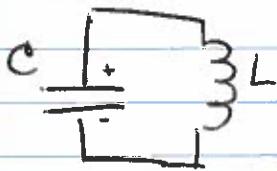


Written assignment #2 solutions

Q₁



$$V_C = \frac{Q}{C}$$

$$V_L = -L \frac{dI}{dt} = L \frac{d^2Q}{dt^2}$$

$$V_C + V_L = 0$$

$$L \frac{d^2Q}{dt^2} + \frac{Q}{C} = 0$$

a) $Q(t) = Q_0 \cos \omega_0 t \quad -L\omega_0^2 + \frac{1}{C} = 0$

$$I(t) = -\omega_0 Q_0 \sin \omega_0 t$$

$$\omega_0 = \frac{1}{\sqrt{LC}}$$

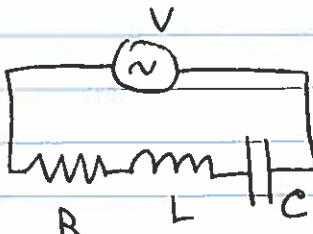
does not depend on the initial charge

b) $L \cdot C = \frac{1}{\omega_0^2} = \frac{1}{(2\pi)^2 (10^7 \text{ Hz})^2} \approx 2.5 \cdot 10^{-16} \text{ } \text{H}^{-2}$

Typical capacitors — $10 \mu\text{F}$ (10^{-11}F) to 1mF (10^{-3}F)
 Typical inductances — $0.1 \mu\text{H}$ (10^{-7}F) to 10mH (10^{-2}H)

(For example $C = 50 \mu\text{F}$ and $L = 5 \mu\text{H}$ will work)

Q₂



$$V(t) = V_0 e^{i\omega t} \quad I(t) = I_0 e^{i\omega t}$$

(V_0, I_0 can be complex!)

a) $V_R = -IR, \quad V_L = -L \frac{dI}{dt} = -iL\omega I_0 e^{i\omega t}$

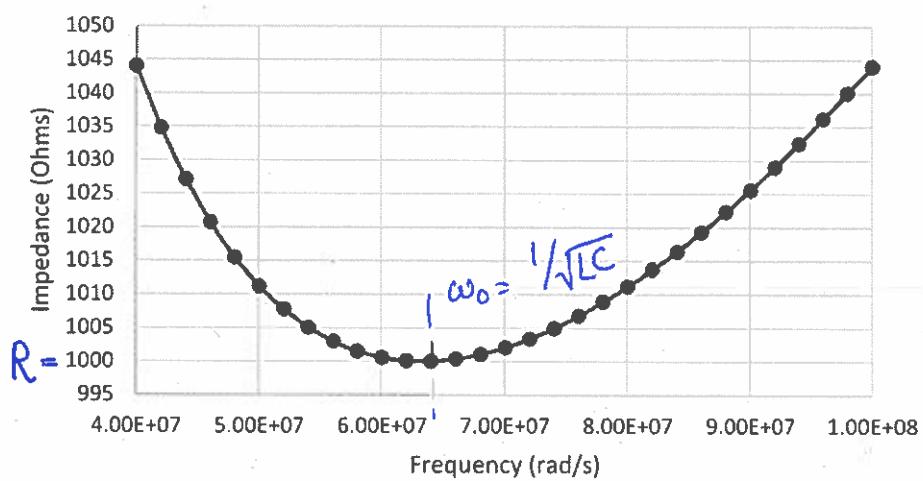
$$V_C = \frac{Q}{C} = -\frac{1}{C} \int I_0 e^{i\omega t} dt = -\frac{1}{i\omega C} I_0 e^{i\omega t}$$

$$V(t) = V_0 e^{i\omega t} - I_0 R e^{i\omega t} - i\omega L I_0 e^{i\omega t} + \frac{1}{i\omega C} I_0 e^{i\omega t} = 0$$

$$V(t) = \underbrace{\left(R + i\omega L + \frac{1}{i\omega C} \right)}_{Z(\omega)} I(t)$$

$$|Z(\omega)| = \sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2}$$

AC circuit impedance as a function of frequency



Q3

$$\frac{d^2x}{dt^2} + \frac{2}{\tau} \frac{dx}{dt} + \omega_0^2 x = -\frac{eE_0}{2m} e^{-i\omega t}$$

a) $x = x_0 e^{-i\omega t}$

$$-\omega^2 x_0 e^{-i\omega t} + \frac{2}{\tau}(-i\omega) x_0 e^{-i\omega t} + \omega_0^2 x_0 e^{-i\omega t} = -\frac{eE_0}{2m} e^{-i\omega t}$$

$$[(\omega_0^2 - \omega^2) \cancel{\star} + \frac{2i\omega}{\tau}] x_0 = -\frac{eE_0}{2m}$$

$$x(t) = -\frac{eE_0}{2m} e^{-i\omega t} \frac{1}{(\omega_0^2 - \omega^2) + 2i\omega/\tau}$$

b) if $\omega_0 \approx \omega$ $(\omega_0^2 - \omega^2) = (\omega_0 - \omega)(\omega_0 + \omega) \approx 2\omega(\omega_0 - \omega)$

$$x(t) \approx -\frac{eE_0 \cancel{\star}}{4m\omega} e^{-i\omega t} \frac{1}{(\omega_0 - \omega) - i/\tau} = -\frac{i e E_0 \tau e^{+i\omega t}}{1 + i(\omega_0 - \omega)\tau}$$

$$d(t) = -e x(t) = +\frac{1}{4m\omega} \frac{i e^2 E_0 \tau}{1 + i(\omega_0 - \omega)\tau} e^{-i\omega t} = \frac{i e^2 E_0 \tau e^{-i\omega t}}{4m\omega (1 + i(\omega_0 - \omega)\tau)}$$

