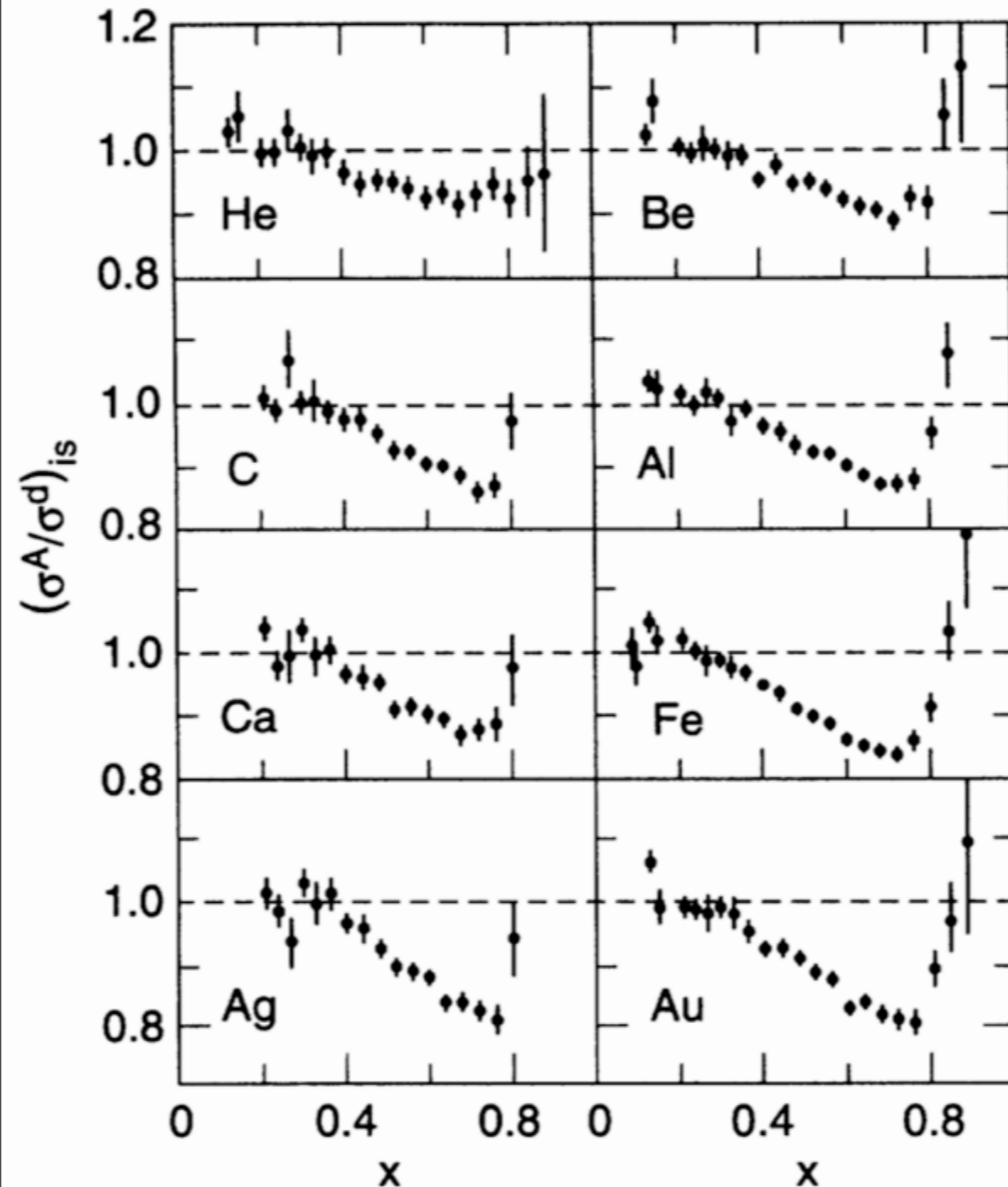




Tagging the EMC Effect with Slow and Fast Protons

K. Griffioen
College of William & Mary
griff@physics.wm.edu

Gordon Research Conference on
Photonuclear Reactions
Tilton, New Hampshire
5 August 2010

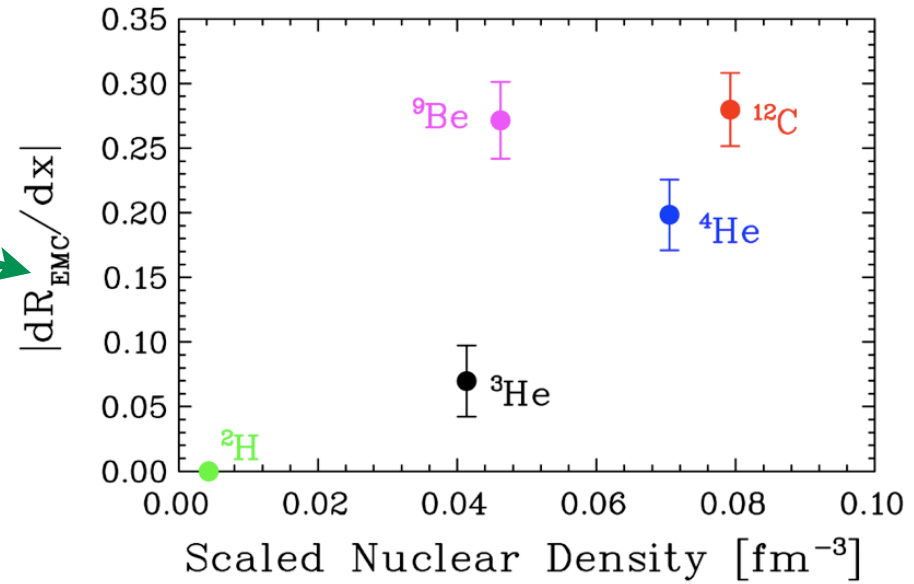
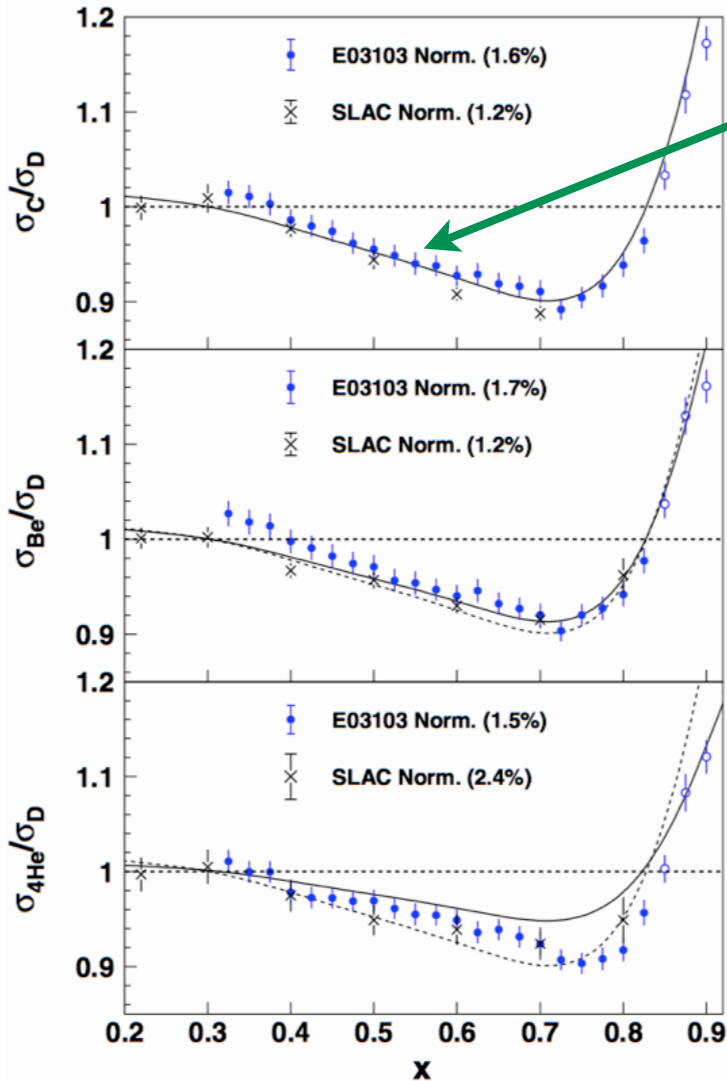


Gomez et al, PRC49(94)4348

- European Muon Collaboration (EMC) discovered a reduced deep-inelastic cross section for $0.4 < x < 0.7$ in nuclei compared to deuterium (Aubert et al., PLB120(83)275)
- SLAC (see plot) showed the effect grows with A and depends almost linearly with x .
- Effect is still not understood, but the likely explanation is a modification of the bound nucleon structure function.



Seely et al, PRL 103(09)202301



- Recent data from JLab are very precise.
- Slope with respect to x is used to characterize the strength of the effect
- ^9Be anomaly is used to argue for a local-density origin of the effect

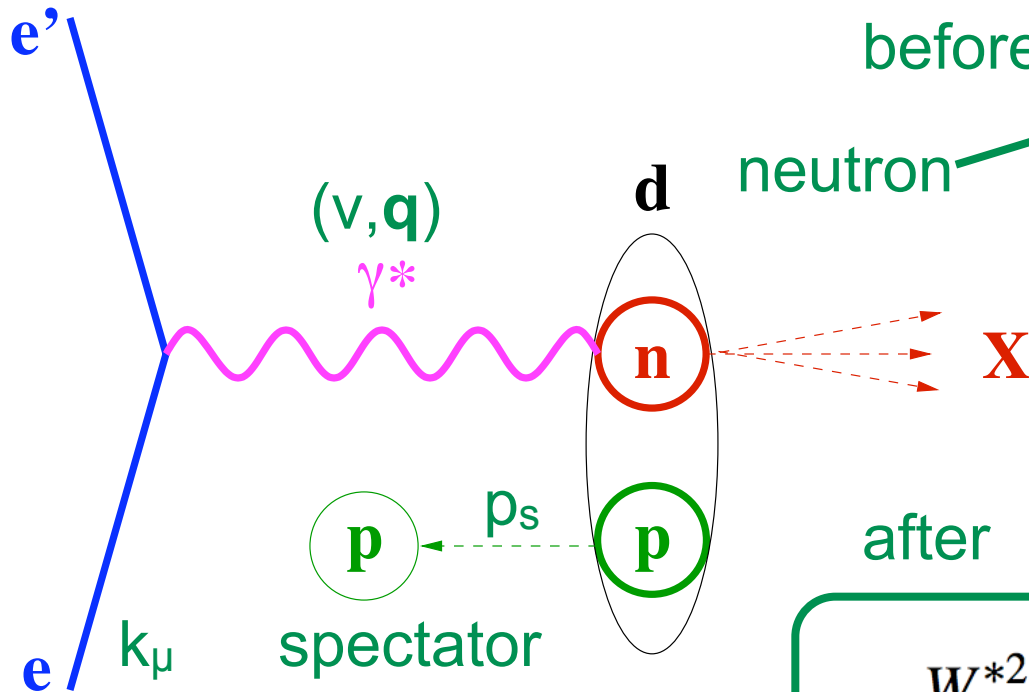


- BoNuS
- Bound Nucleon Structure Experiment
- $d(e, e'p_s)X$ [(deep) inelastic]
- deuterium target, spectator proton
- $70 < p_s < 180$ MeV/c
- JLab Hall B CLAS with an RTPC
- search for F_2^n at high x

- DEEPS
- $d(e, e'p_s)X$ [deep inelastic]
- deuterium target, spectator proton
- $250 < p_s < 600$ MeV/c
- JLab Hall B CLAS
- search for $F_2^{n\{\text{eff}\}}$ for $0.3 < x < 0.7$



$d(e, e' p_s) X$



$$p_N^\mu = (M_d - E_s, -\vec{p}_s)$$

$$E_p + E_n = M_d$$

$$E^* = M_d - \sqrt{M_s^2 + p_s^2}$$

$$M^{*2} = (M_d - E_s)^2 - \vec{p}_s^2$$

$$W^{*2} \approx M^{*2} - Q^2 + 2Mv(2 - \alpha_s)$$

$$\alpha_s = \frac{E_s - p_{s\parallel}}{M_s} \quad y^* = \frac{p_N^\mu q_\mu}{p_N^\mu k_\mu} \approx y$$

$$x^* = \frac{Q^2}{2p_N^\mu q^\mu} \approx \frac{Q^2}{2Mv(2 - \alpha_s)} = \frac{x}{2 - \alpha_s}$$

