



Deep Inelastic and Deep Exclusive Results from Jefferson Lab

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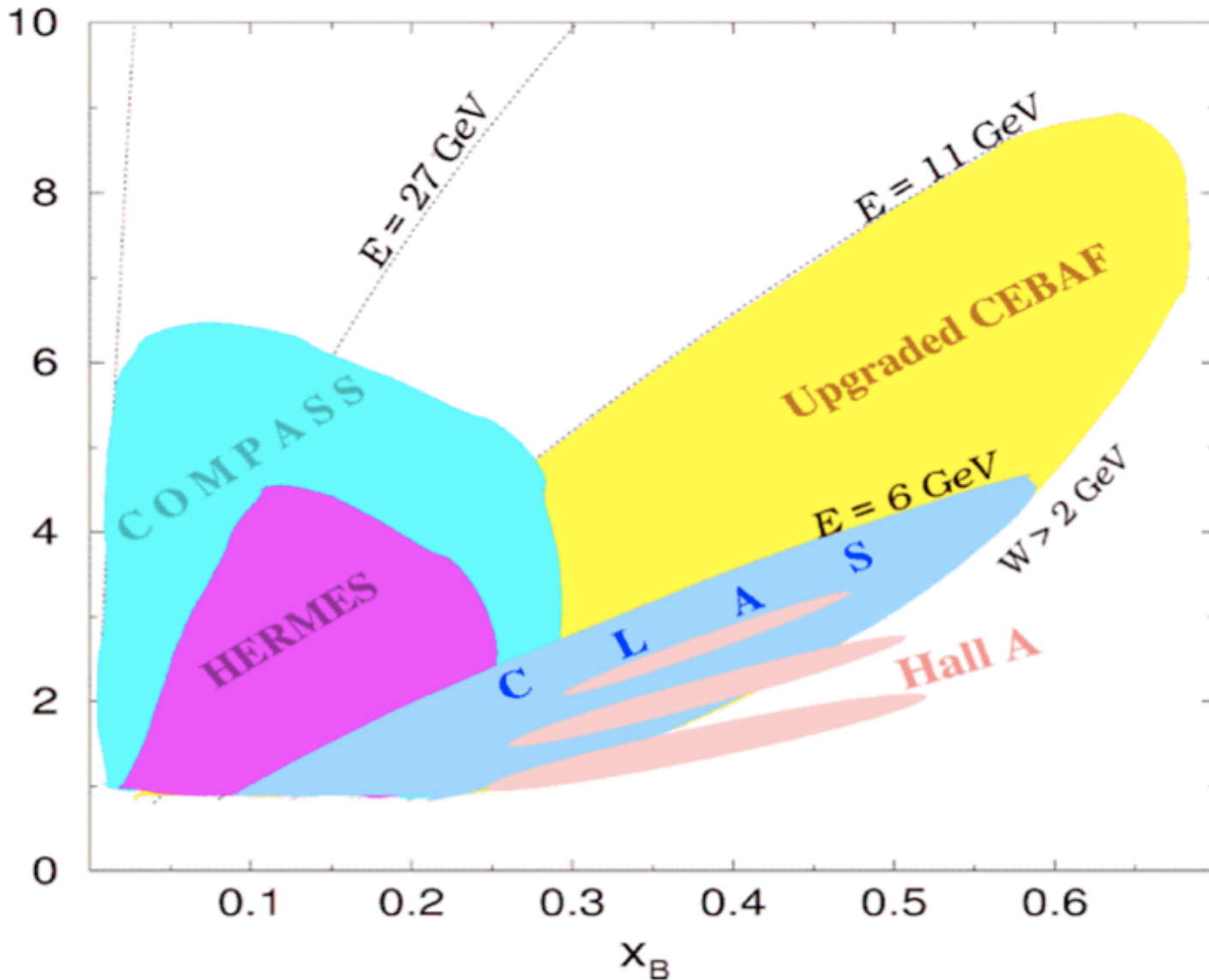
Deep Inelastic Scattering 2011
Newport News, Virginia
11 April 2011



- The only thing we can measure is a cross section.
- But by separating kinematics from nucleon structure, we can identify robust, experimentally determined objects, the structure functions:

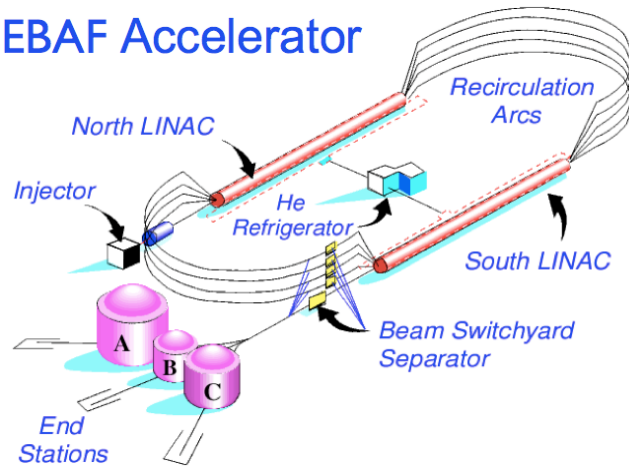
$$\frac{d\sigma}{dx dy d\psi} = \frac{2\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left\{ \boxed{F_T} + \varepsilon \boxed{F_L} + S_{\parallel} \lambda_e \sqrt{1-\varepsilon^2} 2x \boxed{(g_1 - \gamma^2 g_2)} \right. \\ \left. - |S_{\perp}| \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S 2x\gamma \boxed{(g_1 + g_2)} \right\}$$

- Thus, F_T , F_L , g_1 , $g_2(x, Q^2)$ can be extracted for all x , Q^2 .
- Experiment tells us where these can be interpreted in terms of parton distribution functions in pQCD and where complications show up.
- PDFs are known only through model fitting of structure functions.
- More so for transverse momentum dependent distributions and generalized parton distributions

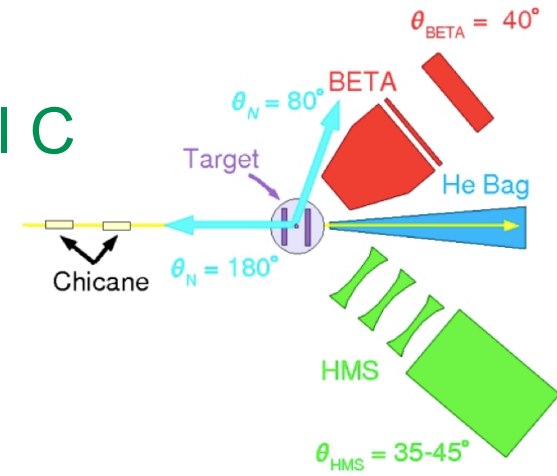




The CEBAF Accelerator

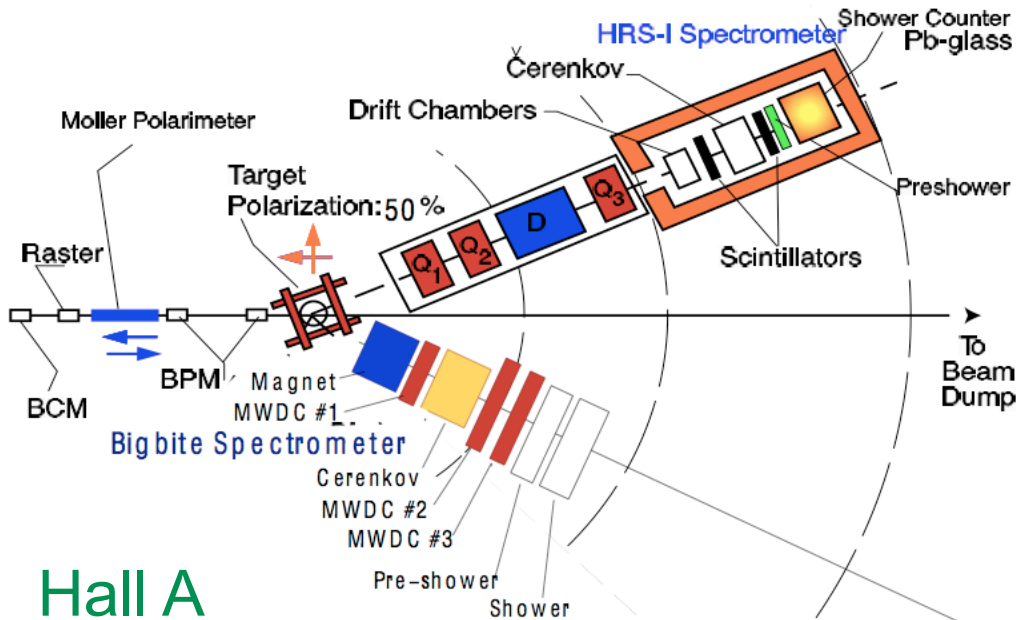
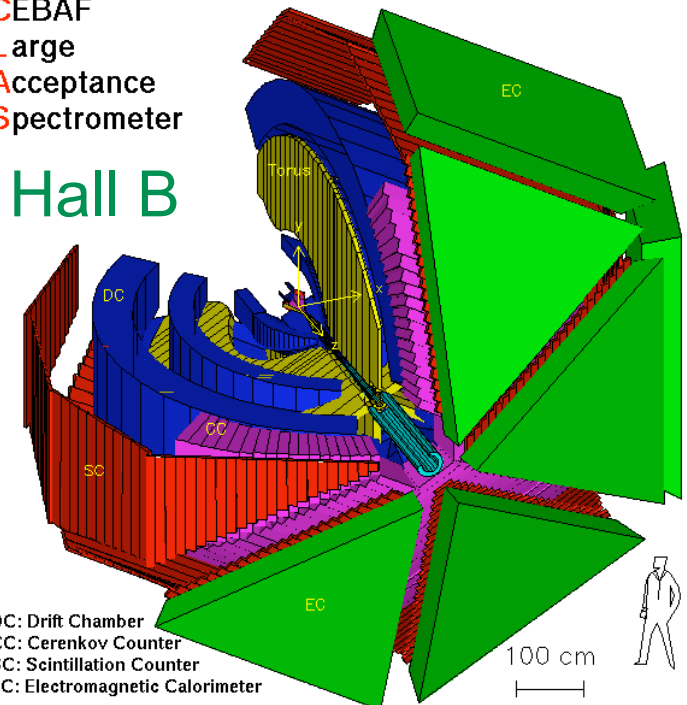


Hall C



CEBAF
Large
Acceptance
Spectrometer

Hall B



Hall A

To Beam Dump



- ★WG1PS2: Jeff Owens *Uncertainties in determining the d quark PDF at large values of x*
- ★WG1PS2: Slava Tkachenko *Model independent extraction of neutron structure functions from deuterium data*
- ★WG1PS7: Peter Monaghan *First Extraction of F_L Moments from World Data*
- ★WG1PS7: Ibrahim Albayrak ... *Deuteron F_L ... and Extractions of the Deuteron and Non-Singlet Moments*
- ★WG1PS9: Silvia Pisano *Results and Achievements at CLAS*
- ★WG1PS9: Simona Malace *Quark-hadron duality*
- ★WG1PS10: Patricia Solvignon *The nuclear dependence of $R=\sigma_L/\sigma_T$ in Deep Inelastic Scattering*
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- ★WG4SINS: Hayk Hakobyan *Quark propagation and hadron formation in the nucleus*
- ★WG4SINS: Sergio Anefalos Pereira *Strangeness production in CLAS*
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- ★WG6PSH2: Nilanga Liyanage *Moments of the neutron g_2 structure function and ... higher-twist effects*
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- ★WG6PSHP1: Yohann Perrin *Coherent deeply virtual Compton scattering off helium (CLAS)*
- ★WG6PSHP1: Andrey Kim *Studies of exclusive processes with a longitudinally polarized target*
- ★WG7PS3: Dave Gaskell, Xin Qian, Yelena Prok, Javier Gomez, Francois-Xavier Girod, Gordon D. Cates *12 GeV*