Q4

$$x(w) = \frac{e_0 e^2 N \tau}{4\pi m \omega} \frac{1}{1 + i(\omega_0 - \omega)\tau} = \frac{e_0 e^2 N \tau}{4\pi m \omega} \frac{(1+i)(1-i\tau(\omega_0 - \omega))}{1 + \tau^2(\omega_0 - \omega)^2}$$

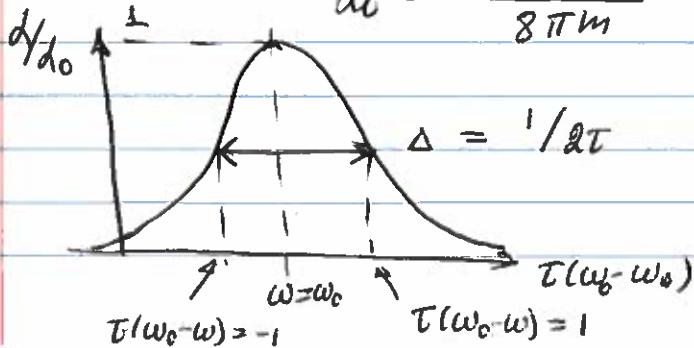
$$= \frac{e_0 e^2 N \tau}{4\pi m \omega} \frac{1 + i\tau(\omega_0 - \omega)}{1 + \tau^2(\omega_0 - \omega)^2}$$

a) Absorption $d(\omega) = +\frac{1}{\lambda} x(\omega) = +\frac{\omega}{2\pi c} \cdot \frac{e_0 e^2 N \tau}{4\pi m \omega} \frac{(1+i)(1-i\tau(\omega_0 - \omega))}{1 + \tau^2(\omega_0 - \omega)^2}$

$$d(\omega) = \frac{e_0 e^2 N \tau}{8\pi m} \frac{1}{1 + \tau^2(\omega_0 - \omega)^2}$$

max absorption $\omega = \omega_0$ (on resonance)

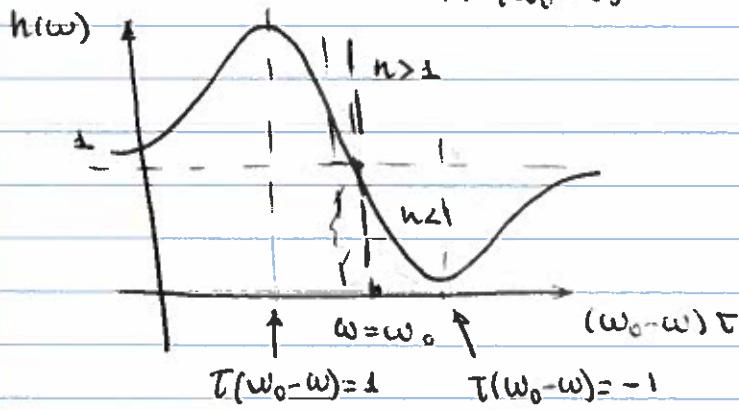
$$d_0 = \frac{e_0 e^2 N \tau}{8\pi m}$$



The longer is the lifetime of an excited atomic state, the narrower is the absorption resonance.

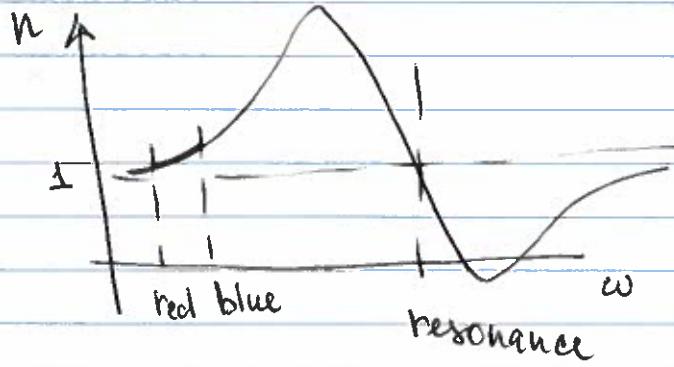
$$b) n(\omega) = 1 + \operatorname{Re}(\chi(\omega)) = 1 + \frac{\epsilon_0 e^2 N \tau}{4 \pi \omega} \frac{\tau(\omega_0 - \omega)}{1 + \tau^2 (\omega_0 - \omega)^2} =$$

$$= 1 + \frac{\lambda d_0}{1 + (\omega_0 - \omega)^2 \tau^2}$$



outside of absorption
the dispersion is
positive $dn/d\omega > 0$

$$c) n_{\text{red}} = 1.5, n_{\text{blue}} = 1.54$$



Since $1 < n_{\text{red}} < n_{\text{blue}}$
we expect to
find resonances
in glass in UV