- plane-wave impulse approximation
- backward-emitted p is spectator
- struck neutron is off-shell
- momenta are equal and opposite
- Lorentz invariants are corrected for initial neutron 4-momentum



$$\frac{d\sigma}{dx^* dQ^2} = \frac{4\pi\alpha_{\text{EM}}^2}{x^* Q^4} \left[\frac{y^{*2}}{2(1+R)} + (1-y^*) + \frac{M^{*2} x^{*2} y^{*2}}{Q^2} \frac{1-R}{1+R} \right] F_2(x^*, \alpha_s, p_T, Q^2) \times S(\alpha_s, p_T) \frac{d\alpha_s}{\alpha_s} d^2 p_T,$$

Cross Section

Off-Shell F_2

Light Cone

Spectral Function

Nonrelativistic w.f.

$$P(\vec{p}_s) = J |\psi_{\text{NR}}(p_s)|^2$$

$$J = 1 + \frac{p_{s\parallel}}{E_n^*} = \frac{(2 - \alpha_s) M_d}{2(M_d - E_s)}$$

$$S(\alpha_s, p_T) \frac{d\alpha_s}{\alpha_s} d^2 p_T = P(\vec{p}_s) d^3 p_s$$

$$S^{\text{LC}}(\alpha_s, p_T) \frac{d\alpha_s}{\alpha_s} d^2 p_T = |\psi_{\text{NR}}(|\vec{k}|^2)|^2 d^3 k$$

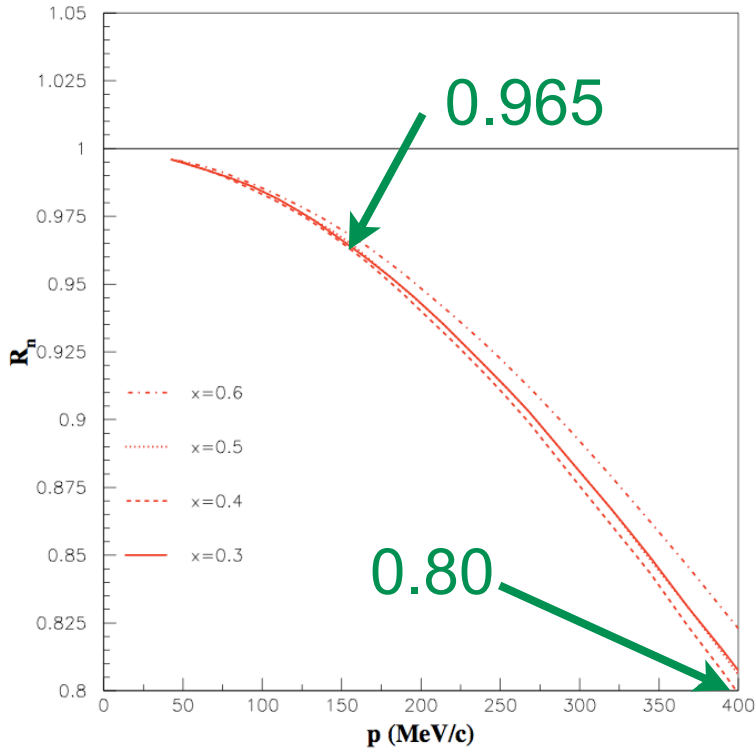
$$|\vec{k}| = \sqrt{\frac{M^2 + p_T^2}{\alpha_s(2 - \alpha_s)} - M^2} \quad \alpha_s = 1 - \frac{k_{\parallel}}{\sqrt{M^2 + \vec{k}^2}}$$

$$k_0 = \sqrt{M^2 + \vec{k}^2} \quad \vec{p}_T = \vec{k}_T$$

$$\int \int \int S^{\text{LC}}(\alpha_s, p_T) \frac{d\alpha_s}{\alpha_s} d^2 p_T = 1$$



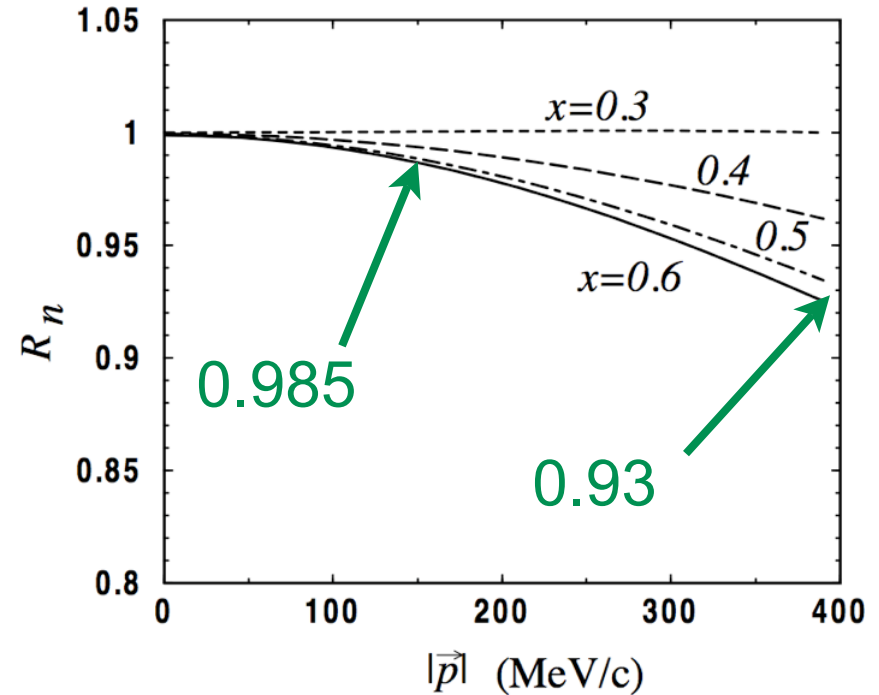
Liuti & Gross PLB356(95)157



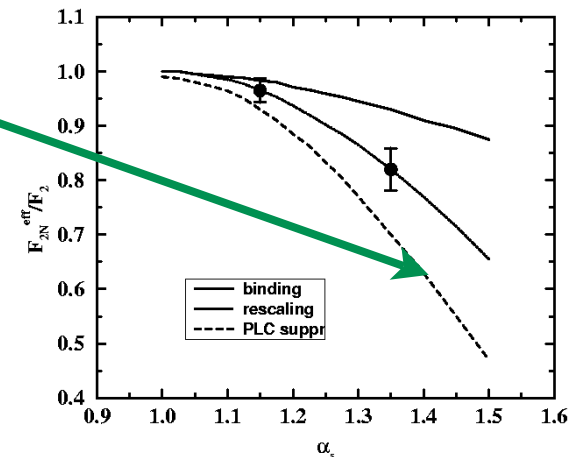
$$R_n \equiv (F_2^n)^{\text{eff}} / (F_2^n)^{\text{free}}$$

- R_n decreases with p_s or α
- at $x^*=0.5$ and $p_s=0.40$ GeV/c, R_n deviates from 1 by 7-20% in these models
- at $\alpha=1.4$ the deviation from unity could be 40%

Melnitchouk et al, PLB335(94)11

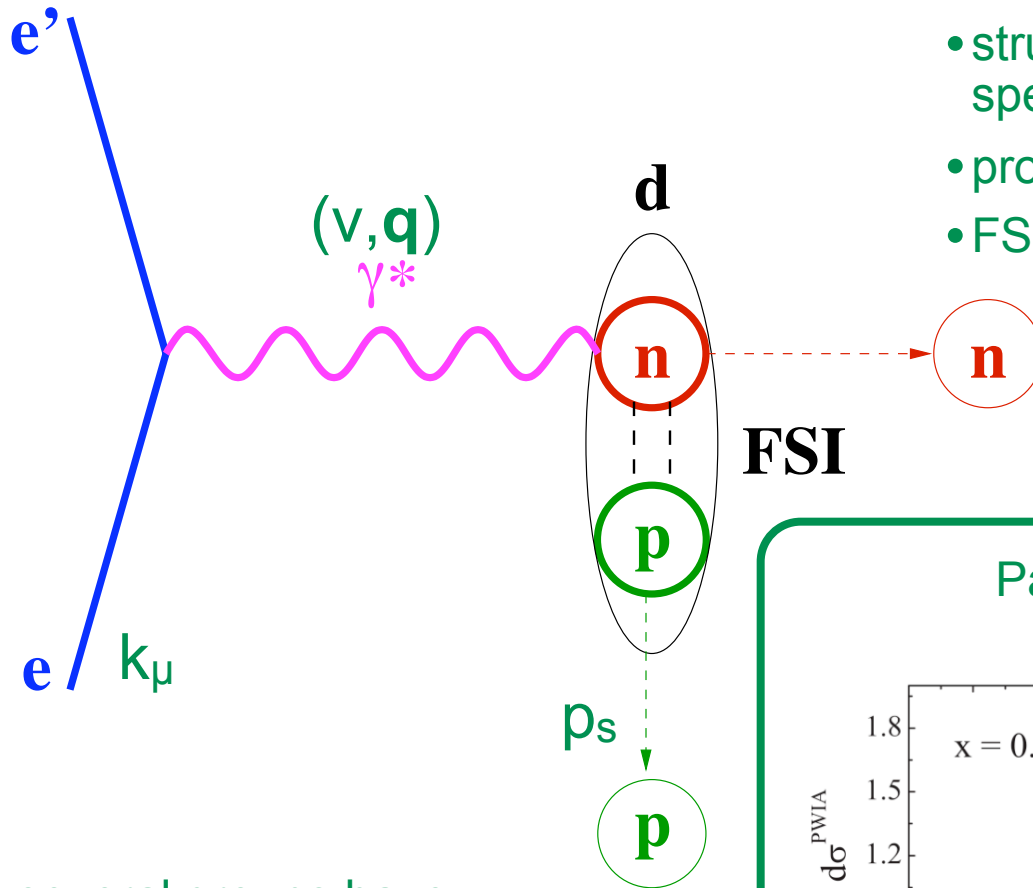


0.6





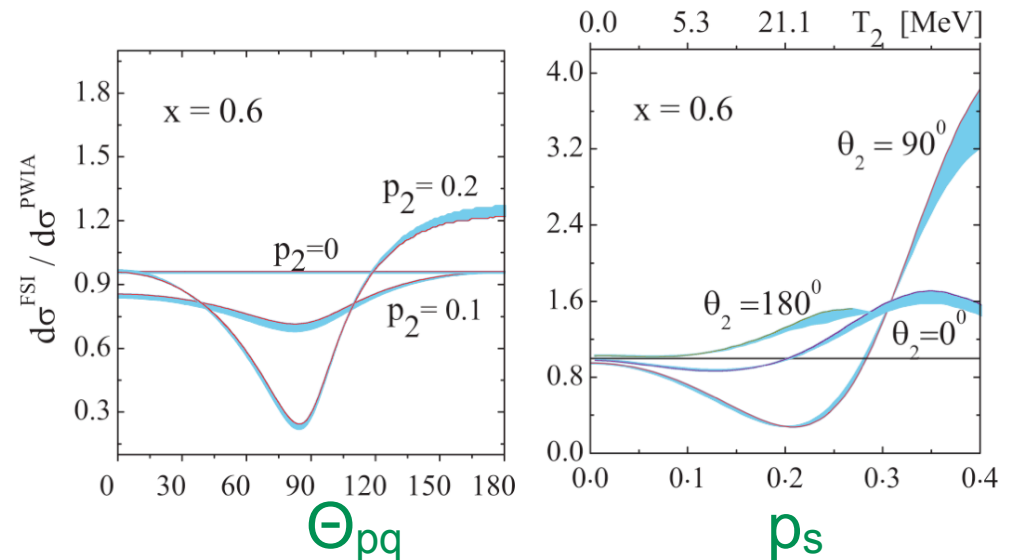
Final State Interactions



- struck neutron can interact with the spectator proton
- proton momentum is enhanced
- FSIs are small at low p_s and large Θ_{pq}

