

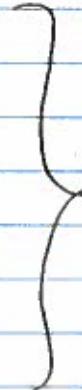
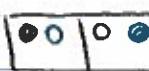
Microscopic definition of entropy

We mentioned that entropy characterizes disorder in a system.

Each macroscopic state can be realized with a number of microscopic configurations.
 Example: 4 beads in 2 boxes
 $\#1 \quad \#2$



- 1 configuration with #2 empty



6 configurations
 with equal numbers
 in both sides



- 1 configuration with #1 empty

W - statistical weight (not work, just the same letter)

$$W_{4-0} = 1$$

$$W_{2-2} = 6$$

$$W_{0-4} = 1$$

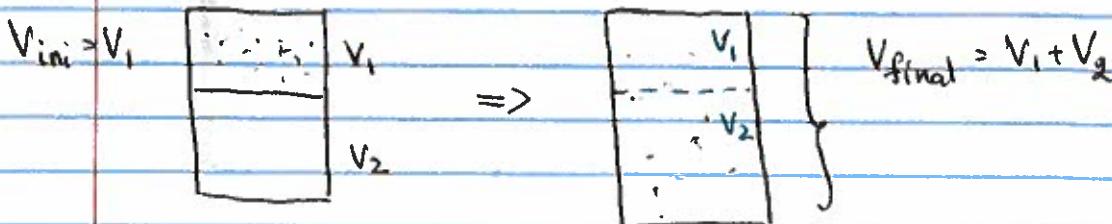
Entropy $S = k_B \ln W$ Boltzmann definition

So $S_{4-0} = 0$ and $S_{0-4} = 0$ but $S_{2-2} = k_B \ln 6$

The quote I really like: Boltzman showed that the law of physics are mostly accurate.

Arrow of time appears not because of fundamental laws of physics, but because every ^{complex} system evolves from ~~more~~ more ordered (^{less} probable) to less ordered (more probable) state \rightarrow entropy grows

We calculated the entropy growth when gas expands using macroscopic description. We can do the same with microscopic description as well.



Probability of a single molecule

For a given volume V_i let's break it into small "cells" of volume V_{min} , and count how many possible "locations" a single molecule can have

$$W_i = \frac{V_i}{V_{\text{min}}}$$

If we have N molecule, each one has the same # of options, so

$$W_i = \left(\frac{V_i}{V_{\text{min}}} \right)^N$$

$$S_i = k_B \ln W_i = k_B \cdot N \ln \frac{V_i}{V_{\text{min}}}$$

If volume is increased $\rightarrow w_1 = \frac{V_{fin}}{V_{min}}$

$$W_{fin} = \left(\frac{V_{fin}}{V_{min}} \right)^N$$

$$S_{fin} = k_B \ln W_{fin} = k_B N \ln \frac{V_f}{V_{min}}$$

$$\Delta S = S_{fin} - S_{ini} = k_B \cdot N \left(\ln \frac{V_f}{V_{min}} - \ln \frac{V_{ini}}{V_{min}} \right) = k_B N \cdot \ln \frac{V_f}{V_i}$$

same as $\Delta S = nR \ln \frac{V_f}{V_i}$ we got last time