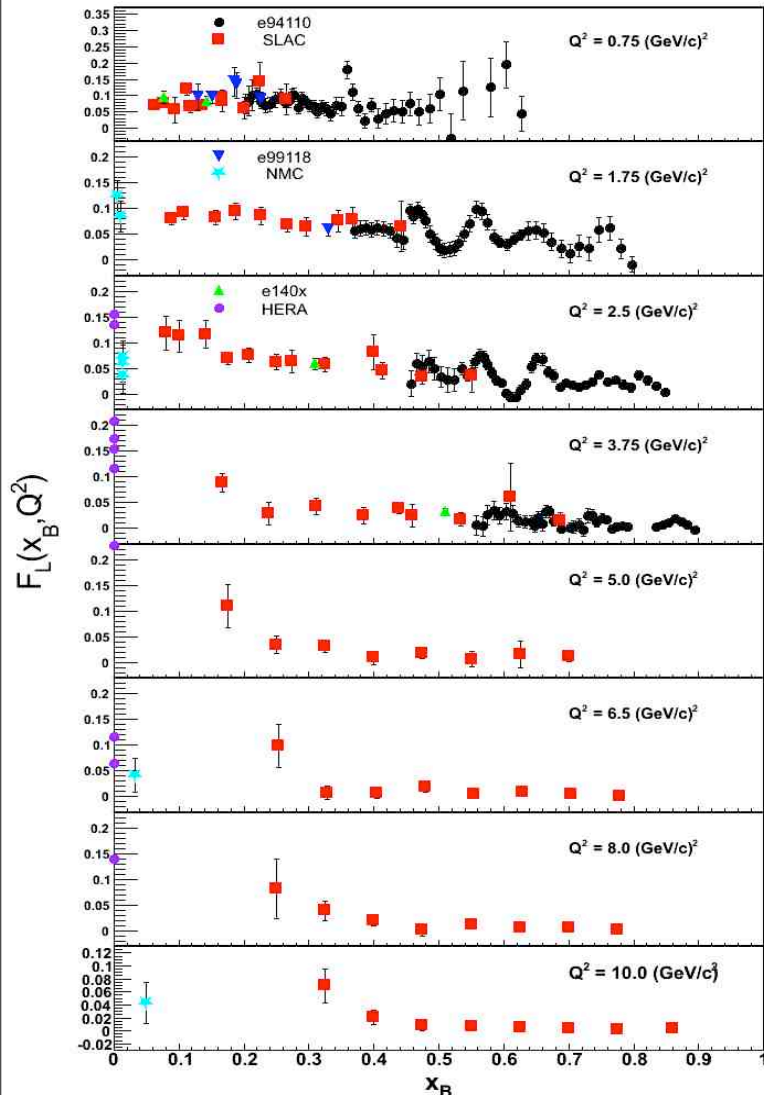


$$\frac{d\sigma}{dx dy d\psi} = \frac{2\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left\{ \boxed{F_T} + \varepsilon \boxed{F_L} + S_{\parallel} \lambda_e \sqrt{1-\varepsilon^2} 2x \boxed{(g_1 - \gamma^2 g_2)} \right. \\ \left. - |\mathbf{S}_{\perp}| \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S 2x\gamma \boxed{(g_1 + g_2)} \right\}$$



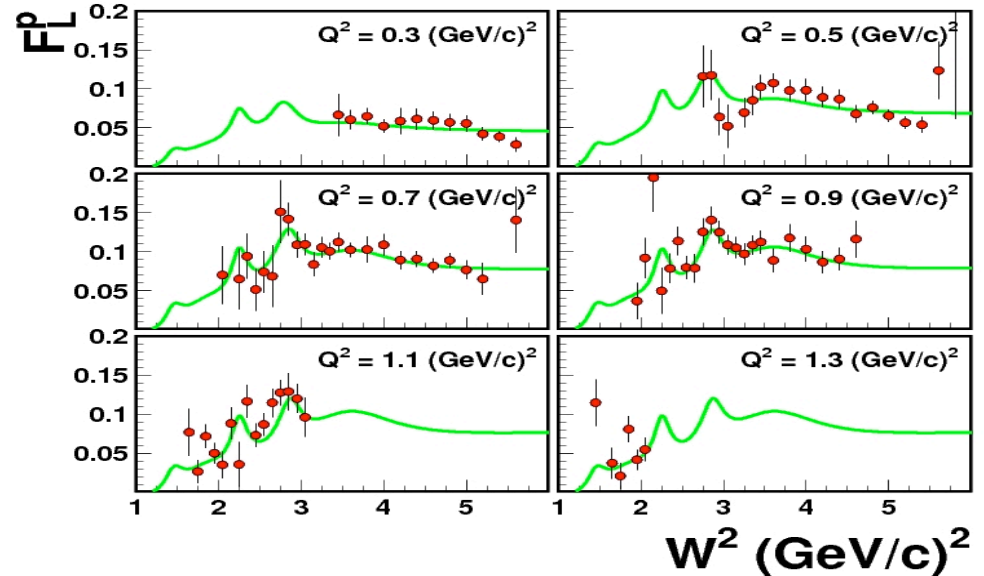
★WG1PS7: Peter Monaghan *First Extraction of F_L Moments from World Data*

★WG1PS7: Ibrahim Albayrak ... *Deuteron F_L ... and Extractions of the Deuteron and Non-Singlet Moments*



E04-110 proton

E00-002, Deuteron



- Rosenbluth separation of F_L and F_T
- Moments require data at all x , including the resonance region



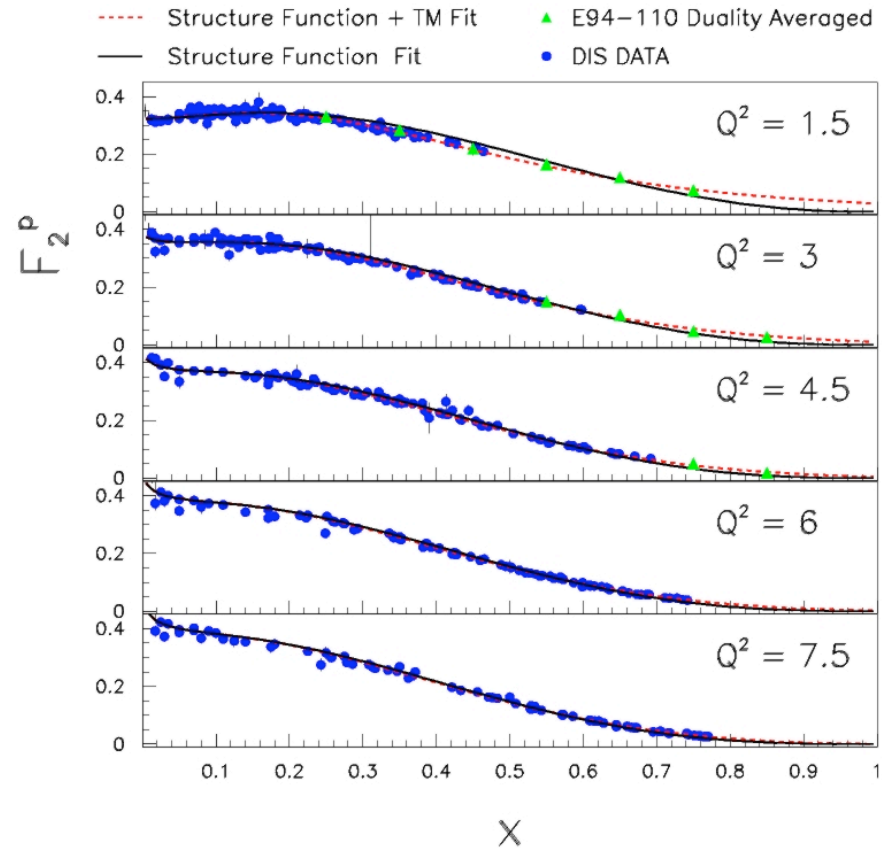
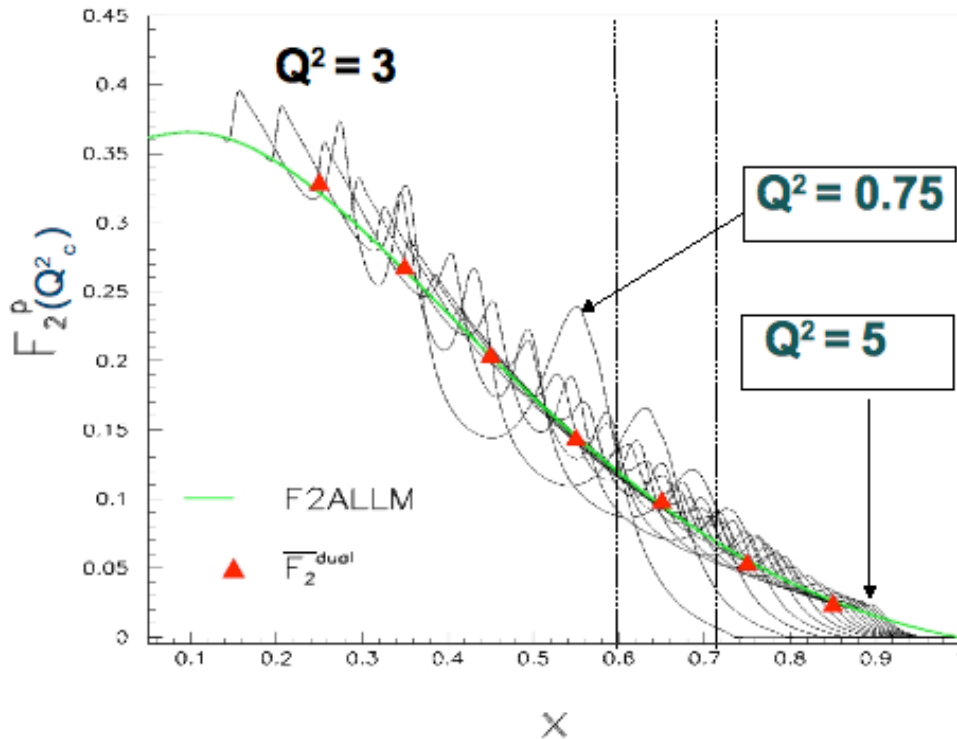
★WG1PS9: Simona Malace *Quark-hadron duality*