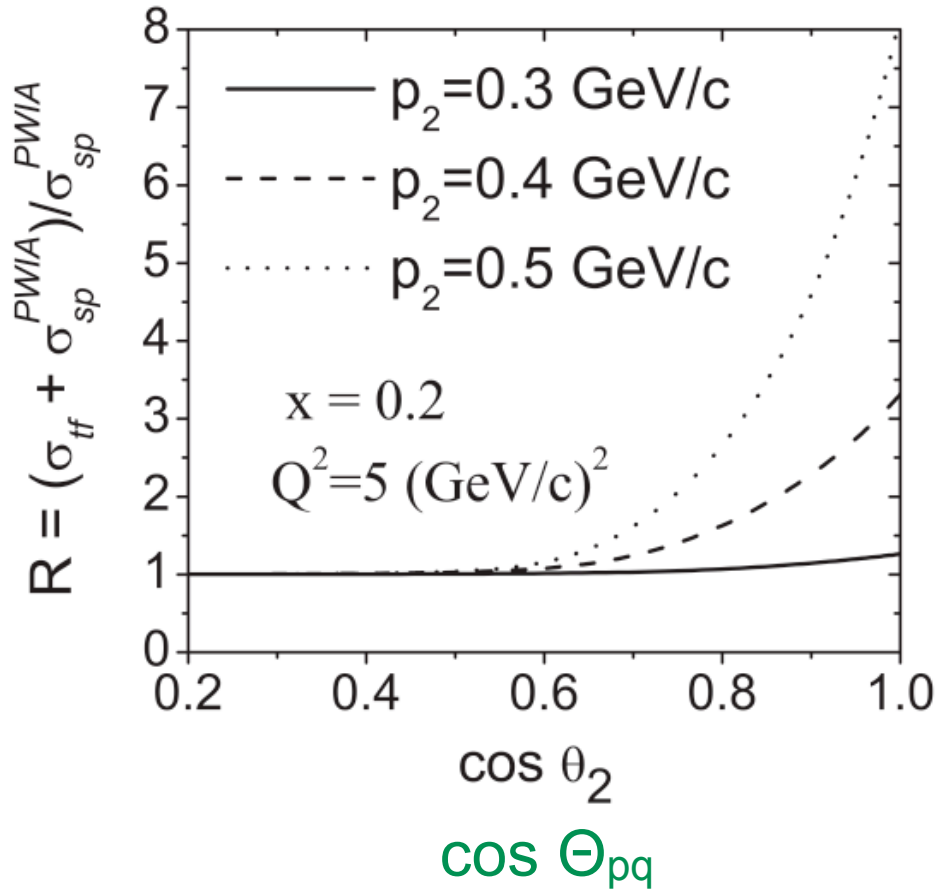
- several groups have calculated FSIs
- see e.g. posters by Wim Cosyn & Sabine Jeschonnek
- $\Theta_{pq} > 120^\circ$ minimizes FSIs

Palli et al, PRC80(09)054610

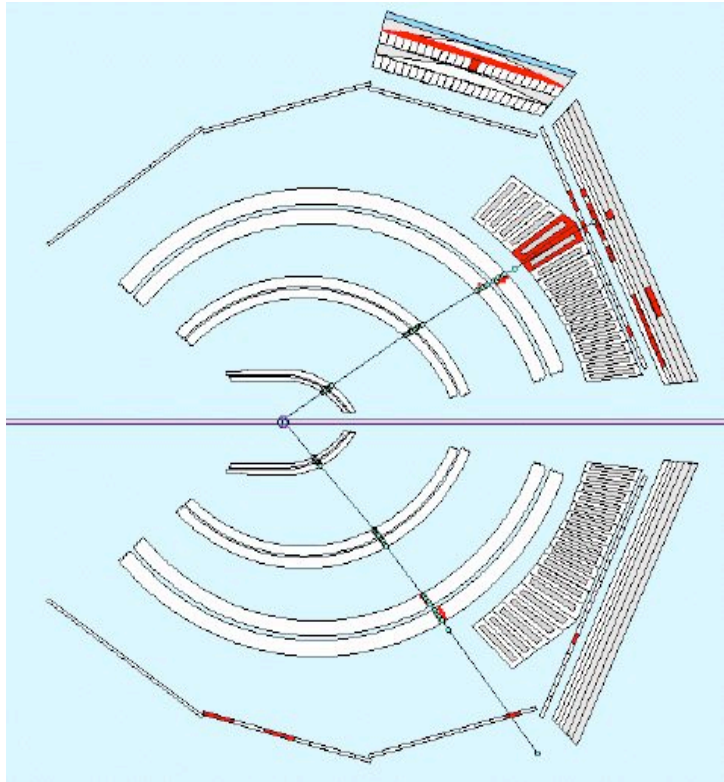




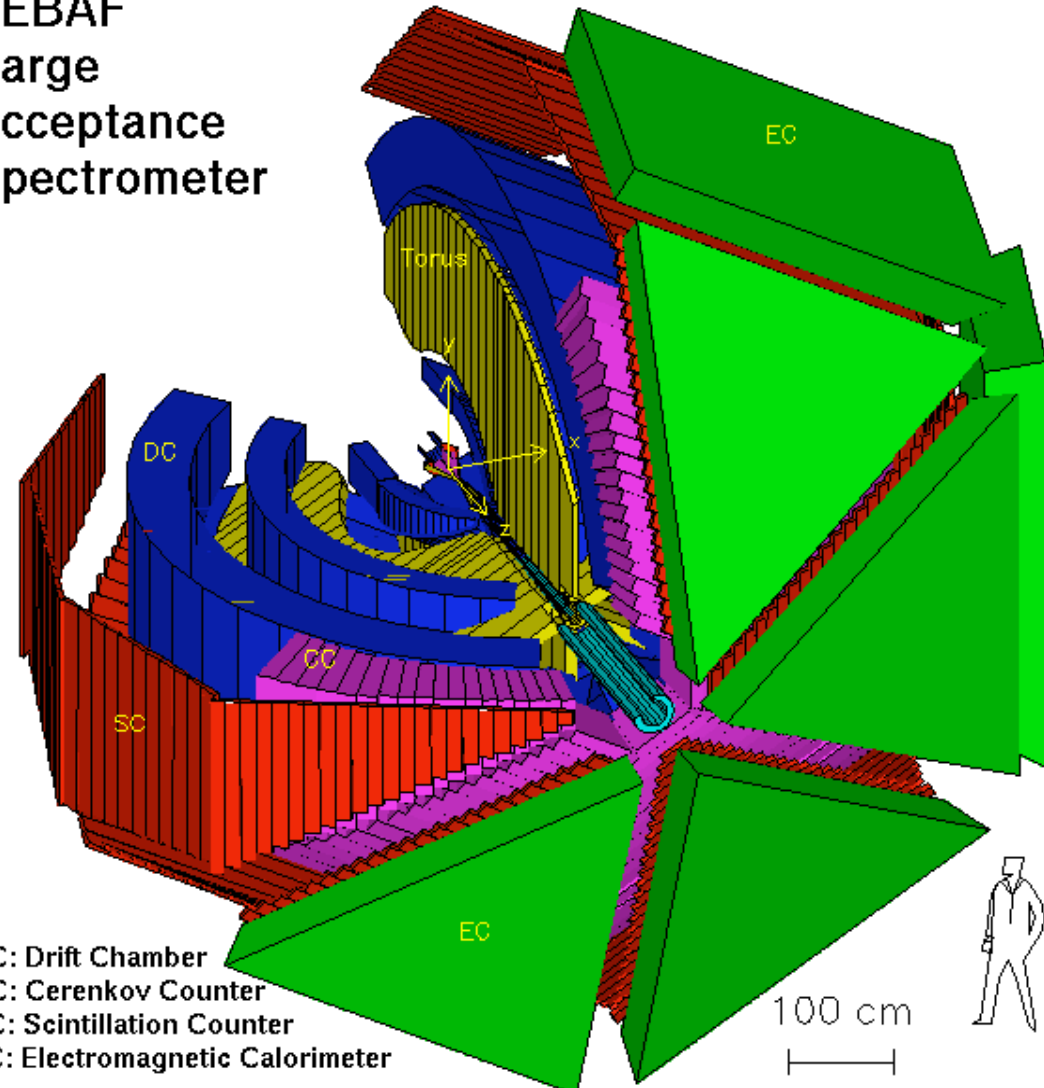
Palli et al, PRC80(09)054610



- target fragmentation enhances the proton yield only at forward angles ($\cos \Theta_{pq} > 0.6$)
- this can be ignored



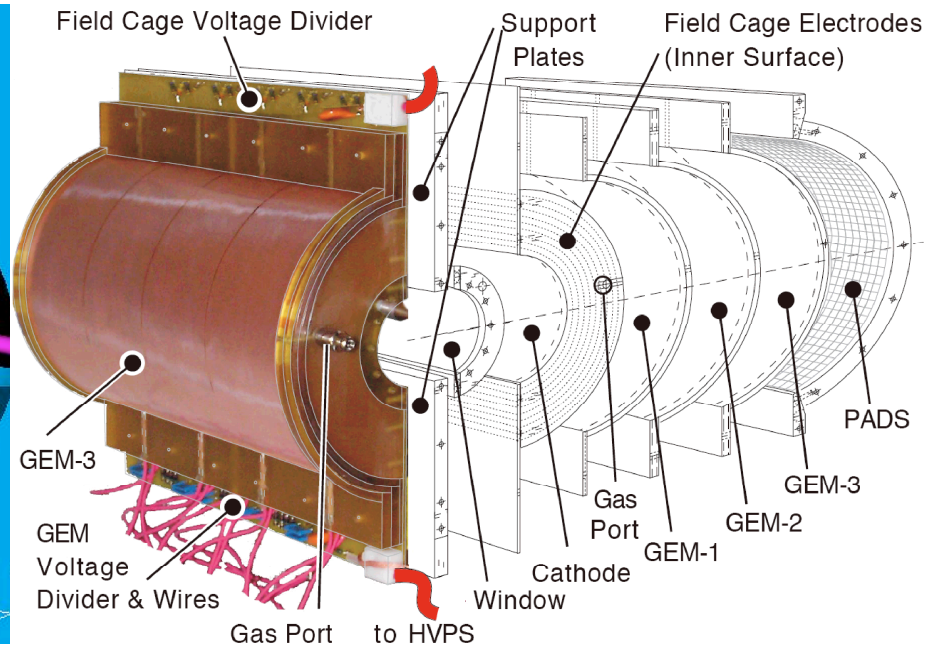
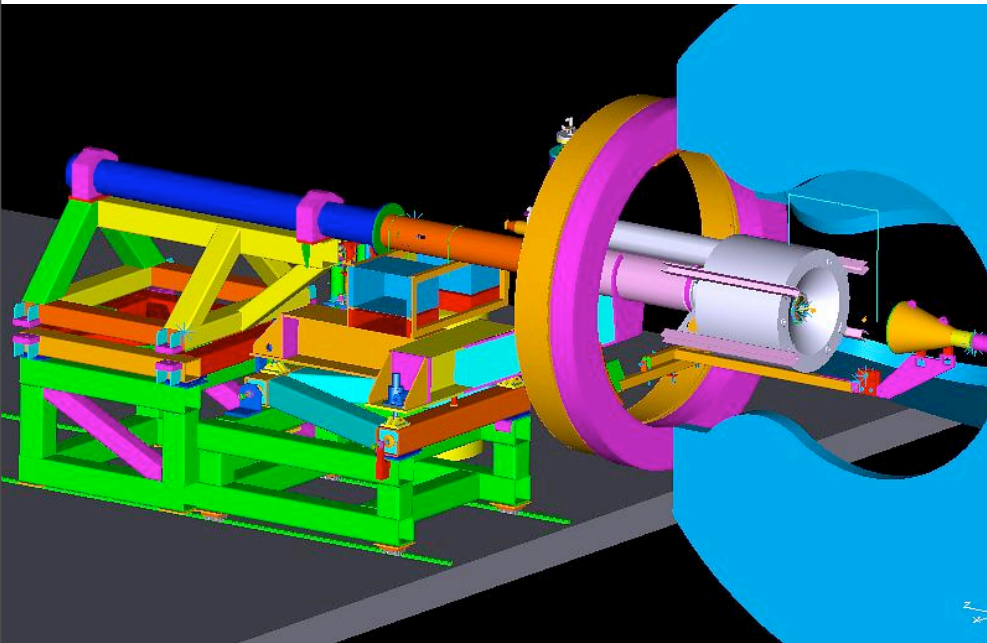
CEBAF Large Acceptance Spectrometer



