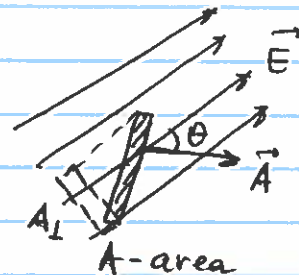
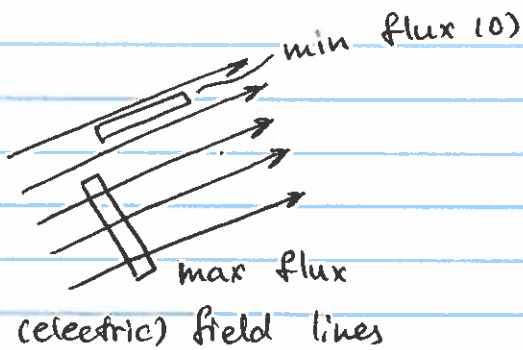


Gauss law for electric field

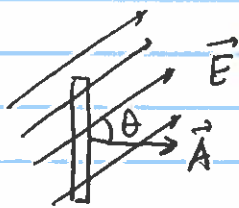
The concept of a flux



\vec{A} - area vector
 $|\vec{A}| = A$
 direction - normal to the surface

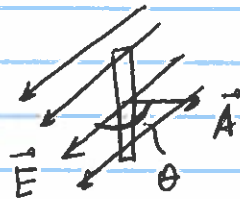
$$\Phi = \vec{E} \cdot \vec{A} = |\vec{E}| |\vec{A}| \cos \theta$$

$\underbrace{|\vec{A}|}_{A_{\perp}}$

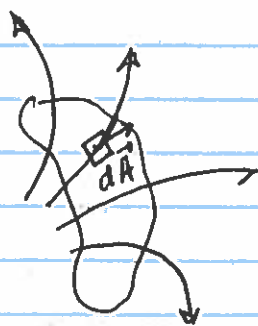


$\theta < 90^\circ$
 $\Phi = |\vec{A}| |\vec{E}| \cos \theta > 0$

$90^\circ < \theta < 180^\circ$
 $\Phi = |\vec{A}| |\vec{E}| \cos \theta < 0$



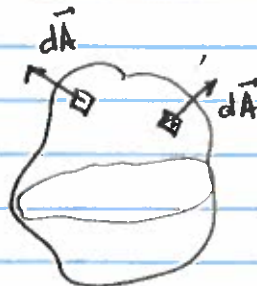
If electric field changes in space



$$d\Phi = \vec{E} d\vec{A}$$

$$\Phi_{\text{tot}} = \int_{\text{surface area}} \vec{E} d\vec{A}$$

If the surface is closed



$$\Phi_{\text{tot}} = \oint \vec{E} d\vec{A}$$

$d\vec{A}$ points outside

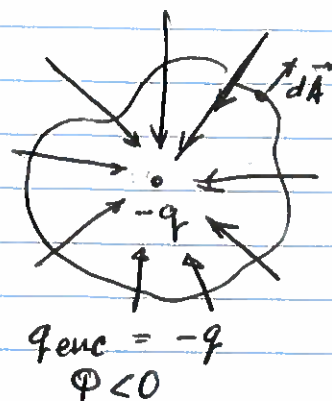
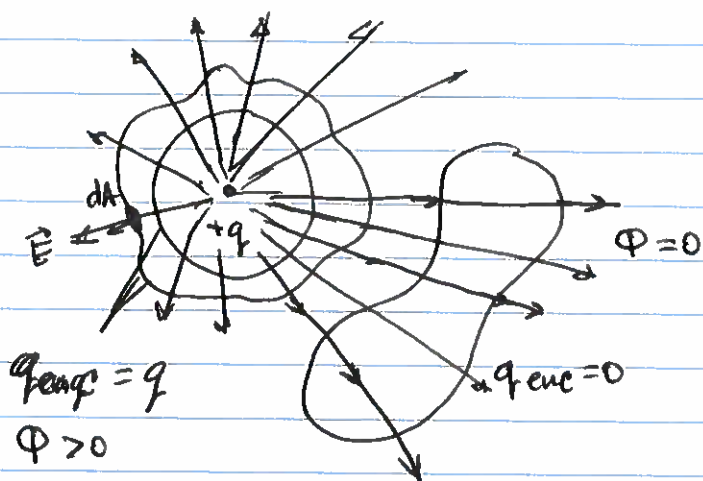
Gauss law

The flux of electric field \vec{E} through any closed surface is equal to the net charge enclosed inside the surface, divided by ϵ_0 .

$$\Phi = \oint \vec{E} \cdot d\vec{A} = \frac{q_{enc}}{\epsilon_0} \quad \left[k = \frac{1}{4\pi\epsilon_0} \right]$$

ϵ_0 - dielectric permittivity of vacuum
(purely historic artefact name)

$\epsilon_0 = 8.84 \cdot 10^{-12}$ F/m (F - Farad, wait till we get to capacitors)



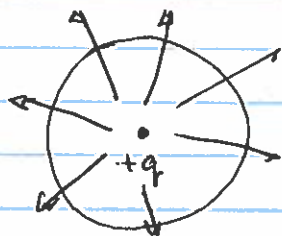
Gauss law is always valid, and is the first out of four main equations of electromagnetic theory - Maxwell's equations.

