



Each force tries to rotate the loop in the same direction

$$\begin{aligned} \tau &= F_{top} \cdot \frac{d}{2} + F_{bottom} \cdot \frac{d}{2} = \\ &= \pi B \cdot d = I(l \cdot d) \cdot B \end{aligned}$$

area

One can show that for any shape
of the flat current loop

$$\vec{\tau} = I \cdot \vec{A} \times \vec{B} = \vec{\mu} \times \vec{B}$$

$\vec{\mu} = I \cdot \vec{A}$ magnetic moment for a
 $U_B = -\vec{\mu} \cdot \vec{B}$ current loop

Many atoms have internal magnetic moment
due to electron(s) orbiting the nucleus.
This magnetic moment tends to be proportional
to the orbital angular momentum \vec{L} : $\vec{\mu} = g_L \vec{L}$
(names s-orbital — angular momentum 0
p-orbital — 1
d-orbital — 2 ...)

That is why atomic energies change
in magnetic field

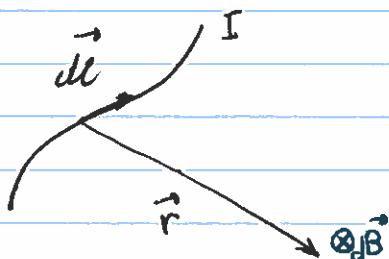
Elementary particles often have intrinsic
magnetic moment — spin
The name "spin" comes from incorrect
early hypothesis that electron and proton
magnetic moment comes from their own
rotation, but in reality spin origin
is relativistic.

Sources of magnetic field

Magnetic field acts on moving charges and currents

\Leftrightarrow Moving charges and electric currents produce magnetic field

Bio-Savart law



$$d\vec{B} = \frac{\mu_0}{4\pi} I \frac{dl \times \hat{r}}{r^2}$$

$$\mu_0 = 4\pi \cdot 10^{-7} \frac{\text{T} \cdot \text{m}}{\text{A}}$$

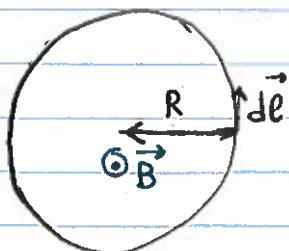
Magnetic field contribution of current element dl

permeability of free space
(a constant needed to connect randomly defined SI units)

$$\vec{B} = \int_{\text{along the wire}} d\vec{B} = \frac{\mu_0 I}{4\pi} \int \frac{dl \times \hat{r}}{r^2}$$

Geometry of the wire

The calculations are the simplest in the plane of the ~~the~~ current, since in this case \vec{B} is perpendicular to this plane



$$d\vec{B} = \frac{\mu_0}{4\pi} I \frac{dl}{R^2} \cdot \hat{z} \quad (\text{out of the page})$$

$$\vec{B} = \oint d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{1}{R^2} \oint \frac{dl}{2\pi R} \hat{z}$$

~~$$\vec{B} = \frac{\mu_0 I}{2\pi R} \hat{z}$$~~