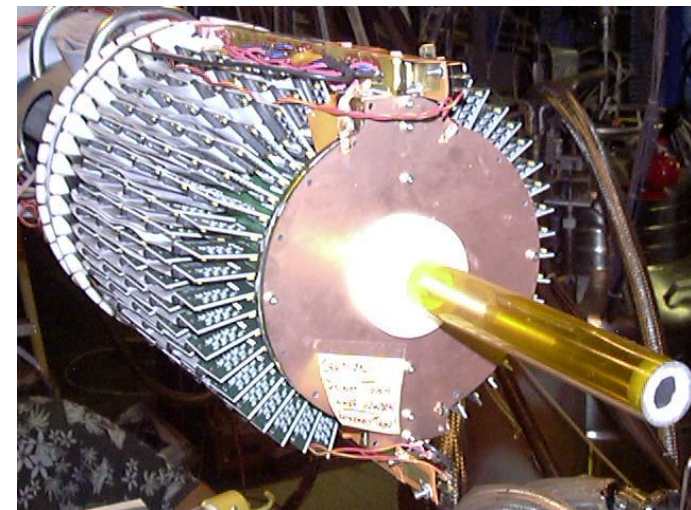
- DC: Drift Chamber
- CC: Cerenkov Counter
- SC: Scintillation Counter
- EC: Electromagnetic Calorimeter

100 cm
10



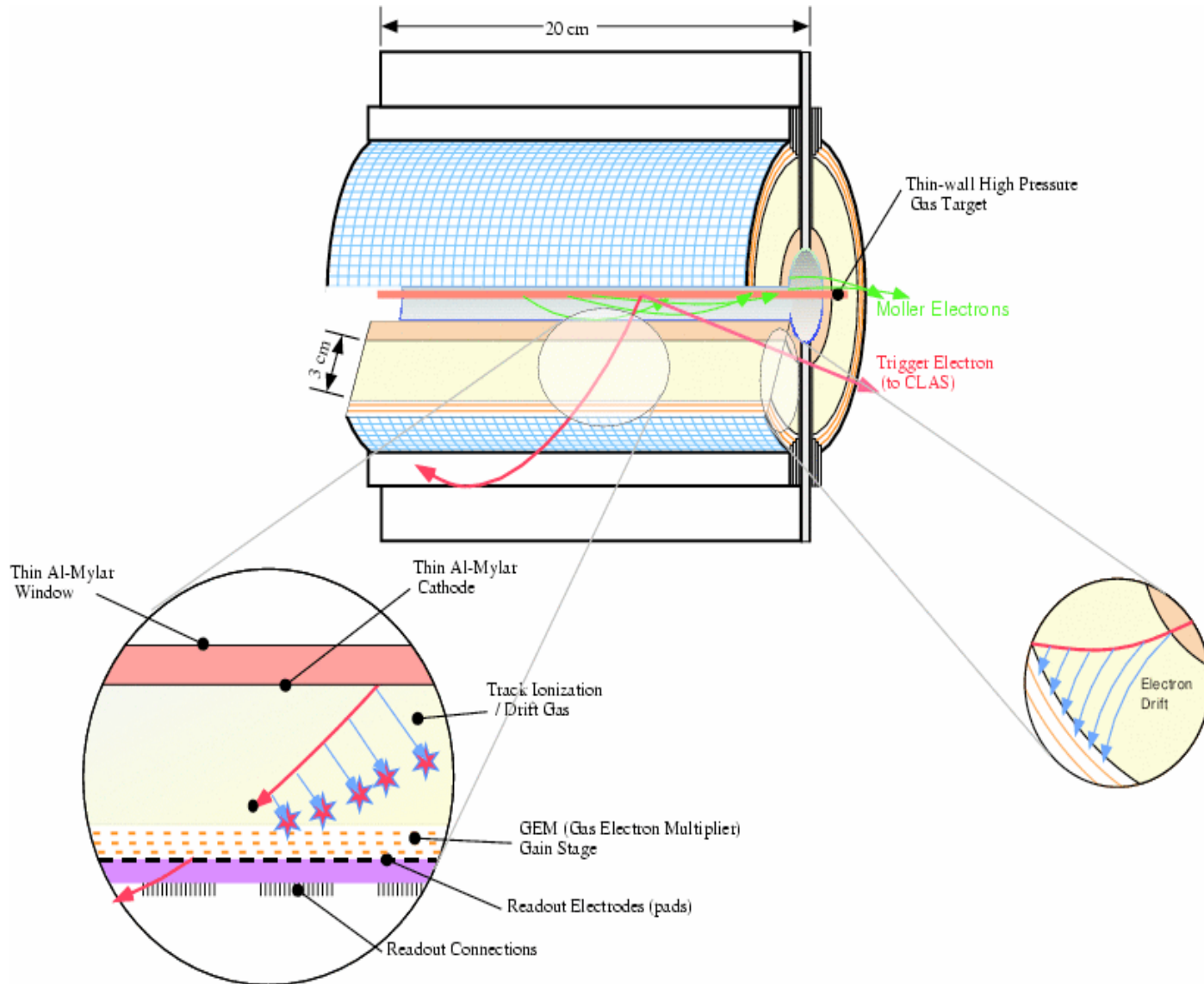


- Bound Nucleon Structure Experiment
- Hall B, JLab, CLAS
- $d(e, e'p_s)X$ with $0.07 < p_s < 0.15$ GeV/c
- $E_{\text{beam}} = 1.1, 2.1, 4.2, 5.3$ GeV
- Radial time projection chamber for p_s
- Data taking in 2005





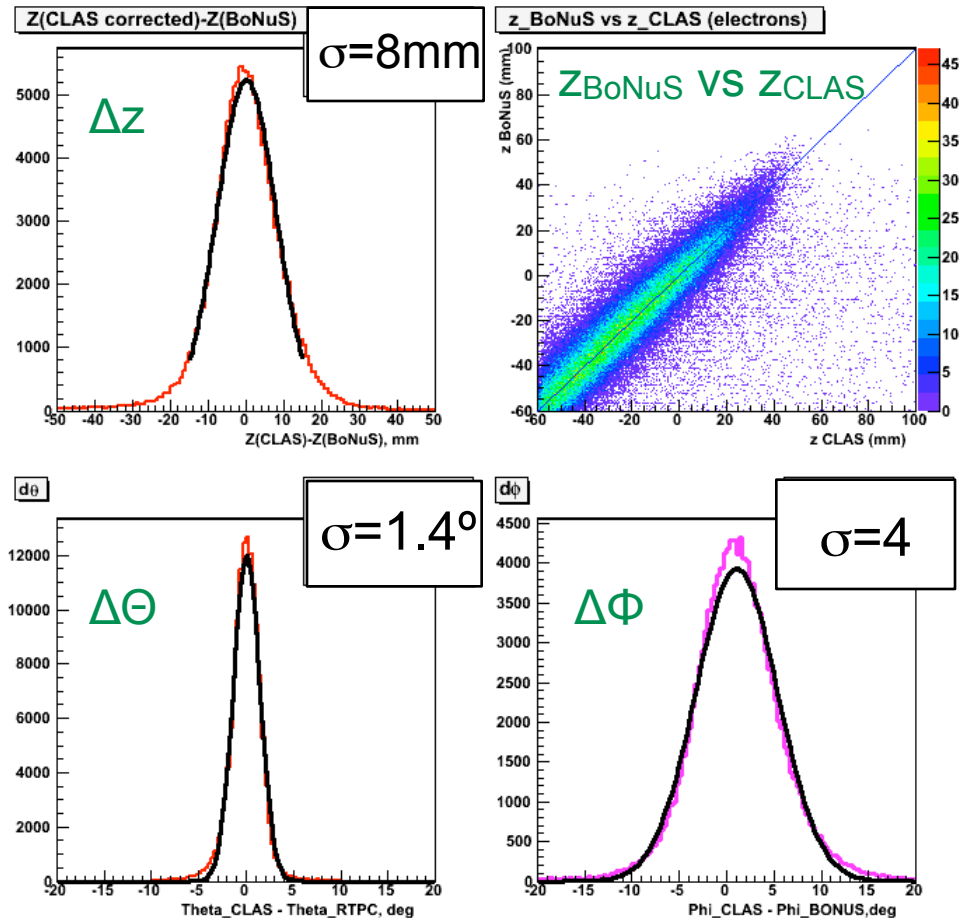
BoNuS Detector





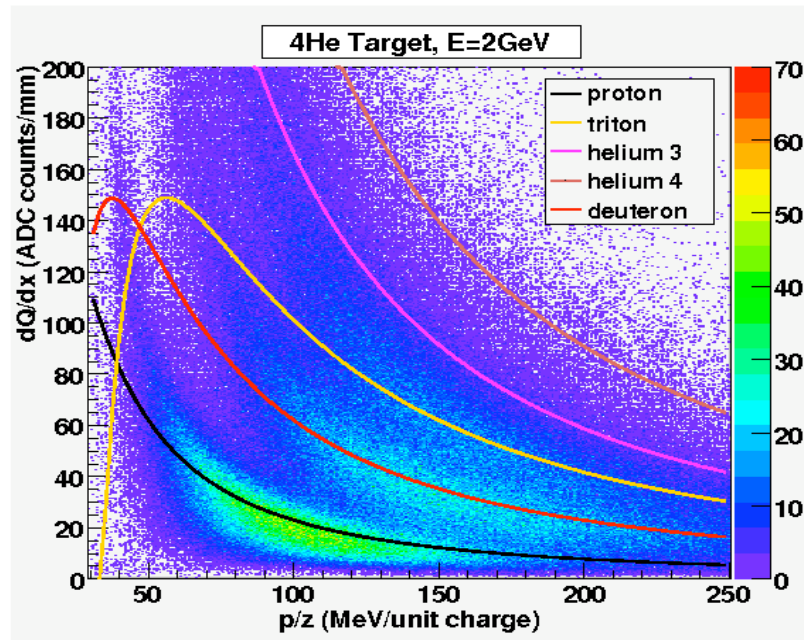
BoNuS RTPC Performance

- upper left: dE/dx vs. p/Z for He target
- lower left: dE/dx vs. p for deuterium target
- below RTPC+CLAS resolution for common e^- events