Hall C
Christy

$$F_2^{\text{TMC}}(x, Q^2) = \frac{x^2}{\xi^2 r^3} F_2^{(0)}(\xi) + \frac{6M^2 x^3}{Q^2 r^4} h_2(\xi) + \frac{12M^4 x^4}{Q^4 r^5} g_2(\xi)$$

$$h_2(\xi, Q^2) = \int_{\xi}^1 du \frac{F_2^{(0)}(u, Q^2)}{u^2} \quad \xi = \frac{2x}{1 + \sqrt{1 + 4x^2 M^2 / Q^2}}$$

$$g_2(\xi, Q^2) = \int_{\xi}^1 du h_2(u, Q^2)$$



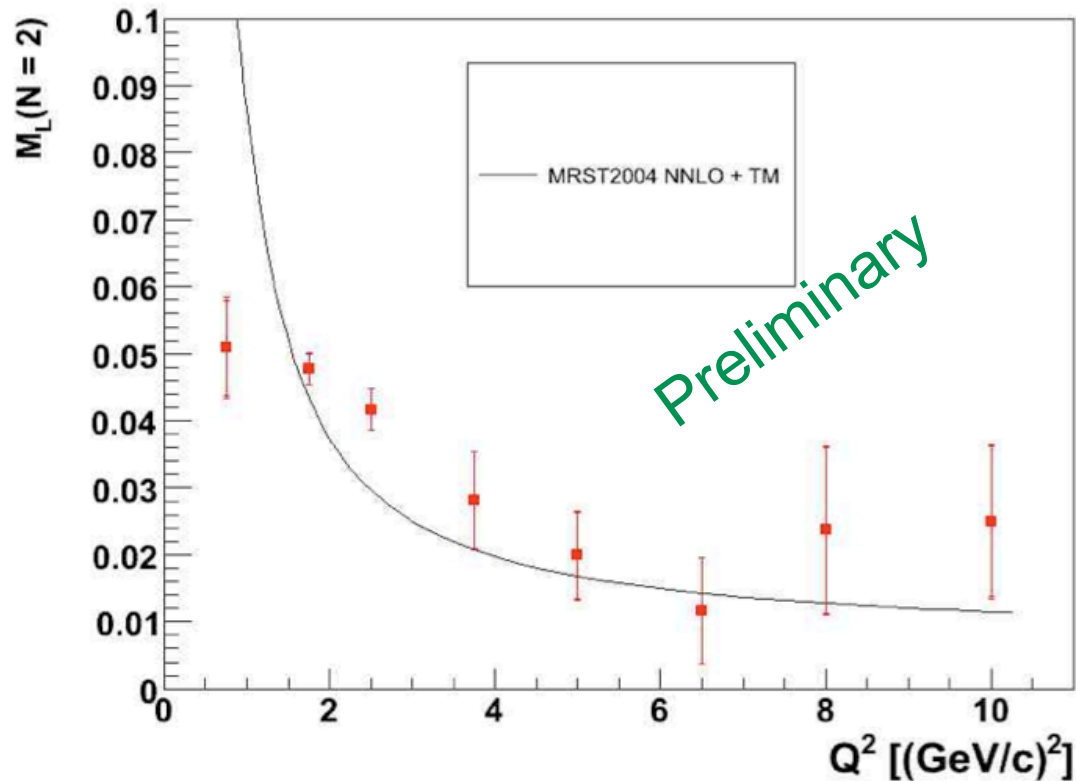


★WG1PS7: Peter Monaghan *First Extraction of F_L Moments from World Data*

Cornwall-Norton Moments

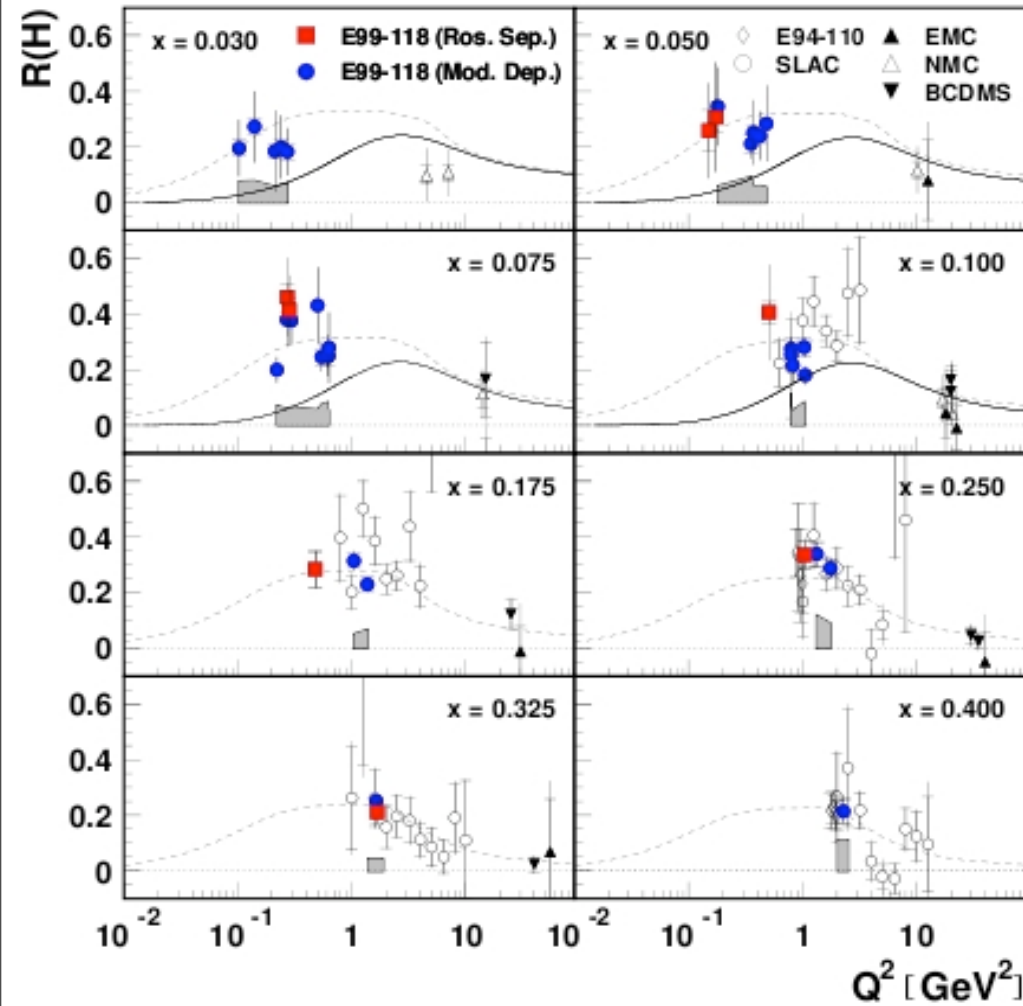
$$M_n^{2,L}(Q^2) \equiv \int_0^1 dx x^{n-2} F_{2,L}(x, Q^2)$$

Analysis of P. Monaghan, et al.





★WG1PS10: Patricia Solvignon *The nuclear dependence of $R=\sigma_L/\sigma_T$ in Deep Inelastic Scattering*

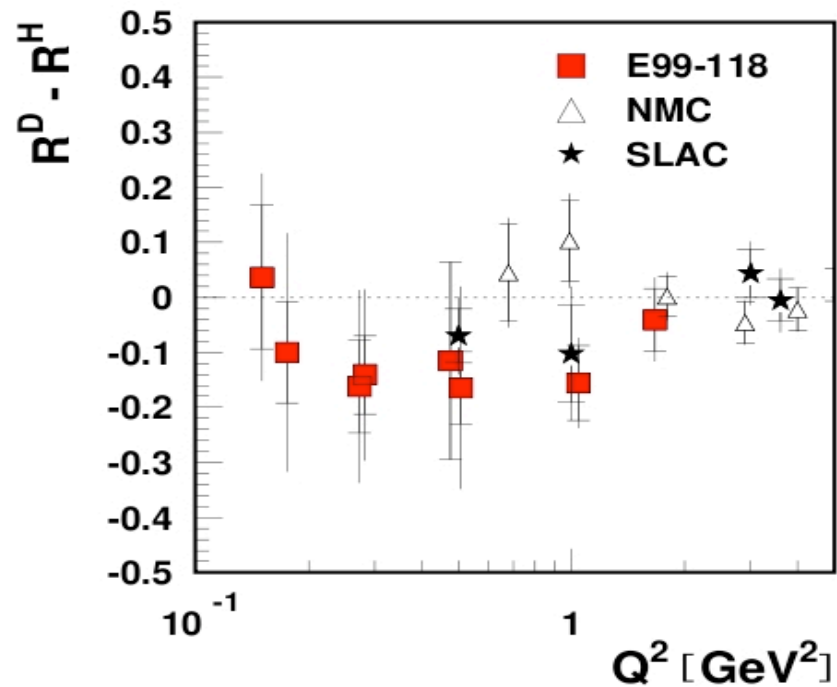


E99-118

p,d targets

$Q^2=0.1-1.7$ GeV²

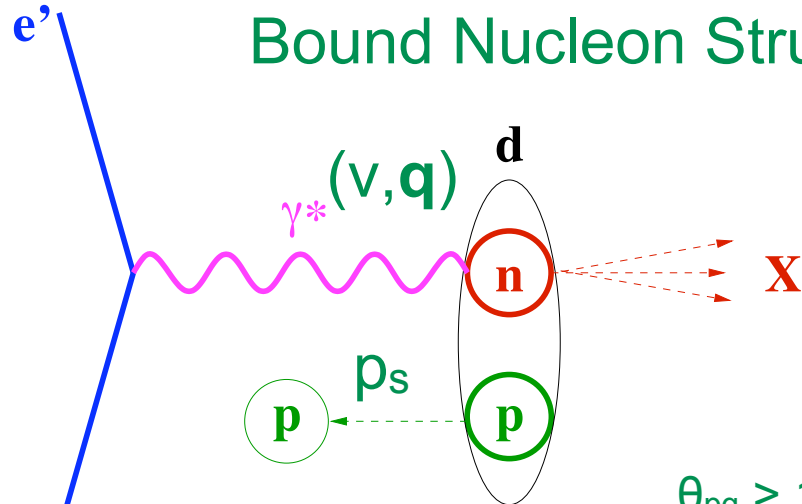
Measure of $R=F_L/F_T$





★WG1PS2: Slava Tkachenko *Model independent extraction of neutron structure functions from deuterium data*

Bound Nucleon Structure Experiment using CLAS



Detect the spectator proton from deuterium following en scattering. Make kinematic corrections using the spectator proton's energy E_s and momentum p_s .