However, sometimes we can use the symmetry of the problem to find electric field. In this case we need to figure out a surface, along which E -field is constant, so then

$$\int \vec{E} \cdot d\vec{A} \rightarrow E \cdot \int dA = E \cdot A$$

Point charge

Electric field points radially away from positive charge (or into a negative charge)
 $|\vec{E}|$ must be the same at the given distance



If we take a sphere, centered at q , \vec{E} -field will be pointing perpendicular to this sphere everywhere, and E -field magnitude is constant

$$d\Phi = \vec{E} \cdot d\vec{A} = E \cdot dA$$

$$\Phi = \int d\Phi = E \int dA = E \cdot A = E \cdot 4\pi r^2 = \frac{q}{\epsilon_0}$$

$$E = \frac{q}{\epsilon_0} \frac{1}{4\pi r^2} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} = \frac{kq}{r^2}$$

of the sphere

$$\vec{E} = \frac{kq}{r^2} \hat{r}$$

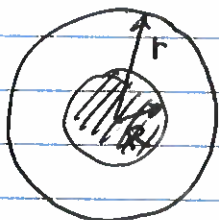
Here we used our physical intuition and the symmetry of the problem to provide extra information about \vec{E} -field direction.

Uniformly charged sphere

Total charge $+Q$, radius R

volume charge density $\rho = Q / \left(\frac{4\pi}{3} R^3 \right) = \frac{3Q}{4\pi R^3}$

Two distinct regions



outside: $r > R$

$$q_{enc} = Q$$

$$\Phi_{tot} = E \cdot 4\pi r^2 = \frac{Q}{4\pi\epsilon_0 r^2}$$

$$E(r > R) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \quad \text{same as point charge}$$



inside: $r < R$

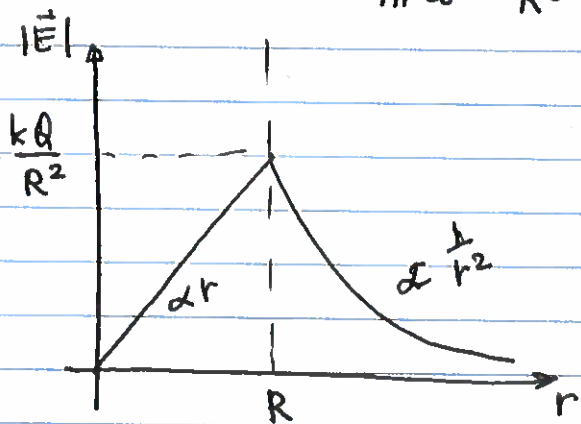
$$q_{enc} < Q \quad q_{enc} = \rho \cdot \frac{4\pi r^3}{3}$$

$$= Q \frac{\frac{4\pi}{3} r^3}{\frac{4\pi}{3} R^3} = Q \frac{r^3}{R^3}$$

$$\Phi_{inside} = |\vec{E}| \cdot 4\pi r^2 = q_{enc} / \epsilon_0 = \frac{1}{\epsilon_0} Q \frac{r^3}{R^3}$$

$$|\vec{E}(\vec{r})|^2 = \frac{1}{4\pi\epsilon_0} \frac{r}{R^3}$$

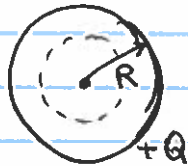
$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{\vec{r}}{R^3}$$



Thin charged shell

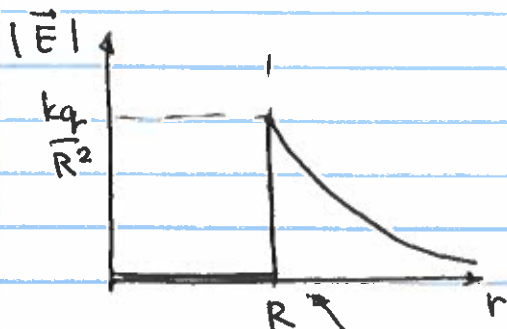
Outside $r > R$ $q_{enc} = Q$

$$\vec{E}_{outside} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$



Inside $r < R$ $q_{enc} = 0$

$$\vec{E}_{inside} = 0$$

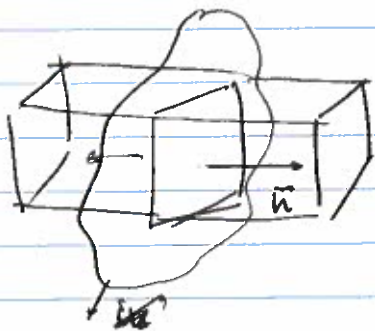


the discontinuity indicates the location of the charge

Electric field inside any charged sphere is zero because of Gauss's law!

Uniformly charged plane

constant charge density σ

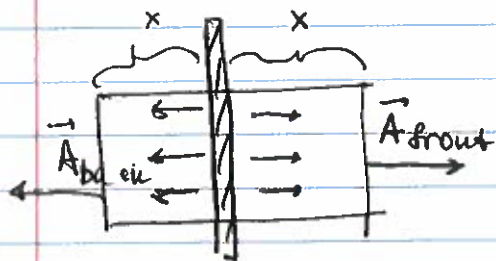


If the plane is infinite,
 \vec{E} -field can only point
 perpendicular to the plane.

Good choice for Gaussian
 surface:

cube or cylinder with sides
 parallel to the plane

Top view



$$q_{enc} = \sigma \cdot A$$

$$\Phi_{tot} = \Phi_{sides} + \Phi_{front} + \Phi_{back}$$

$$\Phi_{front} = \cancel{q_{enc} A} \vec{E} \cdot \vec{A}_{front} = E \cdot A$$

$$\Phi_{back} = \cancel{q_{enc} A} \vec{E} \cdot \vec{A}_{back} = E \cdot A$$

$$\Phi_{sides} = 0$$

$$\Phi_{tot} = 2E \cdot A = \frac{q_{enc}}{\epsilon_0} = \frac{\sigma \cdot A}{\epsilon_0} \quad E = \frac{\sigma}{2\epsilon_0}$$

Note, this is the same electric
 field we calculated when near
 the surface of a charged disk.