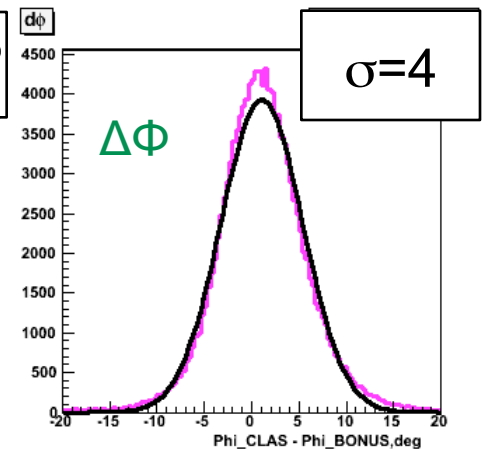
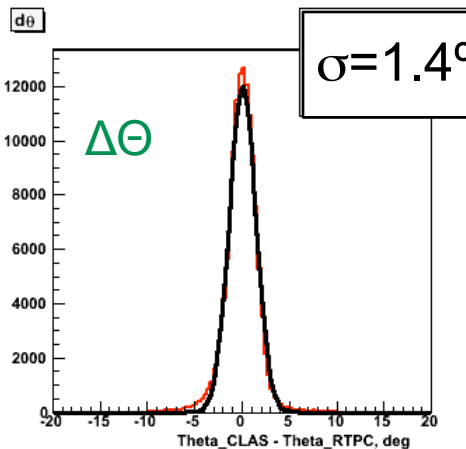
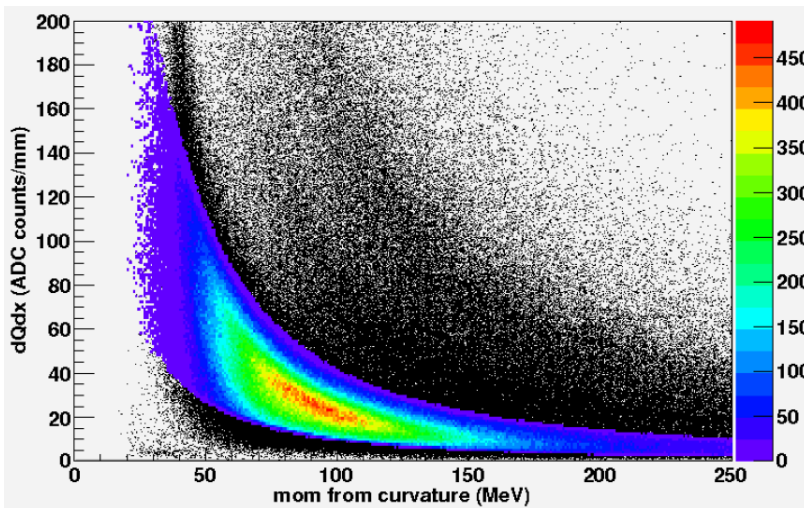
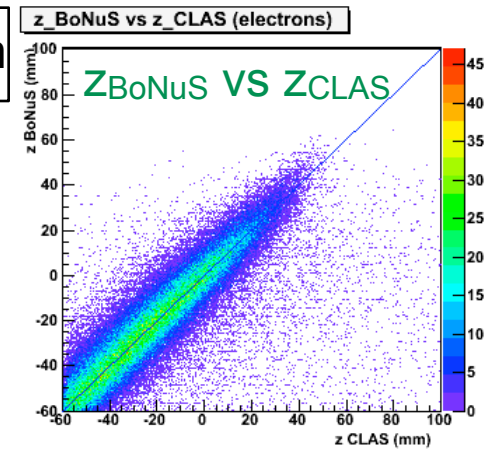
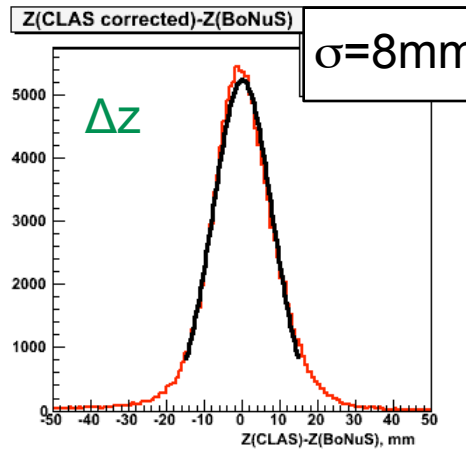




BoNuS RTPC Performance



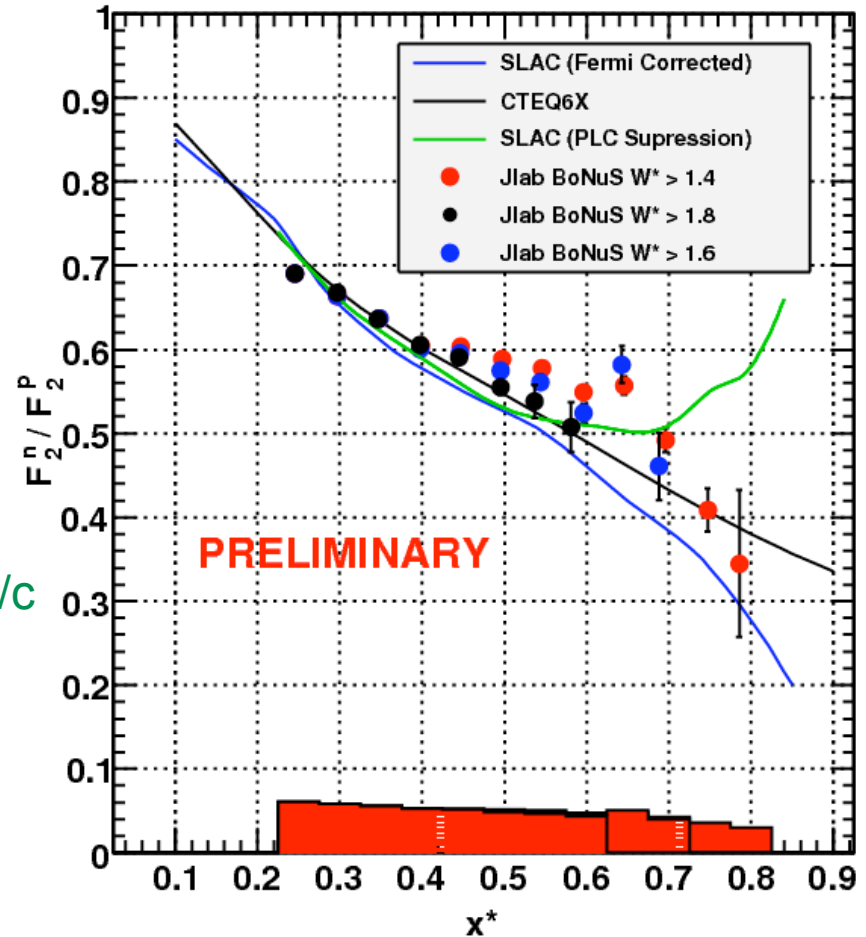
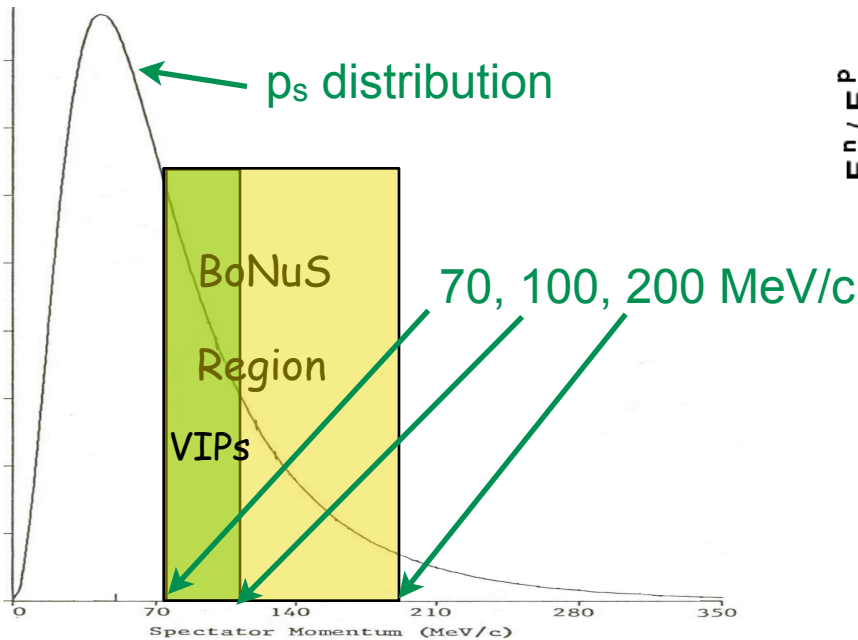
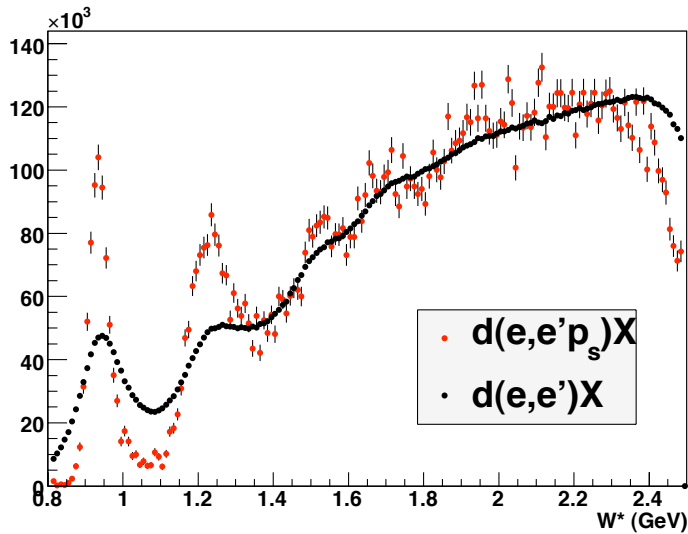
- upper left: dE/dx vs. p/Z for He target
- lower left: dE/dx vs. p for deuterium target
- below RTPC+CLAS resolution for common e^- events

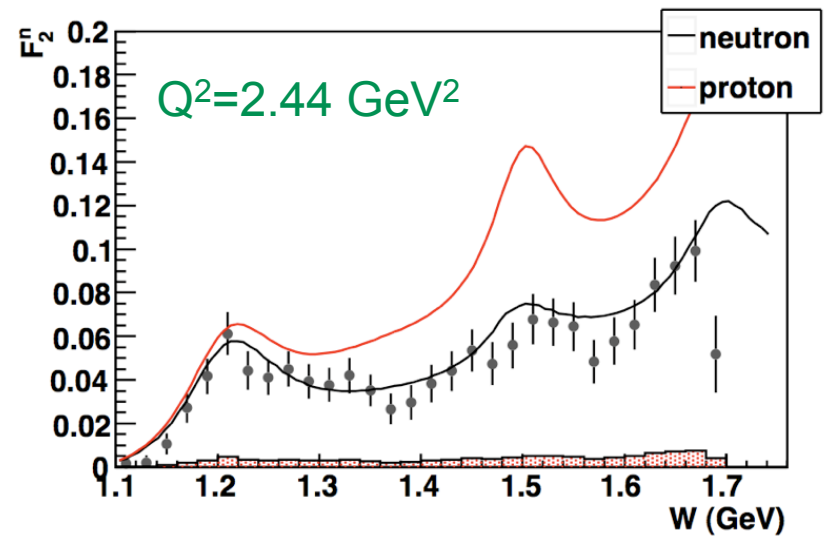
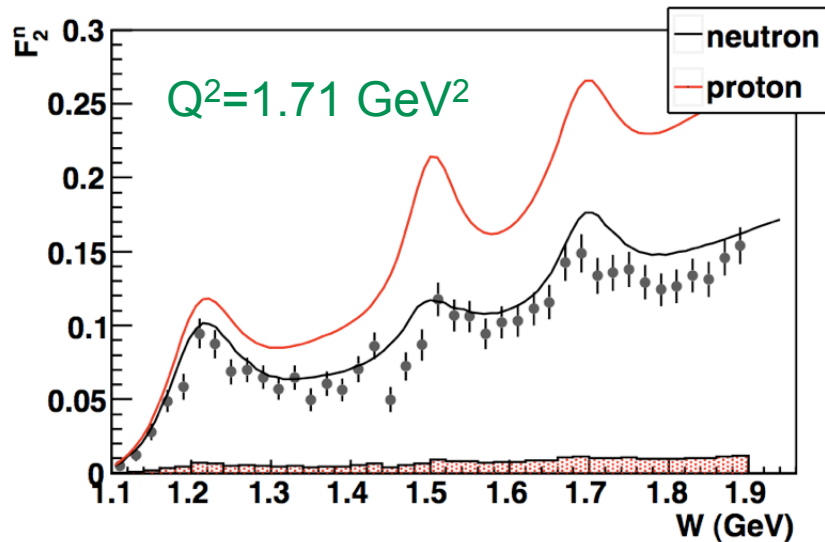
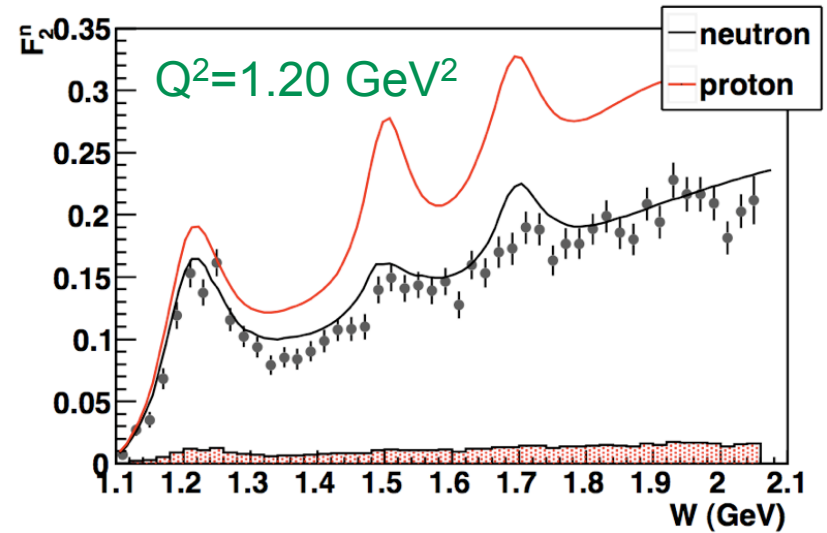
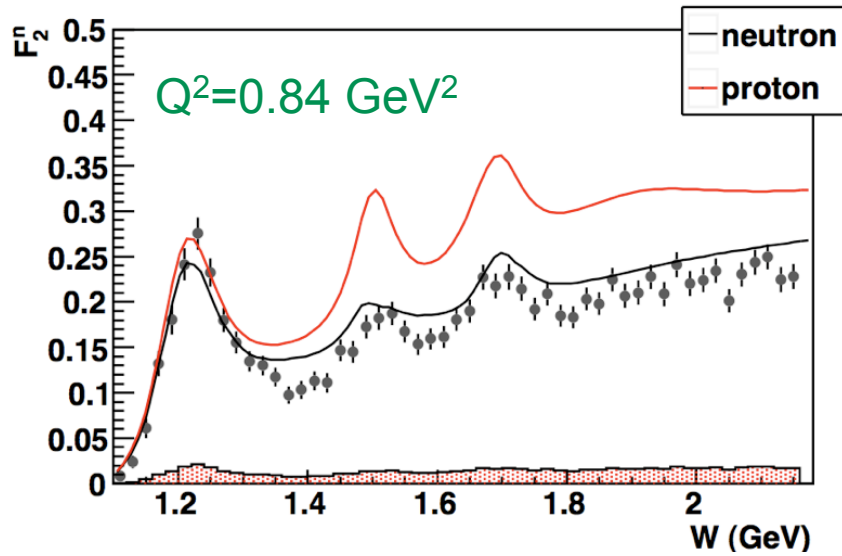