$$\alpha_s = \frac{E_s - p_{s\parallel}}{M_s}$$

$$\theta_{pq} > 100^\circ$$

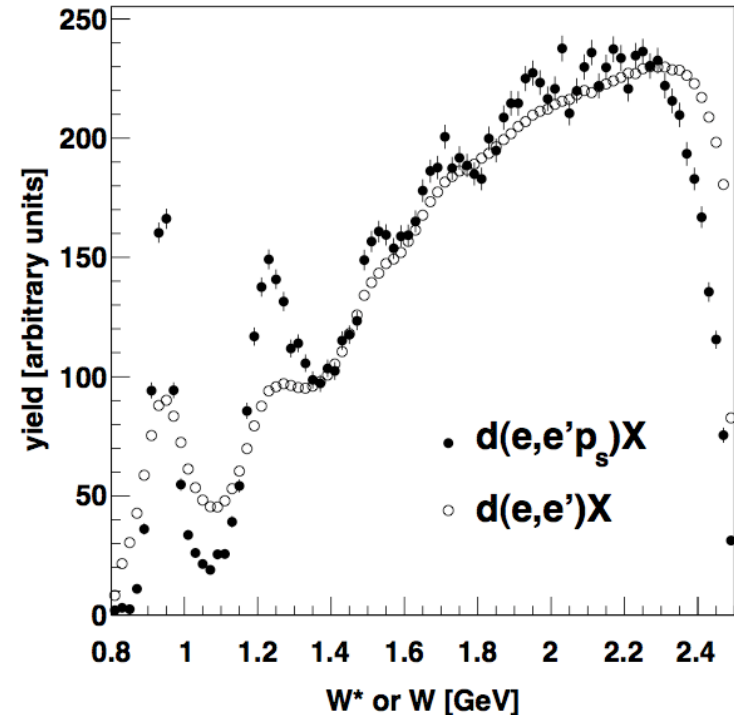
$$70 < p_s < 100 \text{ MeV/c}$$

$$E_{\text{beam}} = 4 \text{ \& 5 GeV}$$

$$M^{*2} = (M_d - E_s)^2 - \vec{p}_s^2$$

$$W^{*2} \approx M^{*2} - Q^2 + 2Mv(2 - \alpha_s)$$

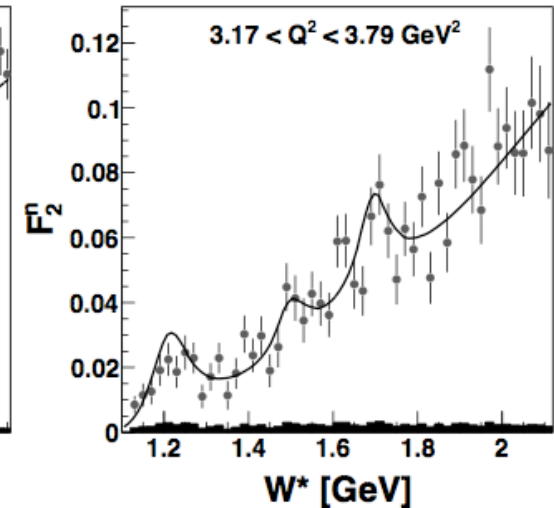
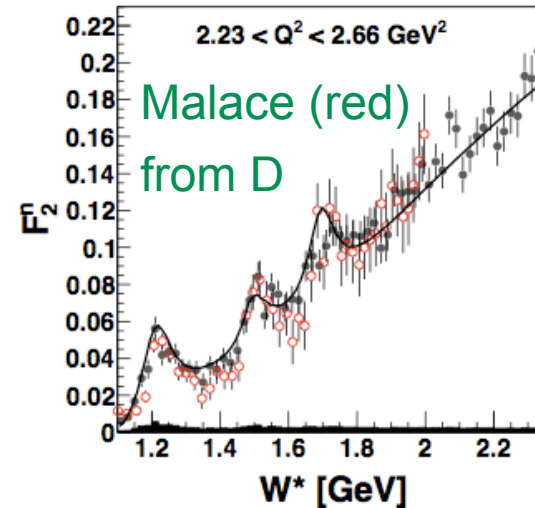
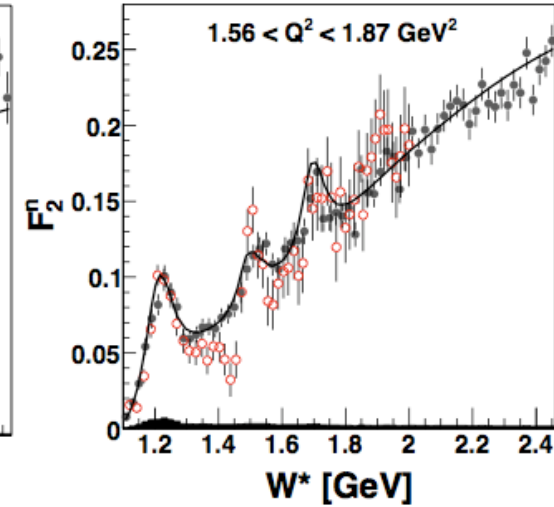
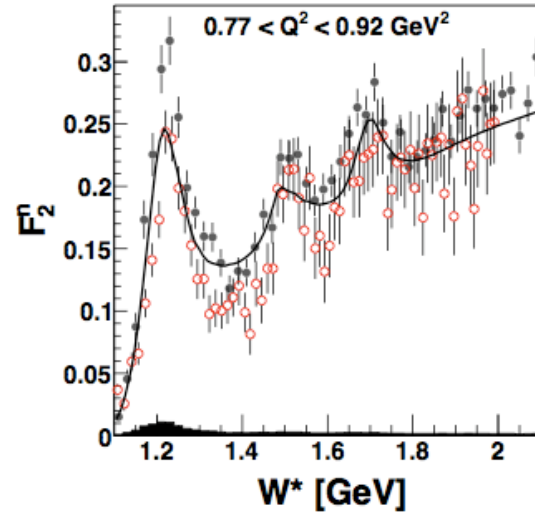
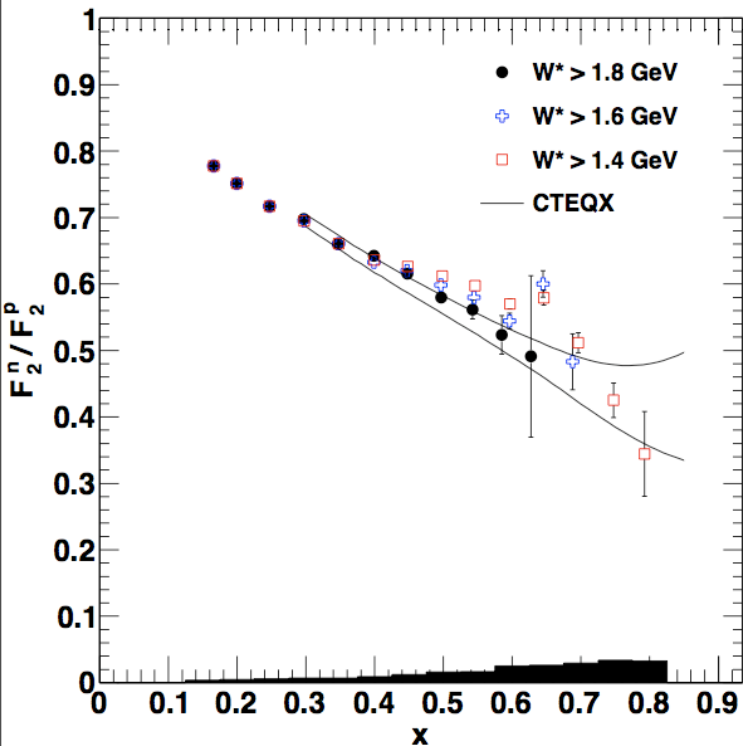
$$x^* = \frac{Q^2}{2p_N^\mu q^\mu} \approx \frac{Q^2}{2Mv(2 - \alpha_s)} = \frac{x}{2 - \alpha_s}$$





★WG1PS2: Slava Tkachenko *Model independent extraction of neutron structure functions from deuterium data*

- First data from a 'free' neutron target
- Black line (right) is Bosted/Christy model
- Black lines (below) are the CTEQX errors band



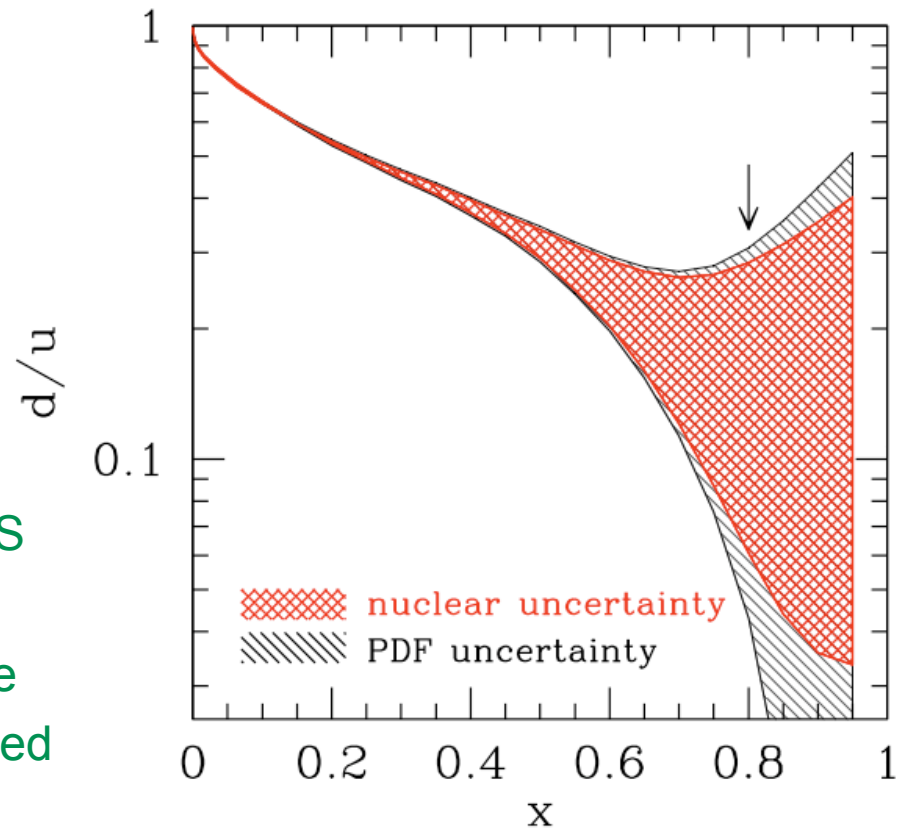


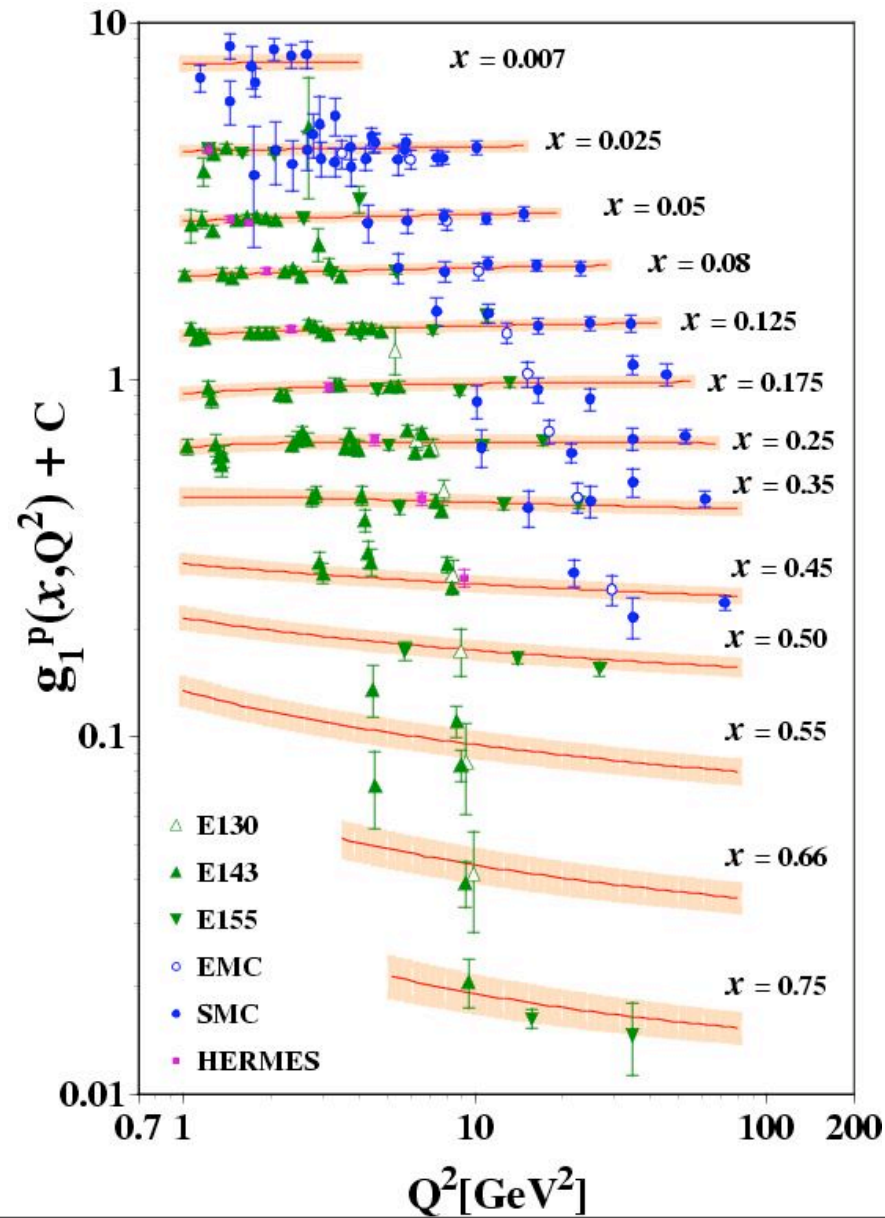
★WG1PS2: Jeff Owens *Uncertainties in determining the d quark PDF at large values of x*

CTEQ-JLab (CJ) collaboration

Accardi, Christy, Keppel, Melnitchouk, Monaghan, Morfin,
Owens, Zhu

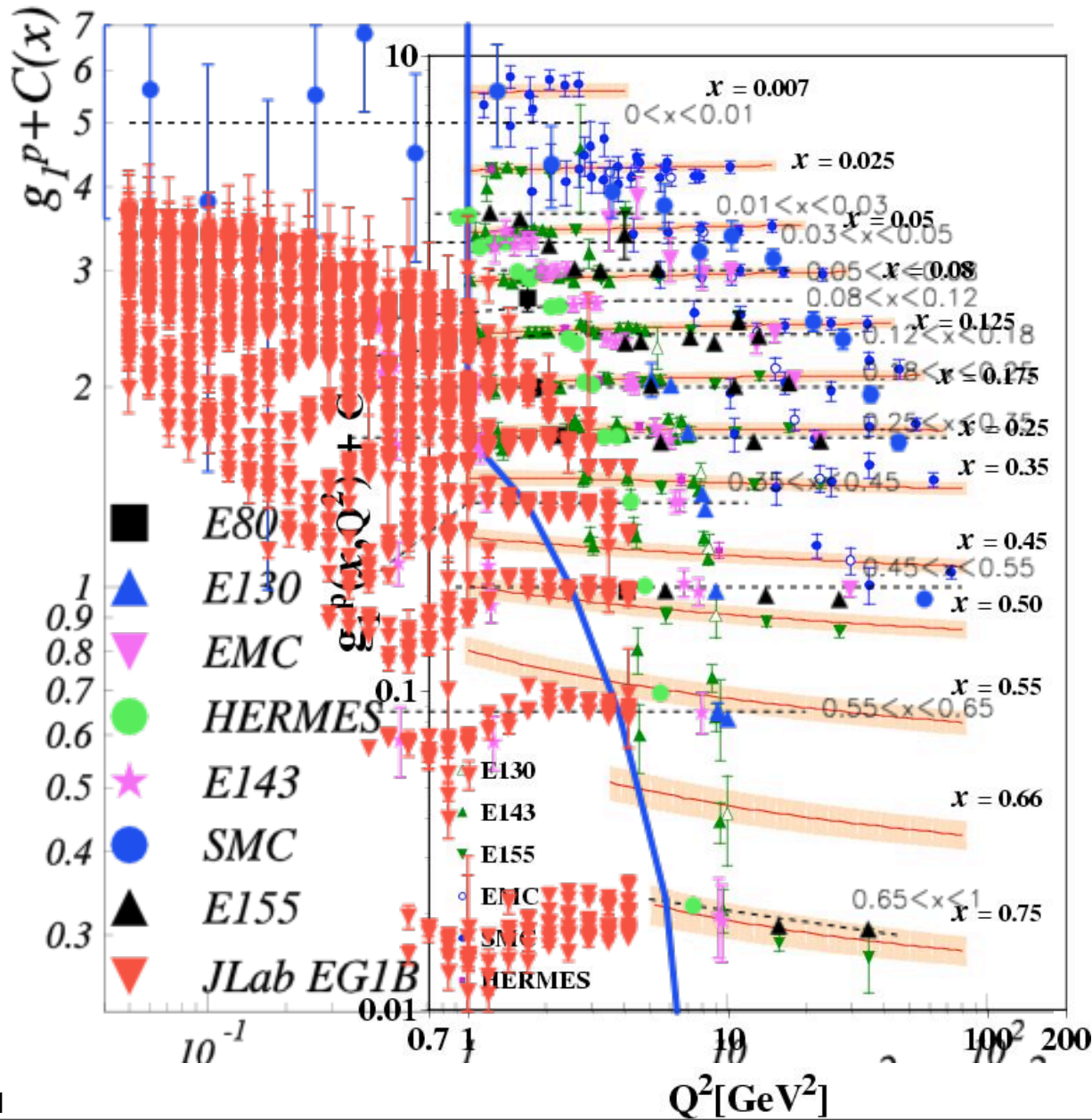
- Focus on large- x , small- Q^2 region
 - fully exploit SLAC and JLab data
 - reduced PDF uncertainty at large x
- Flexible d-quark parametrization
 - extract d/u ratio at $x \rightarrow 1$
- Large nuclear uncertainty
 - d-quark (and gluons)!
 - need BONUS12, MARATHON, PV-DIS
- Polarized PDFs:
 - New JLab Theory/Experiment initiative
 - Use current & future data over extended x & Q^2 range







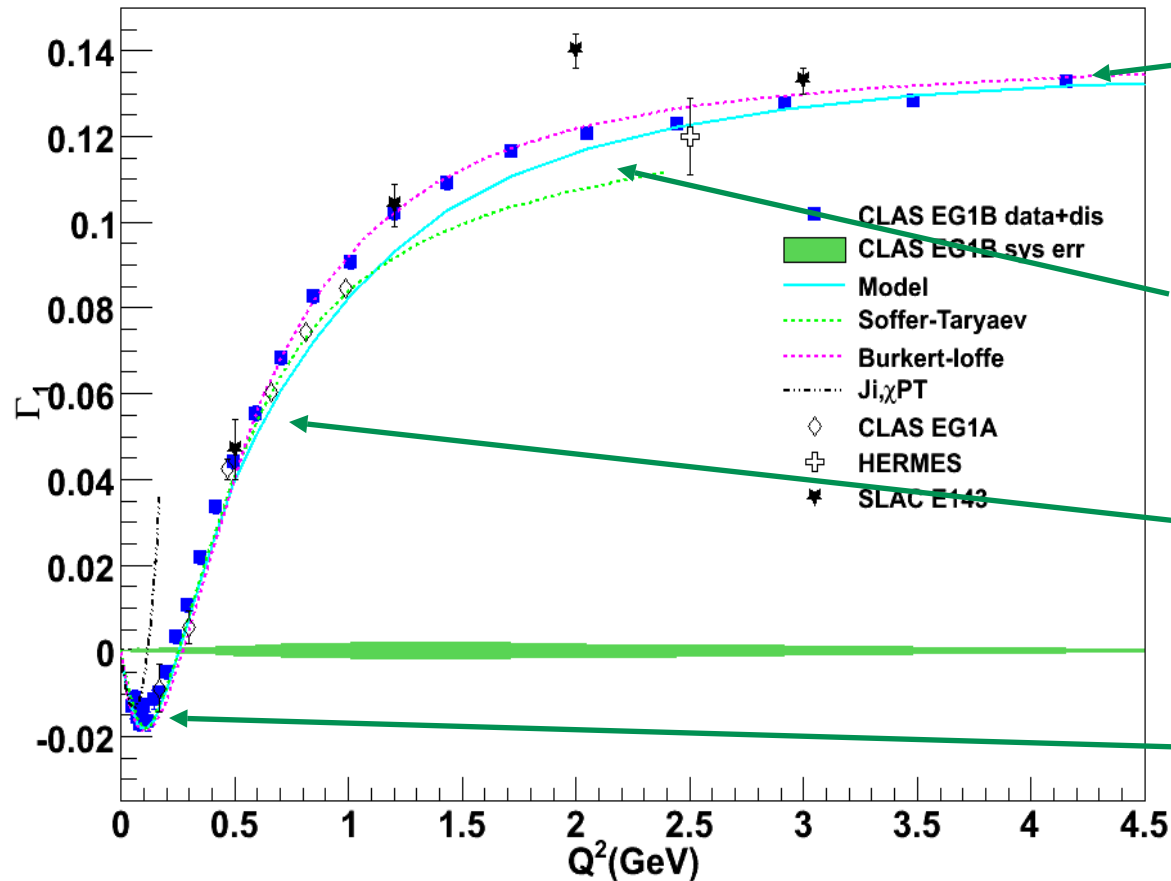
★WG7PS3:Yelena Prok





$\Gamma_1(\mathbf{P})$

$$\Gamma_1^{p,d}(Q^2) = \int_0^1 g_1^{p,d}(x, Q^2) dx$$



scaling: $\ln Q^2$

higher twist: $(1/Q^2)^n$
target mass corrections

dominating resonances

χ PT: $(Q^2)^n$

Jefferson Lab @ 6 GeV explores the transition from DIS to χ PT



CLAS data make moments possible

Osipenko et al, NPA845(10)1
Nachtmann moments for ¹²C

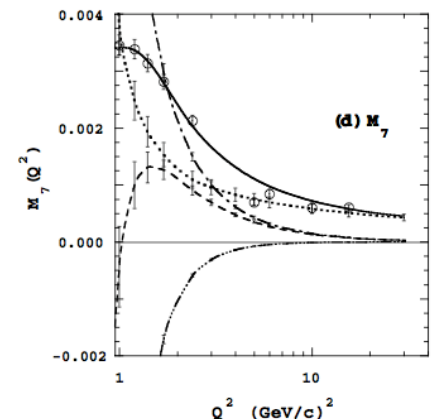
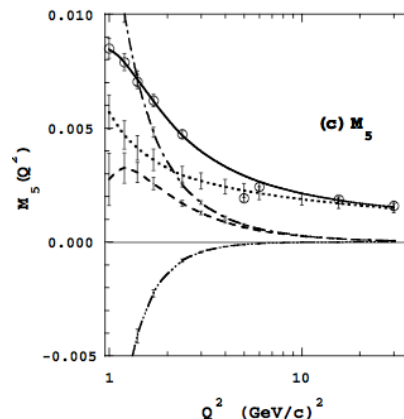
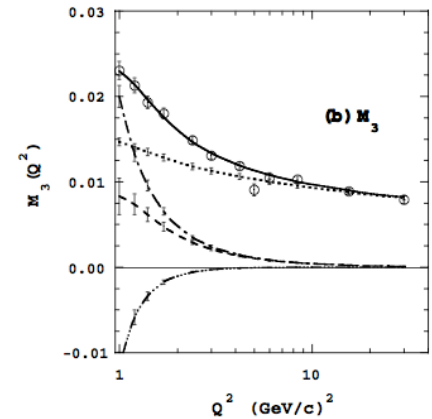
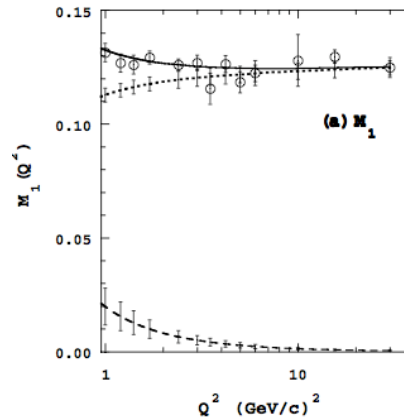
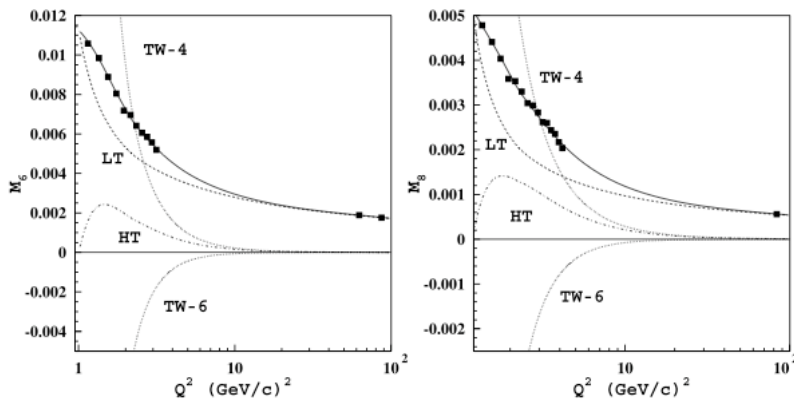
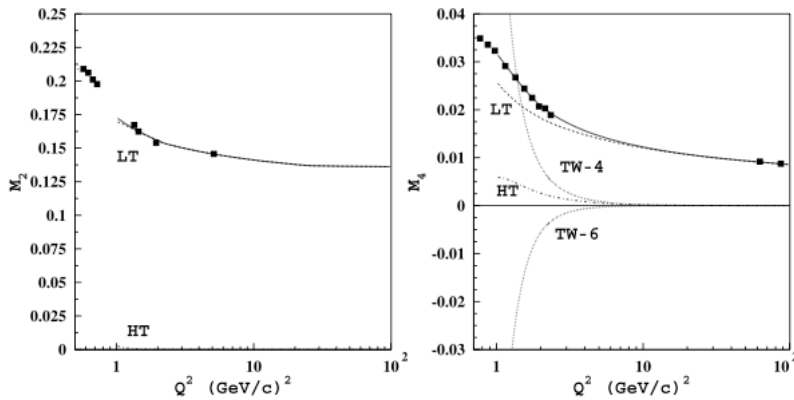
$$M_1(Q^2) = \int_0^1 dx \frac{\xi^2}{x^2} \left\{ g_1(x, Q^2) \left(\frac{x}{\xi} - \frac{1}{9} \frac{M^2 x \xi}{Q^2} \right) - g_2(x, Q^2) \frac{4}{3} \frac{M^2 x^2}{Q^2} \right\}$$