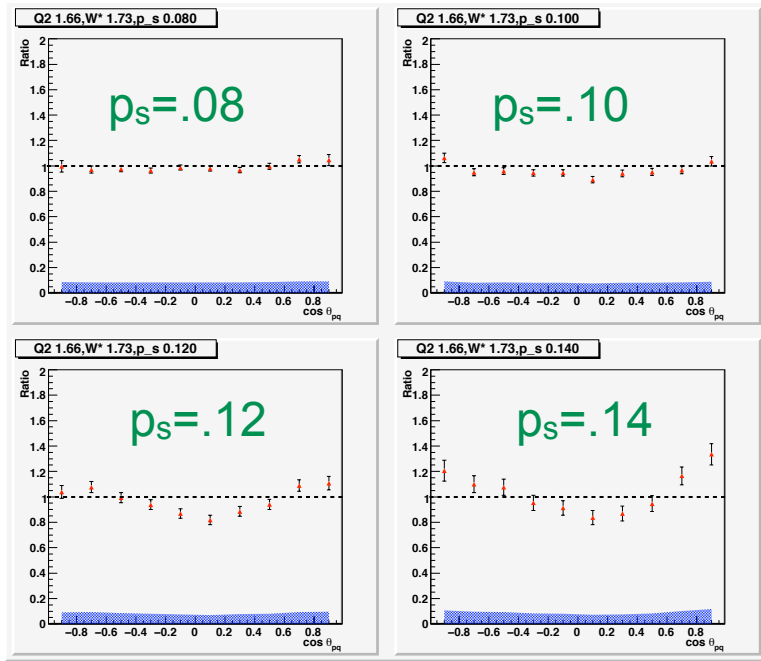
BoNuS Performance

- Very Important Protons $70 < p_s < 100$ MeV/c
- Corrections make resonances stand out
- F_2^n/F_2^p is measured at high x^*

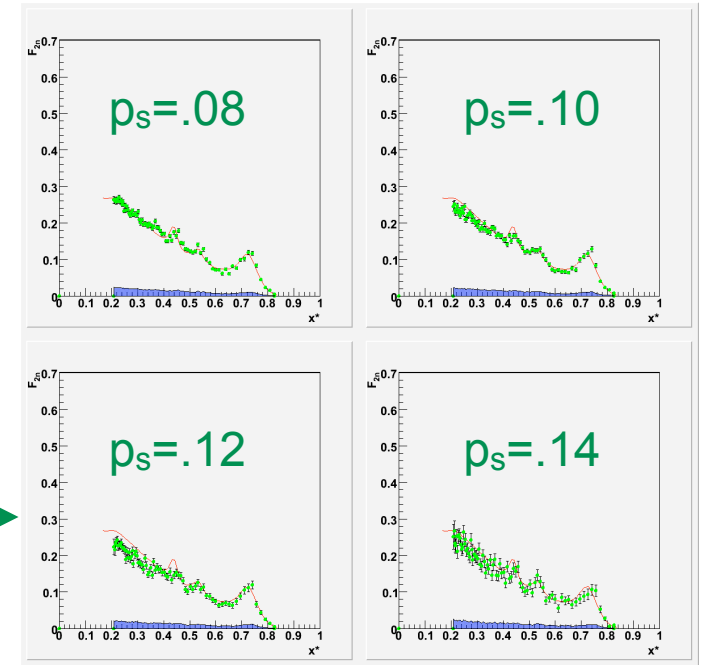




4 of 16 spectra: $0.8 < Q^2 < 4.5$; $E_{\text{beam}} = 4.2 \text{ \& } 5.3 \text{ GeV}$; Bosted/Christy world fits



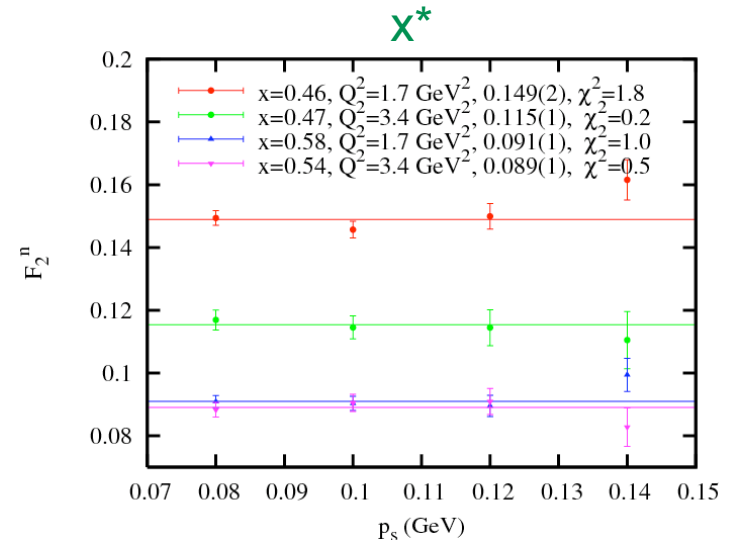
Data / PWIA
←



F_2^n
→

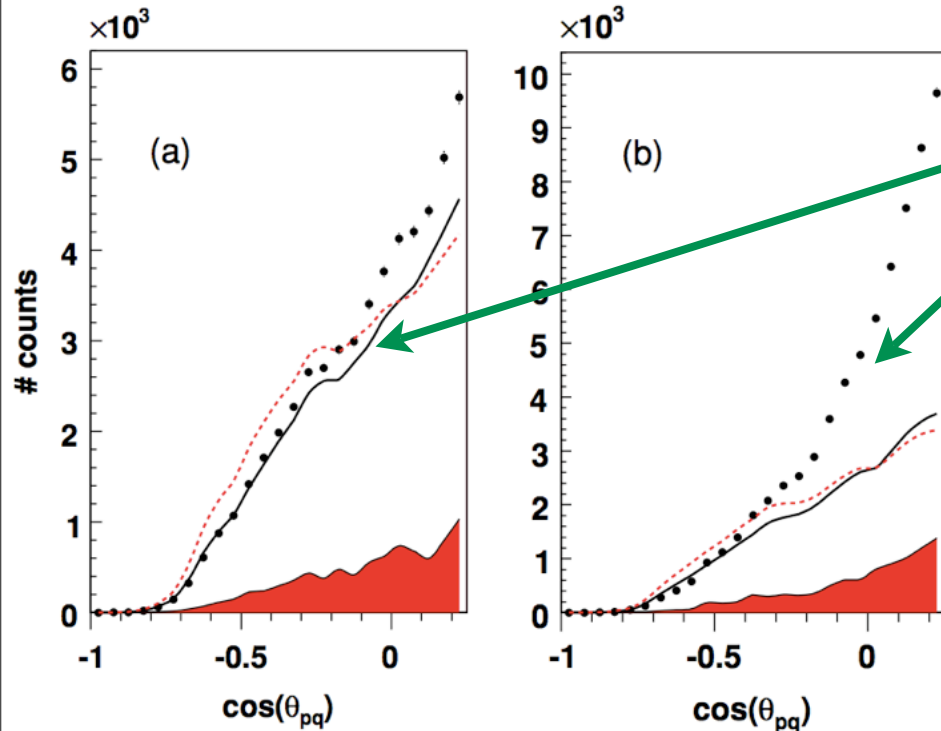
$\cos\theta_{pq}$

- FSIs start to appear for $p_s > 100$ MeV/c
- F_2^n vs. x shows no p_s -dependence
- There is no observed off-shell modification of F_2^n for $p_s < 0.14$ GeV/c



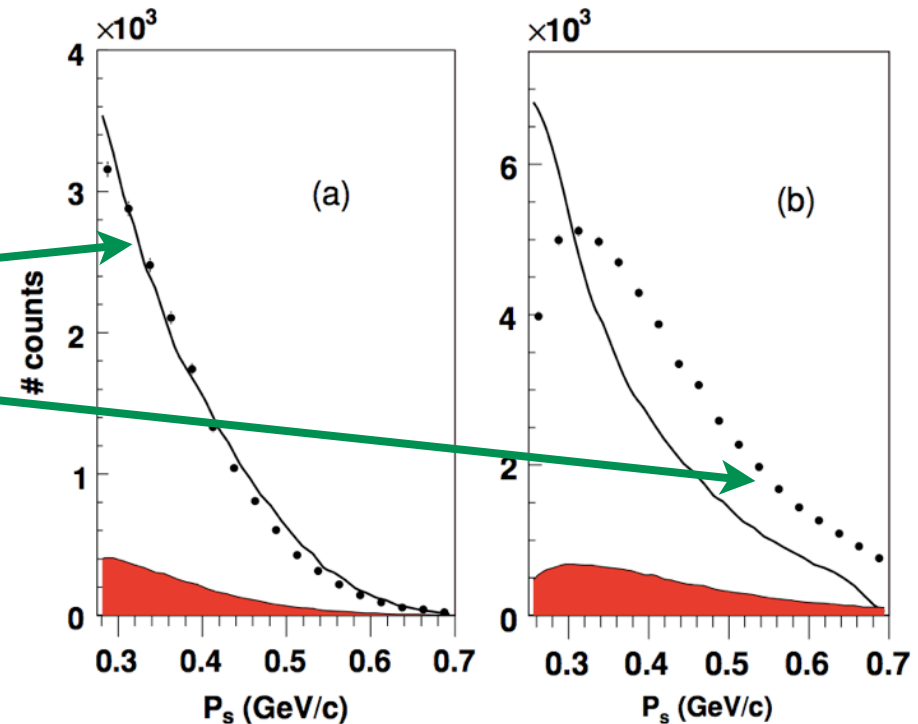


Klimenko *et al.*, PRC73(06)035212



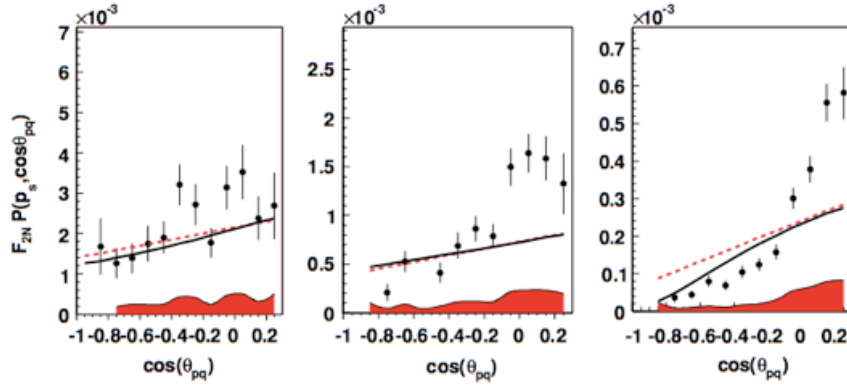
- (a) $p_s = 0.28-0.32$ GeV/c
- (b) $p_s = 0.36-0.42$ GeV/c
- (black line) PWIA w/ light cone w.f.
- (red dashed line) nonrelativistic w.f.
- Falloff is due to CLAS acceptance

- (a) $-1 < \cos \Theta_{pq} < -0.3$
- (b) $-0.3 < \cos \Theta_{pq} < 0.3$
- (black line) PWIA w/ light cone w.f.
- (red dashed line) nonrelativistic w.f.
- Θ_{pq} is the polar angle w.r.t. the q -vector

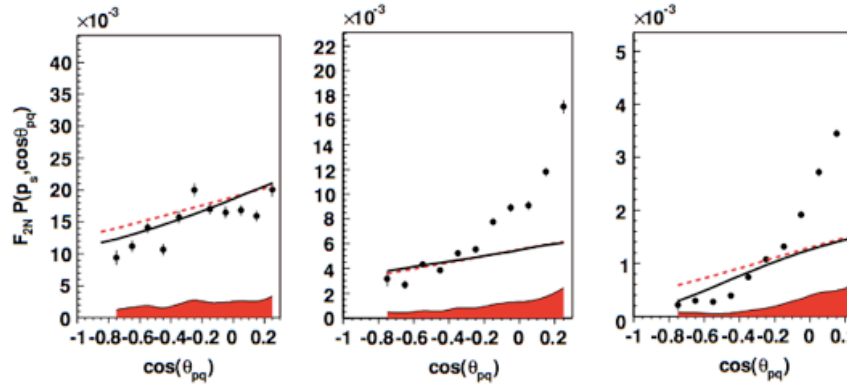




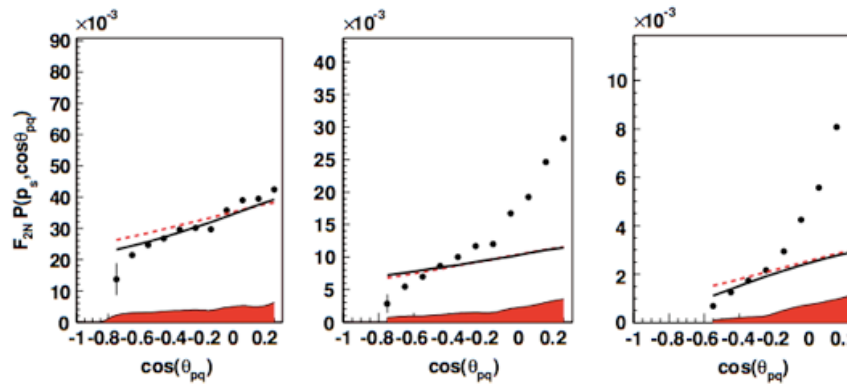
$W^*=0.94$



$W^*=1.50$



$W^*=2.00$



$p_s = .30$

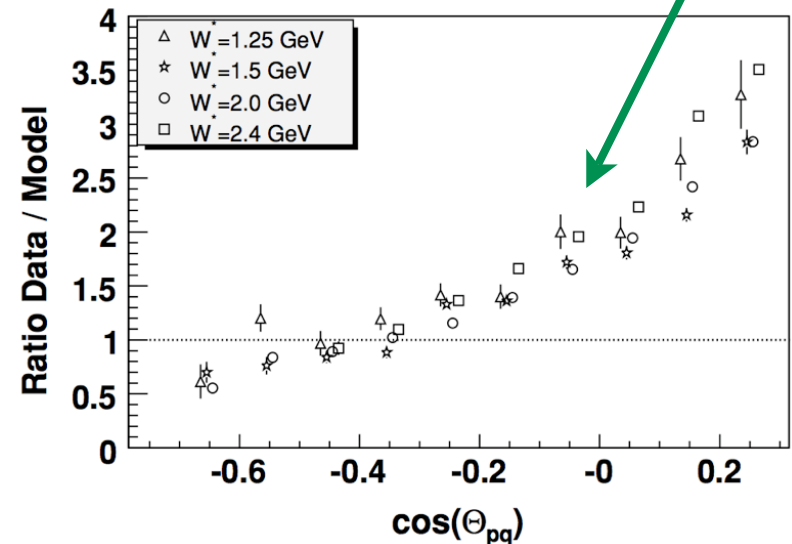
$p_s = .39$

$p_s = .56$

Klimenko *et al.*, PRC73(06)035212

- (red dashed line) nonrelativistic w.f.
- (black line) light cone
- Θ_{pq} is the polar angle w.r.t. \mathbf{q}

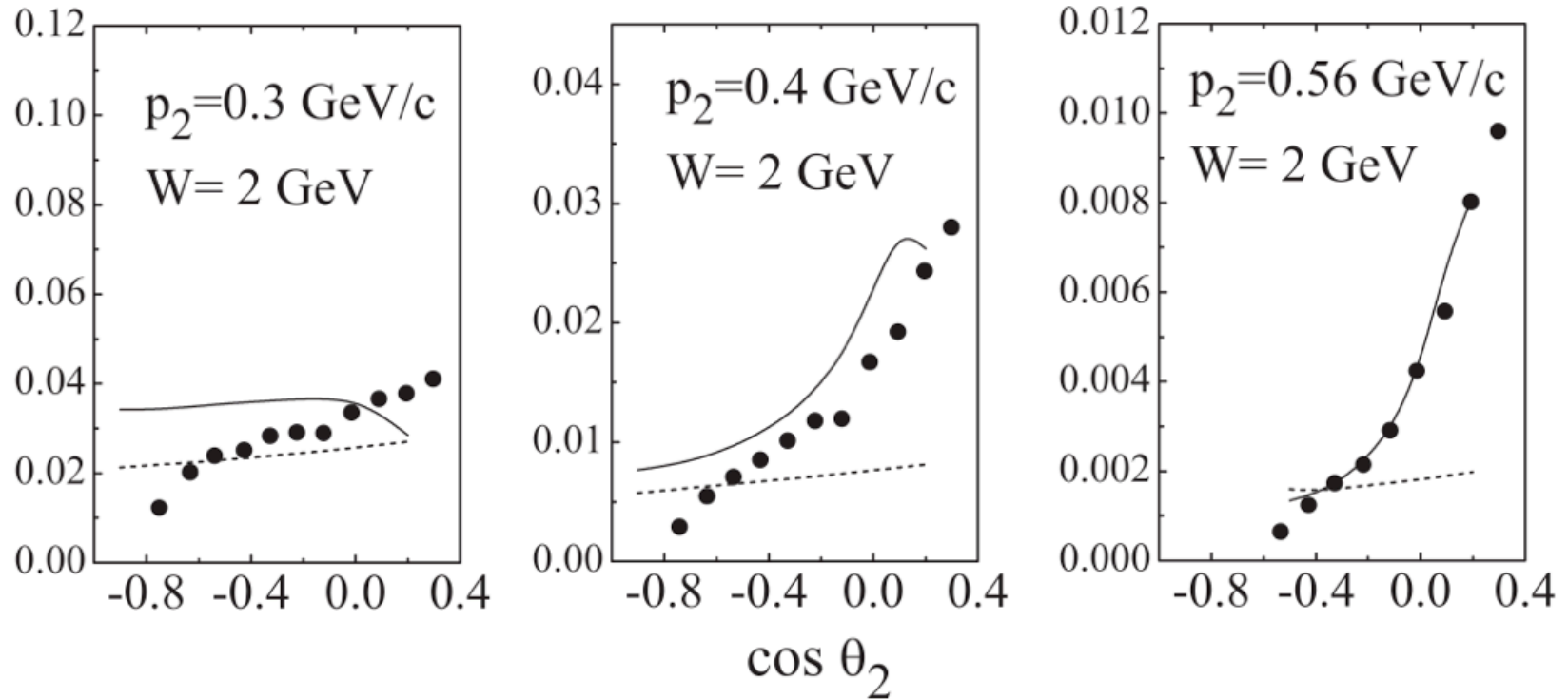
- Ratio of data to PWIA w/ light cone wf
- Ratio increases slightly with W^*
- FSI at very backward angles?





Palli et al, PRC80(09)054610

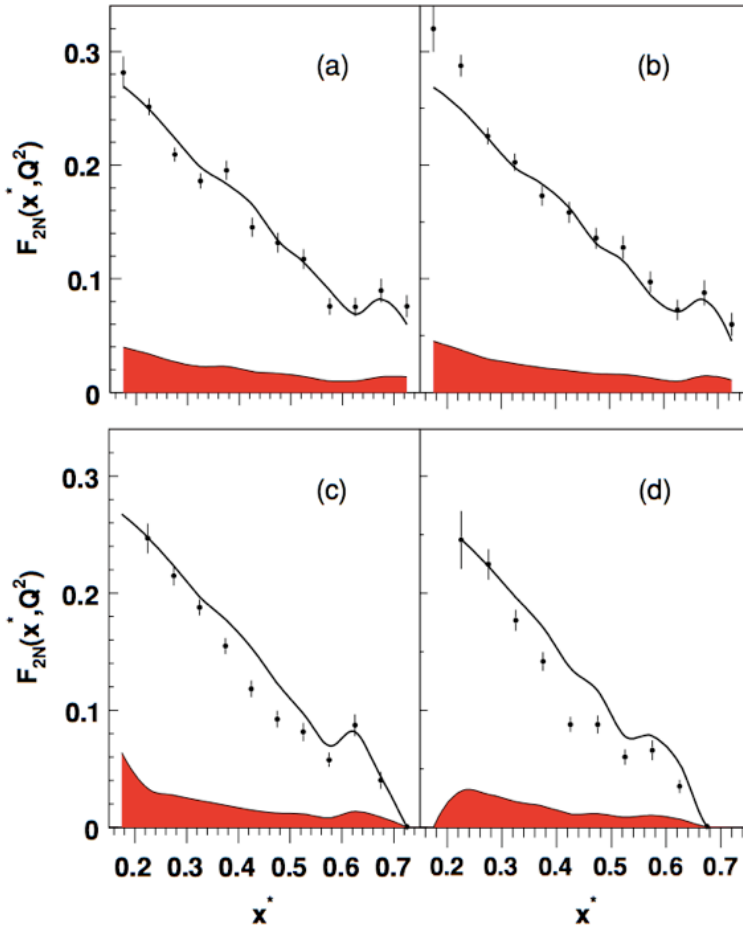
Reduced Cross Section $[(\text{GeV}/c)^{-3}]$



- A FSI model that works well for $p_s = 0.56 \text{ GeV}/c$ misses the mark at $p_s = 0.4$ and $0.3 \text{ GeV}/c$
- We see no discrepancy at $p_s = 0.56$ from an off-shell F_2 for $\cos \Theta_{pq} < -0.4$
- Dotted lines show the PWIA expectation (i.e. no FSI, no modified F_2)