$$\xi = 2x / (1 + \sqrt{1 + 4M^2 x^2 / Q^2})$$

Osipenko et al, PRD71(05)054007
Polarized Nachtmann moments

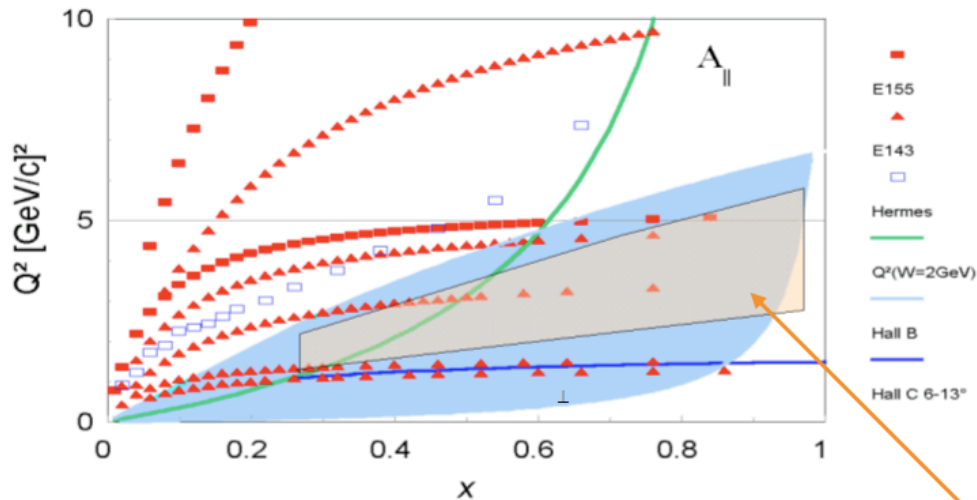
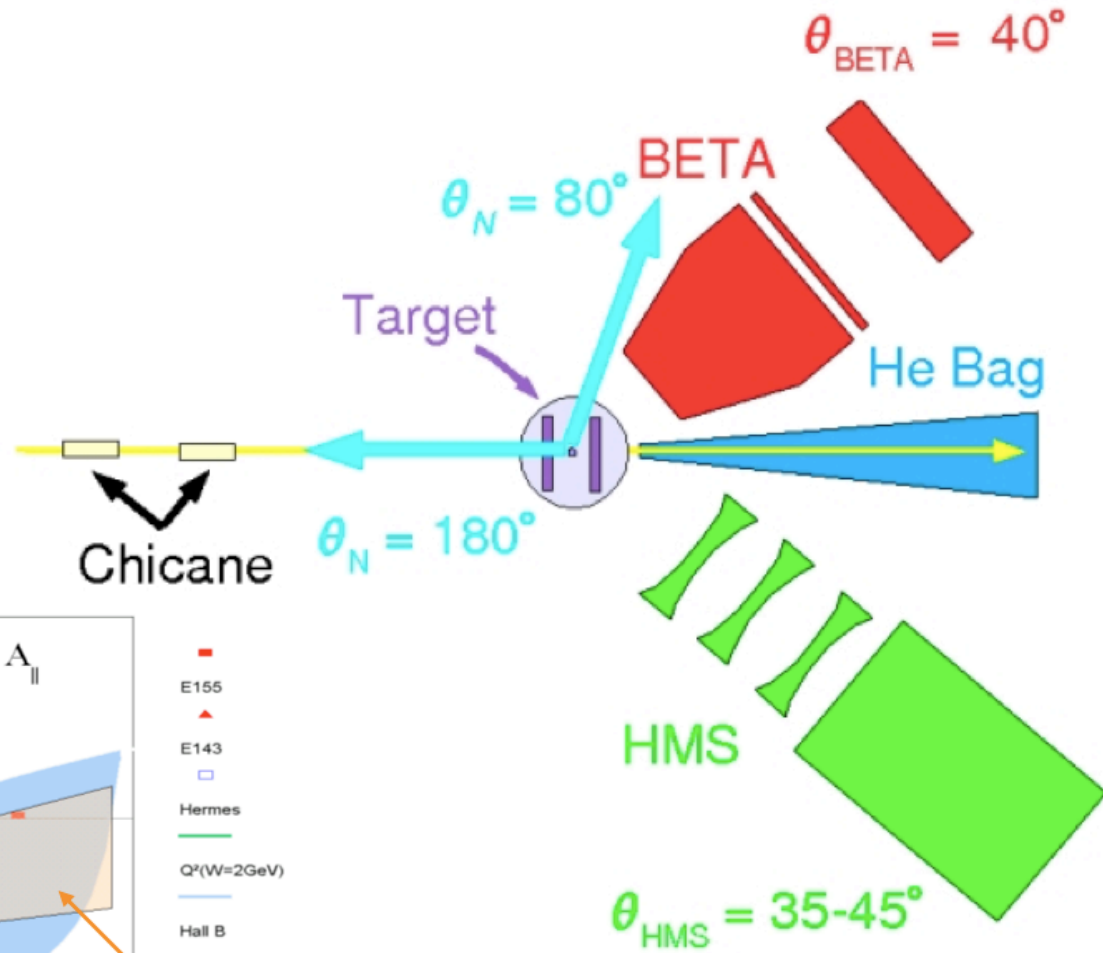
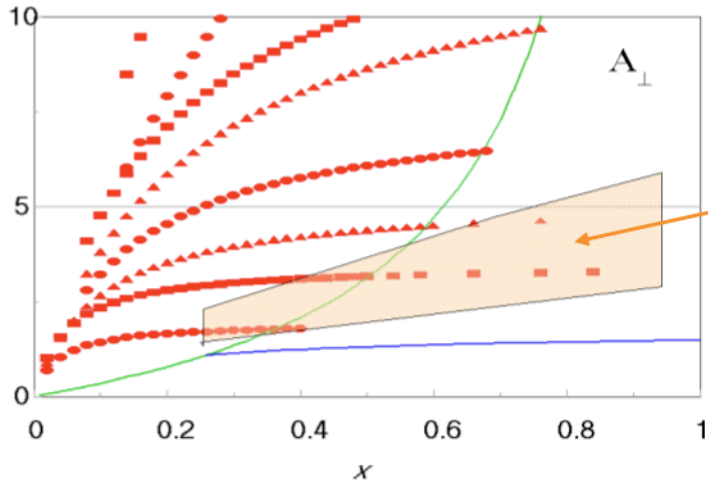
M. Osipenko et al. / Nuclear Physics A 845 (2010) 1–32

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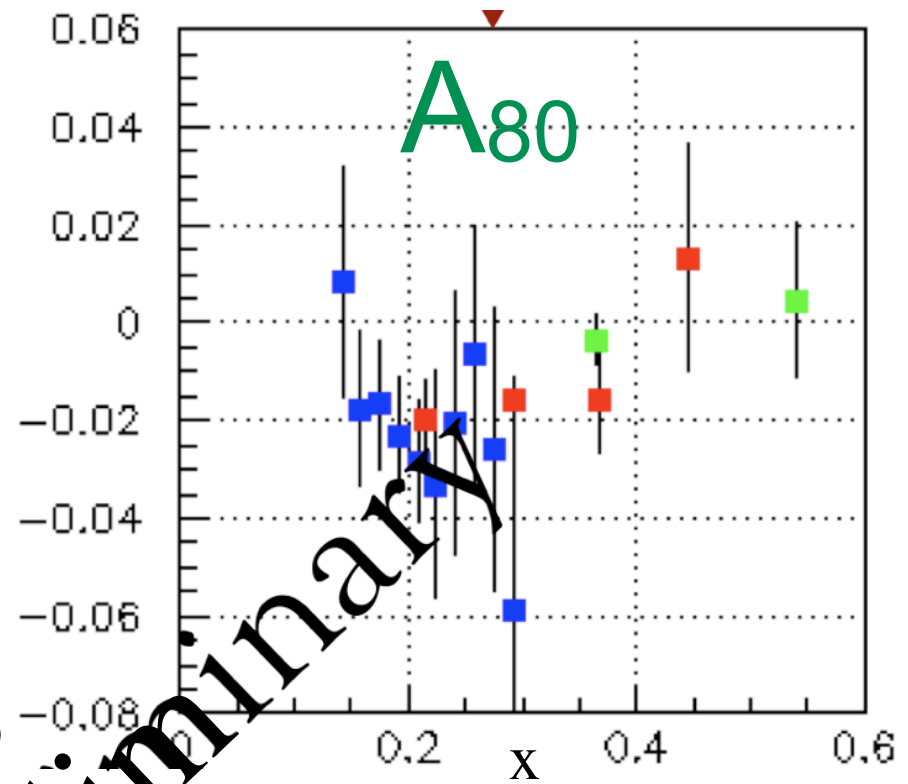
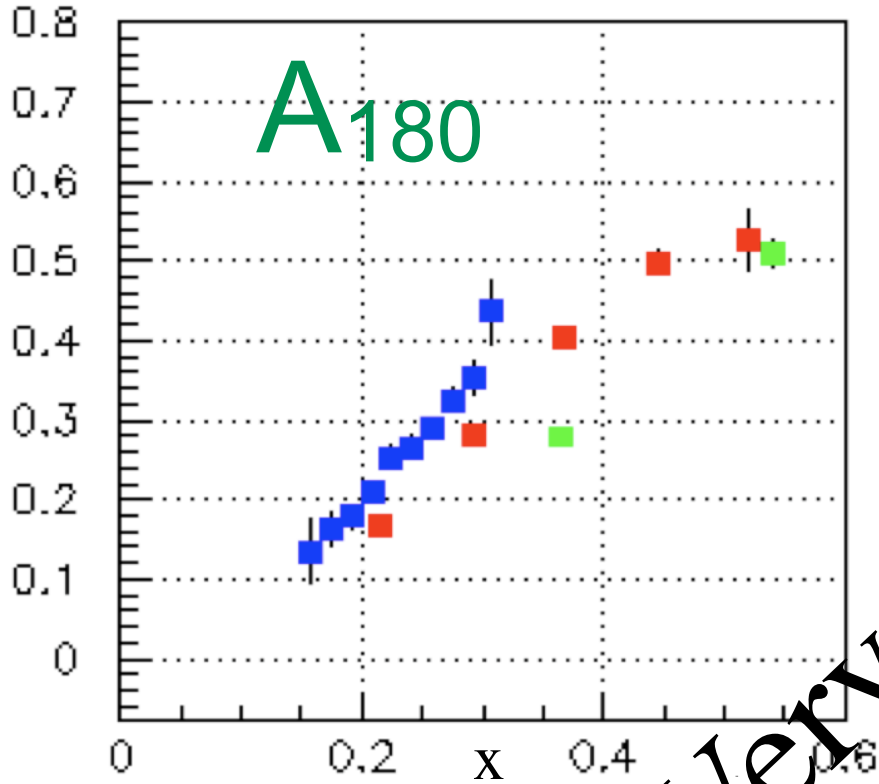
★WG6PSH1: Hovhannes Baghdasaryan *Preliminary proton spin asymmetry results from SANE*



- E155
- ▲ E143
- Hermes
- $Q^2(W=2\text{GeV})$
- Hall B
- Hall C 6-13°



★WG6PSH1: Hovhannes Baghdasaryan *Preliminary proton spin asymmetry results from SANE*