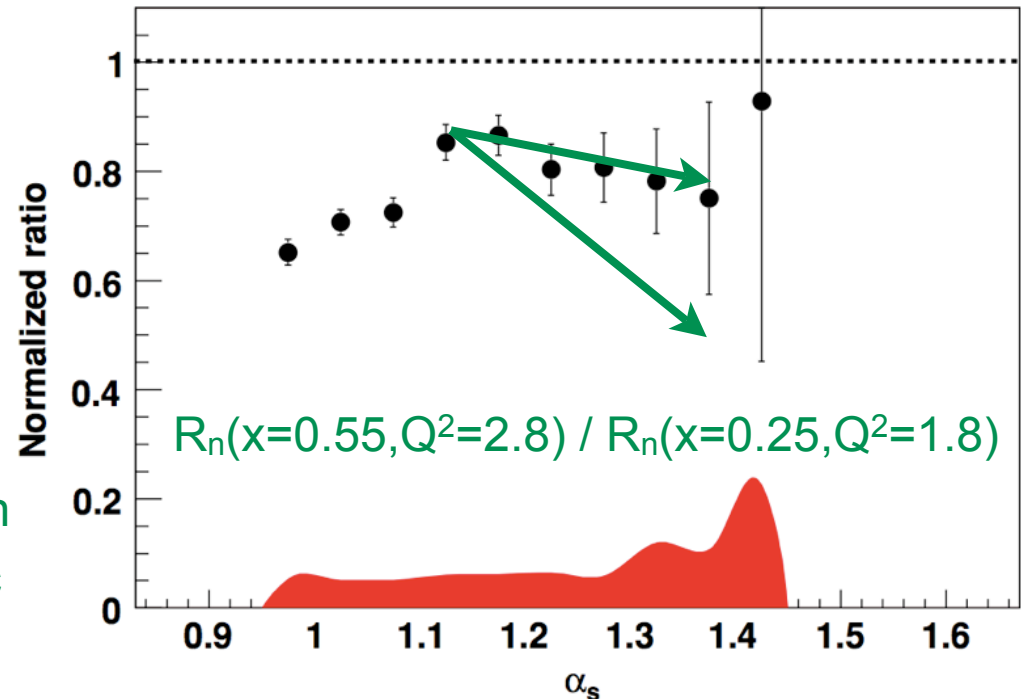


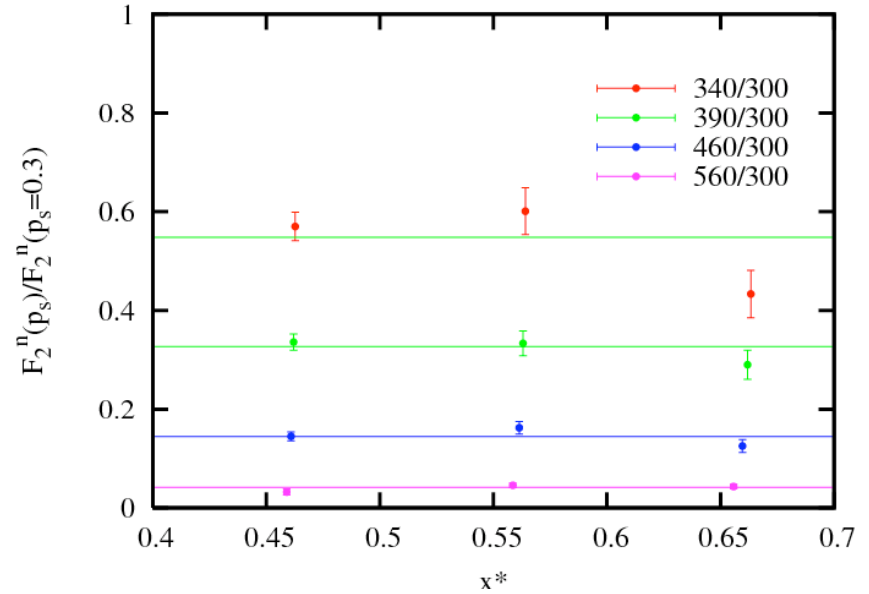
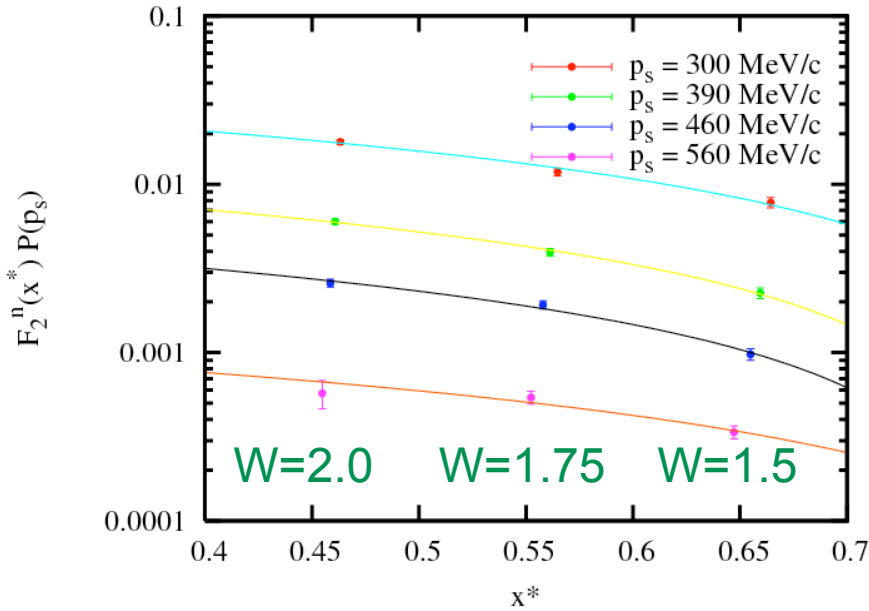
Klimenko *et al.*, PRC73(06)035212



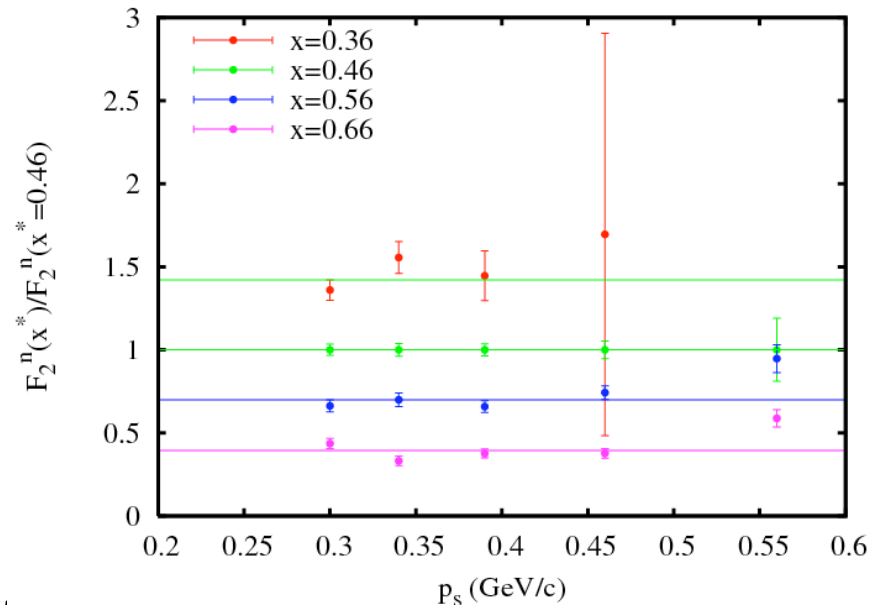
- (black line) PWIA with light cone w.f.
- x^* is the corrected momentum fraction
- (a-d): $p_s = 0.30, 0.34, 0.46, 0.56$ GeV/c
- at most a hint of p_s -dependence

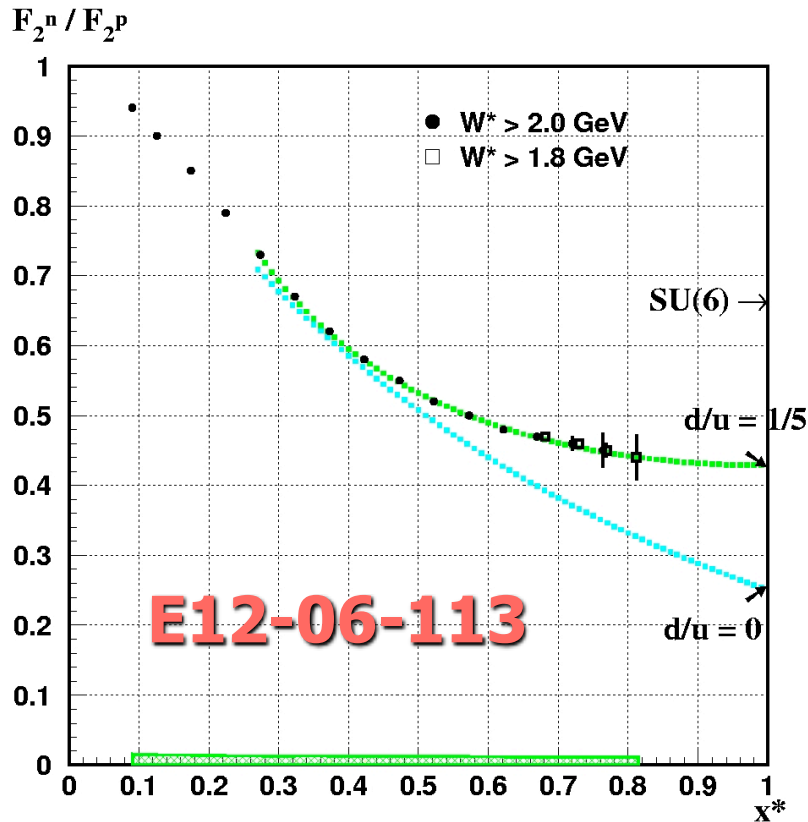
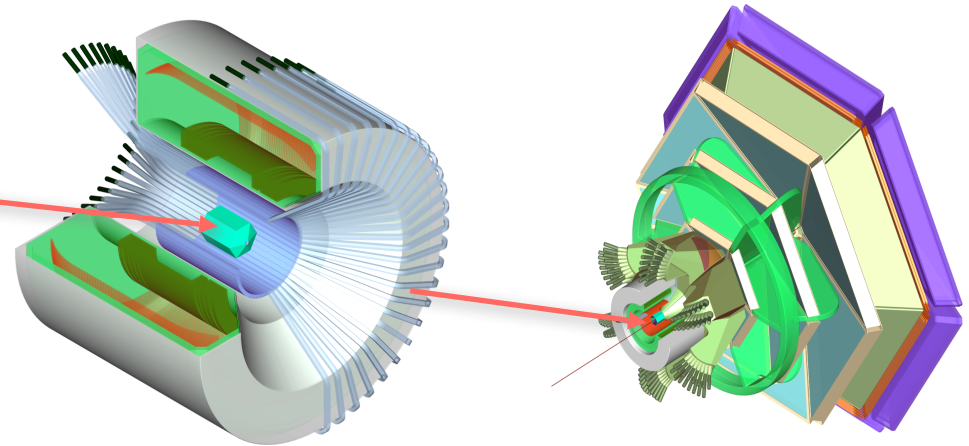
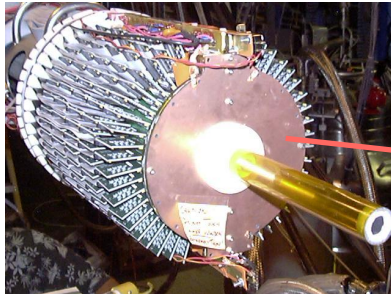
- Ratio of large to small EMC effect
- p_s dependence cancels in PWIA
- PWIA + off-shell predicts 1 at $\alpha_s = 1$
- FSI can upset norm
- No clear off-shell α_s -dependence





- pick $-0.45 < \cos < -0.3$ where integrated data match PWIA
- above: $F_2(x^*) P(p_s)$
- top right: $F_2(x^*) P(p_s) / F_2(x^*) P(300)$
- bottom right: $F_2(x^*) / F_2(0.46)$
- no sign of off-shell F_2 vs. x^* (top right) or vs. p_s (bottom right)





Data taking:

- 35 days on D_2
- 5 days on H_2
- $\mathcal{L} = 2 \times 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$

DIS region:

- $Q^2 > 1 \text{ GeV}^2$
- $W^* > 2 \text{ GeV}$
- $p_s < 100 \text{ MeV}/c$
- $\theta_{pq} > 110^\circ$
- $x^*_{\text{max}} = 0.80$

$W^* > 1.8 \text{ GeV}$: $x^*_{\text{max}} = 0.83$



- **BoNuS:**
 - we have measured F_2^n on a “free” neutron target
 - no effects from Fermi motion and the EMC
 - no evidence for off-shell structure for $p_s < 140$ MeV/c
- **DEEPS:**
 - understanding FSIs is crucial to extract an off-shell F_2^n
 - data tend to favor models with a small off-shell F_2^n
- **CLAS12:**
 - new BoNuS proposal is conditionally approved
 - new DEEPS proposal is still a twinkle in our eyes