$Q^2 = 1.7$ GeV²

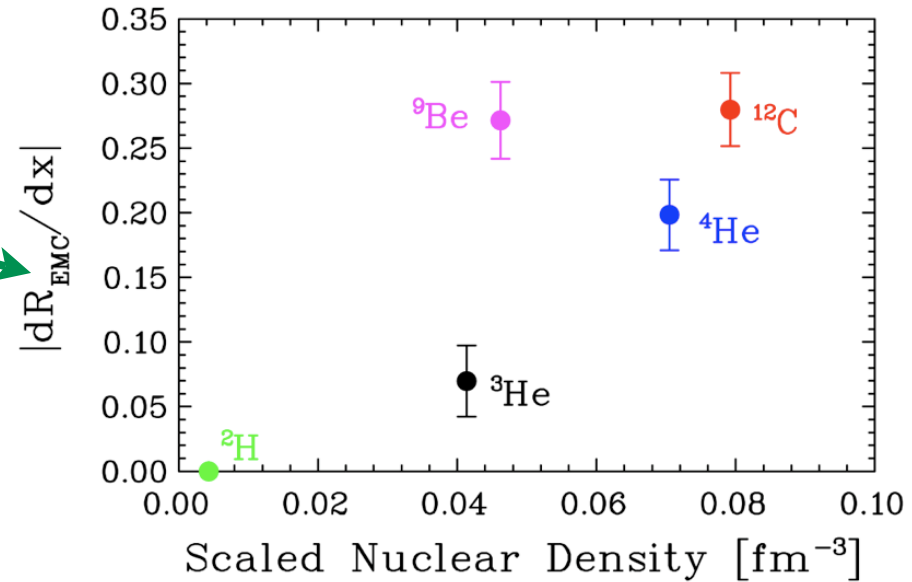
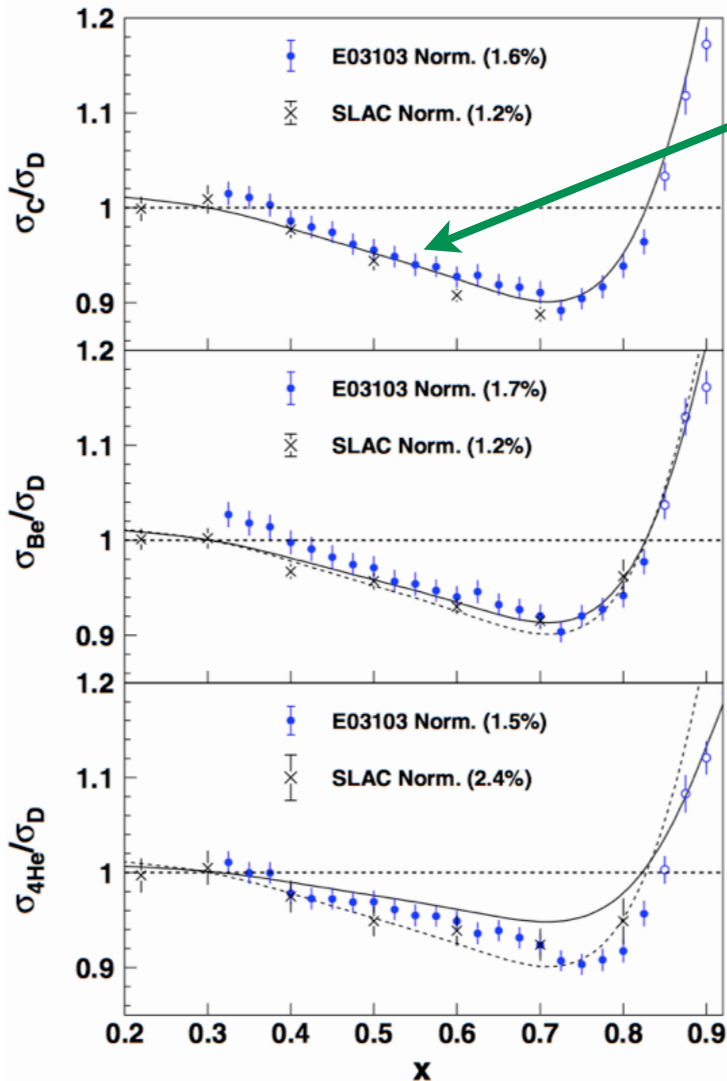
$Q^2 = 2.5$ GeV²

$Q^2 = 3.5$ GeV²

Very Preliminary



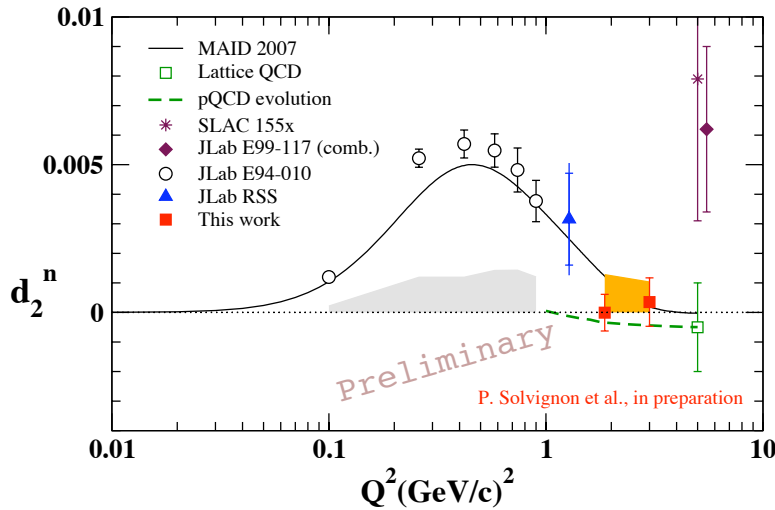
Seely et al, PRL 103(09)202301



- Recent data from JLab are very precise.
- Slope with respect to x is used to characterize the strength of the effect
- ^9Be anomaly is used to argue for a local-density origin of the effect



★WG6PSH2: Nilanga Liyanage *Moments of the neutron g_2 structure function and ... higher-twist effects*



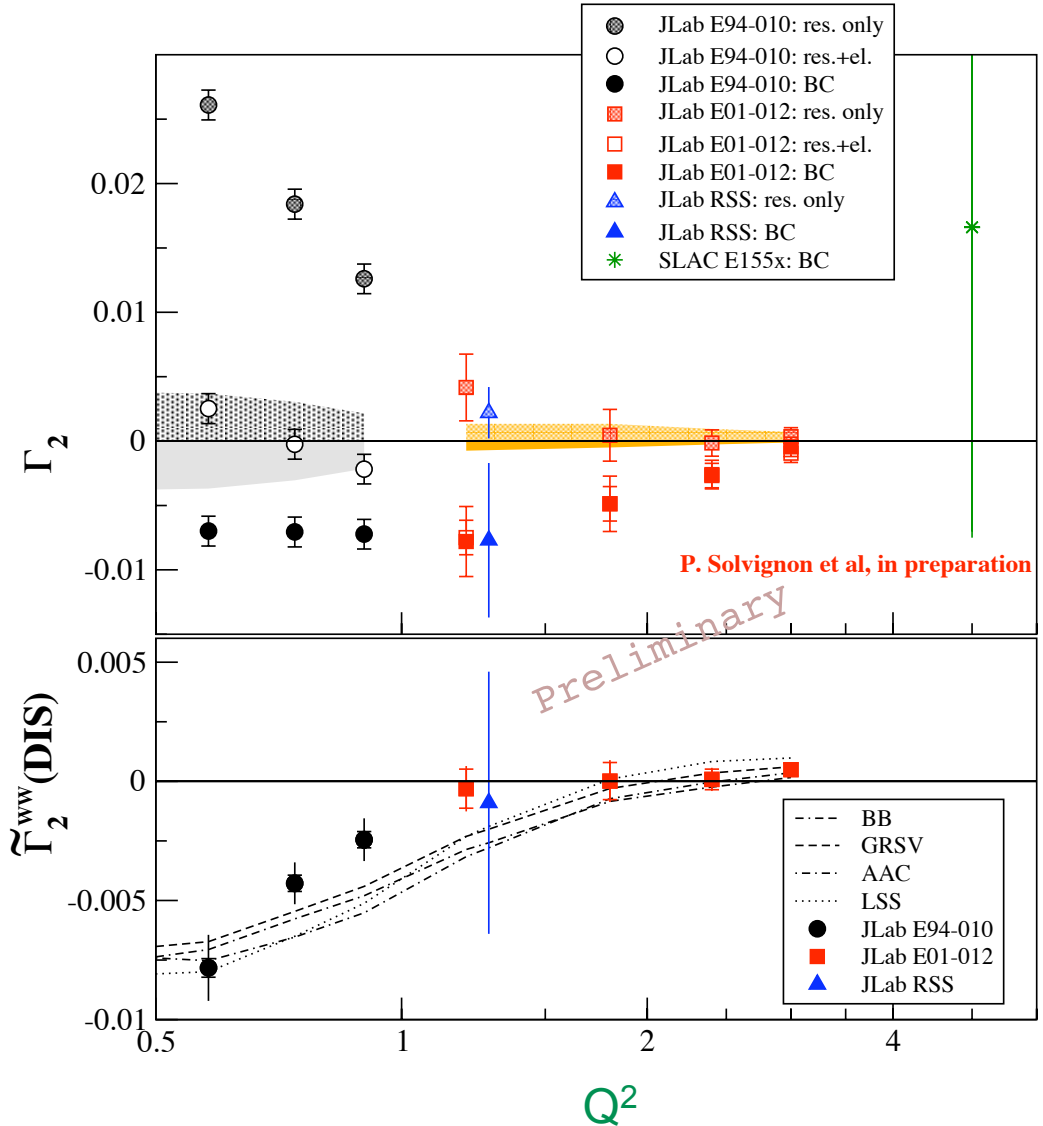
- higher twist coefficient

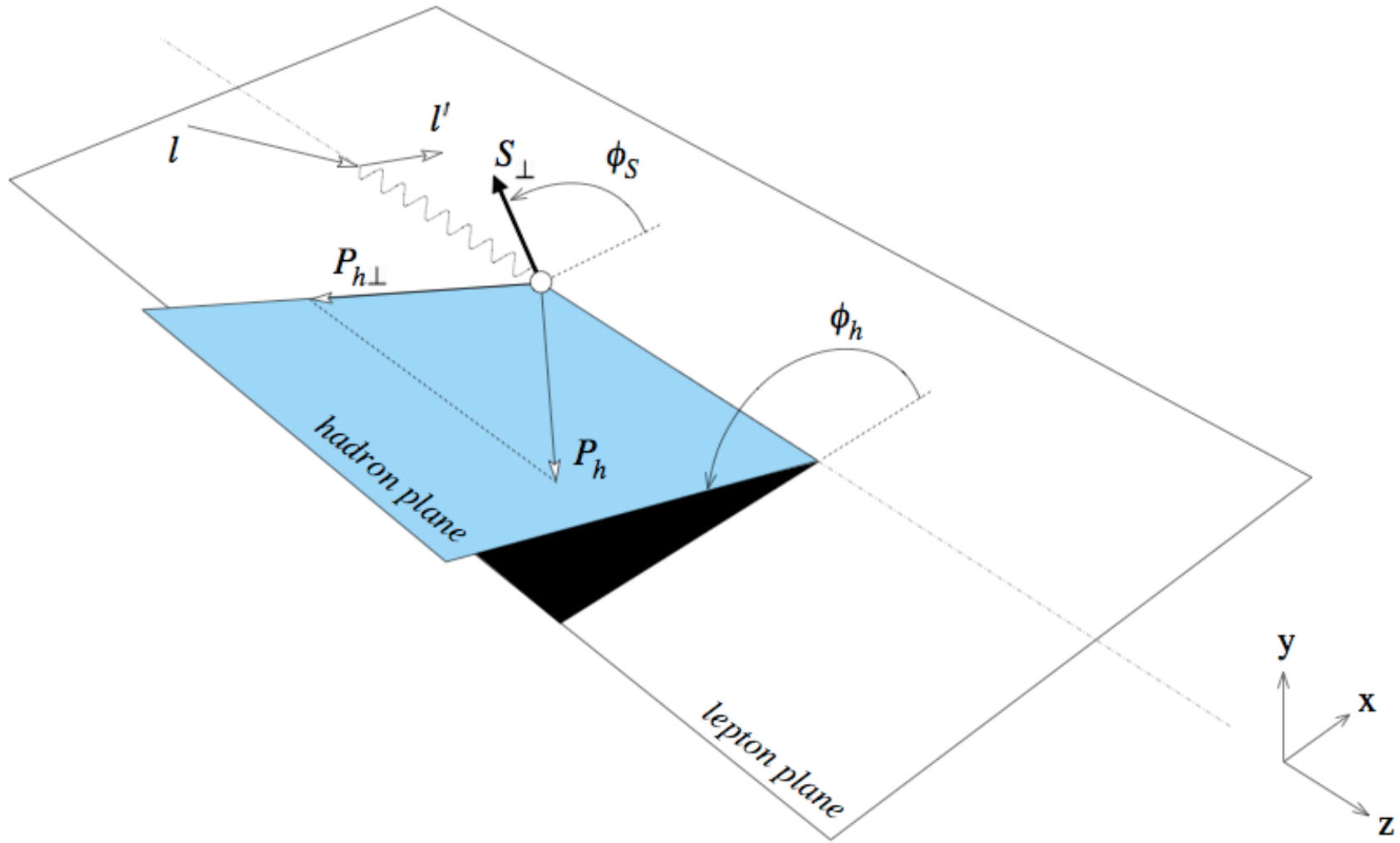
$$d_2(Q^2) = \int_0^1 dx x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)]$$

- Burkhardt-Cottingham Sum Rule

$$\int_0^1 g_2(x, Q^2) dx = 0$$

- Γ_2^{WW} is sum of g_2^{WW} for $W > 2$ GeV







Bacchetta, et al., JHEP 2(2007)093

Unpolarized and Longitudinally polarized

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \boxed{F_{UU,T}} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right.$$

$$+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$+ S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h \boxed{F_{UL}^{\sin\phi_h}} + \varepsilon \sin(2\phi_h) \boxed{F_{UL}^{\sin 2\phi_h}} \right]$$

$$\left. + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} \boxed{F_{LL}} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \right\}$$

$A_{UL} = \{\text{UL terms}\} / \{\text{UU terms}\}$

$A_{LL} = \{\text{LL terms}\} / \{\text{UU terms}\}$

← = Higher Twist



Bacchetta, et al., JHEP 2(2007)093

Transverse target polarizations

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \right.$$

0 at high Q^2 = Higher Twist

$$+ |\mathbf{S}_\perp| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \left. \right]$$

$$+ \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ |\mathbf{S}_\perp| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right.$$

$$\left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},$$



The observables are the structure functions such as $F^{\sin\phi}_{UL}$, not the transverse momentum distributions (TMDs) or fragmentation functions (FFs). Four-fold differential data in x , z , Q^2 and P_T are essential to allow modeling of TMDs and FFs.

$$C[wfD] = x \sum_a e_a^2 \int d^2\mathbf{p}_T d^2\mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{h\perp}/z) w(\mathbf{p}_T, \mathbf{k}_T) f^a(x, p_T^2) D^a(z, k_T^2),$$

$$F_{UL}^{\sin\phi_h} = \frac{2M}{Q} C \left[-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x h_L H_1^\perp + \frac{M_h}{M} g_{1L} \frac{\tilde{G}^\perp}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x f_L^\perp D_1 - \frac{M_h}{M} h_{1L}^\perp \frac{\tilde{H}}{z} \right) \right]$$

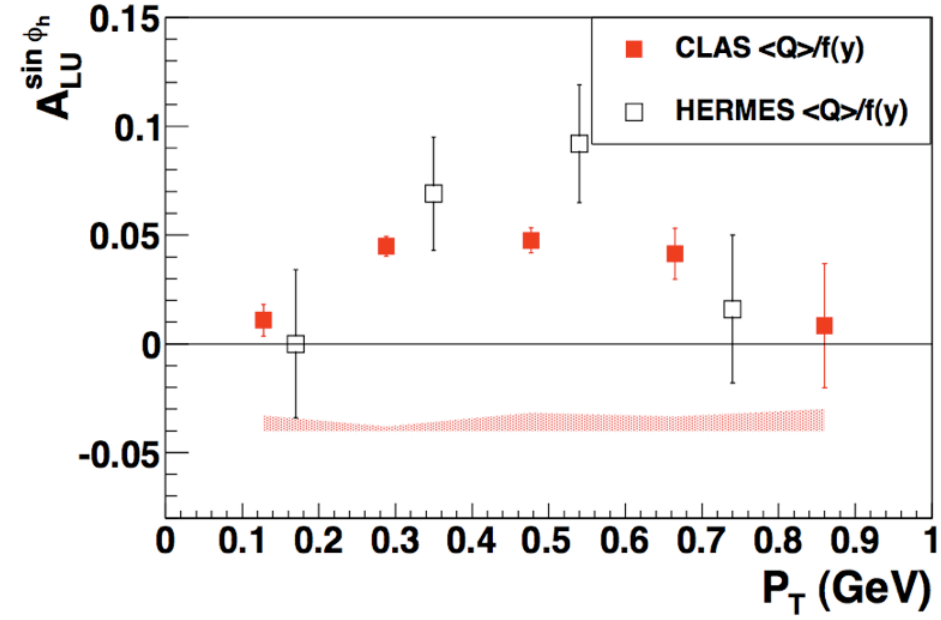
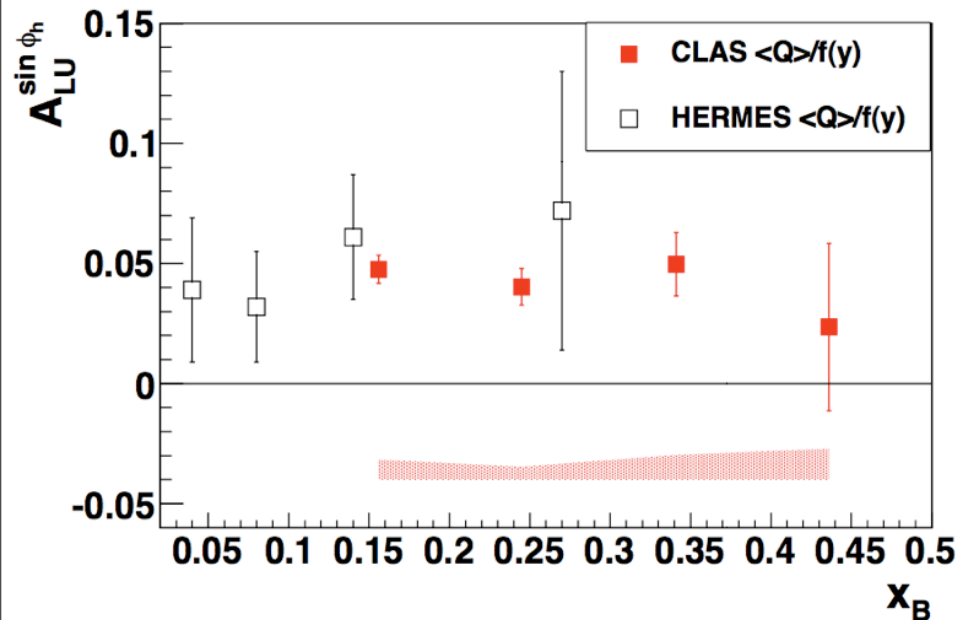
$$F_{UL}^{\sin 2\phi_h} = C \left[-\frac{2(\hat{\mathbf{h}} \cdot \mathbf{k}_T)(\hat{\mathbf{h}} \cdot \mathbf{p}_T) - \mathbf{k}_T \cdot \mathbf{p}_T}{MM_h} h_{1L}^\perp H_1^\perp \right],$$

$$F_{LL} = C[g_{1L}D_1]$$

$$F_{LL}^{\cos\phi_h} = \frac{2M}{Q} C \left[\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x e_L H_1^\perp - \frac{M_h}{M} g_{1L} \frac{\tilde{D}^\perp}{z} \right) - \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x g_L^\perp D_1 + \frac{M_h}{M} h_{1L}^\perp \frac{\tilde{E}}{z} \right) \right]$$



- Mher Aghasyan et al., E01-113 in preparation for publication
- CLAS data for $A_{UL}^{\pi^0}$
- Unpolarized liquid hydrogen target
- Beam energy of 5.776 GeV
- $Q^2 > 1$; $0.4 < z < 0.7$



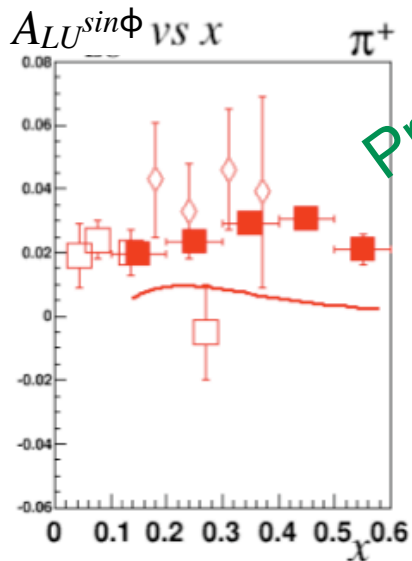


★WG6PST3: Wes Gohn *Beam single spin asymmetries in SIDIS from an unpolarized proton*

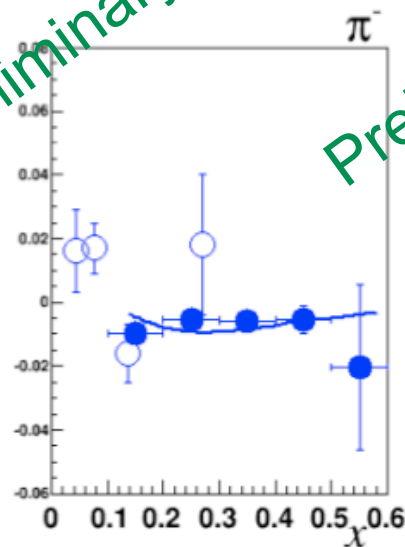
- CLAS data from E1f run period (2003)
- Unpolarized liquid hydrogen target
- Longitudinal beam polarization of 75%
- Beam energy of 5.498 GeV

CLAS, Avakian et al, PRC69(04)042201

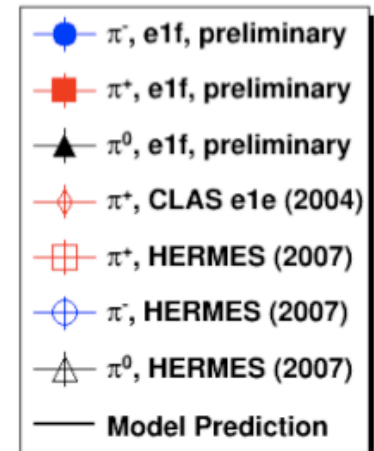
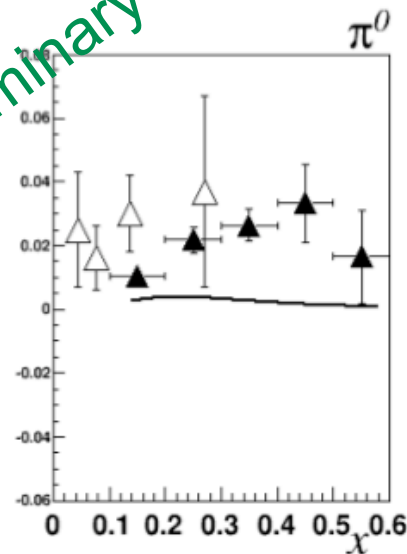
HERMES, Airapetian et al, PLB648(07)164



Preliminary

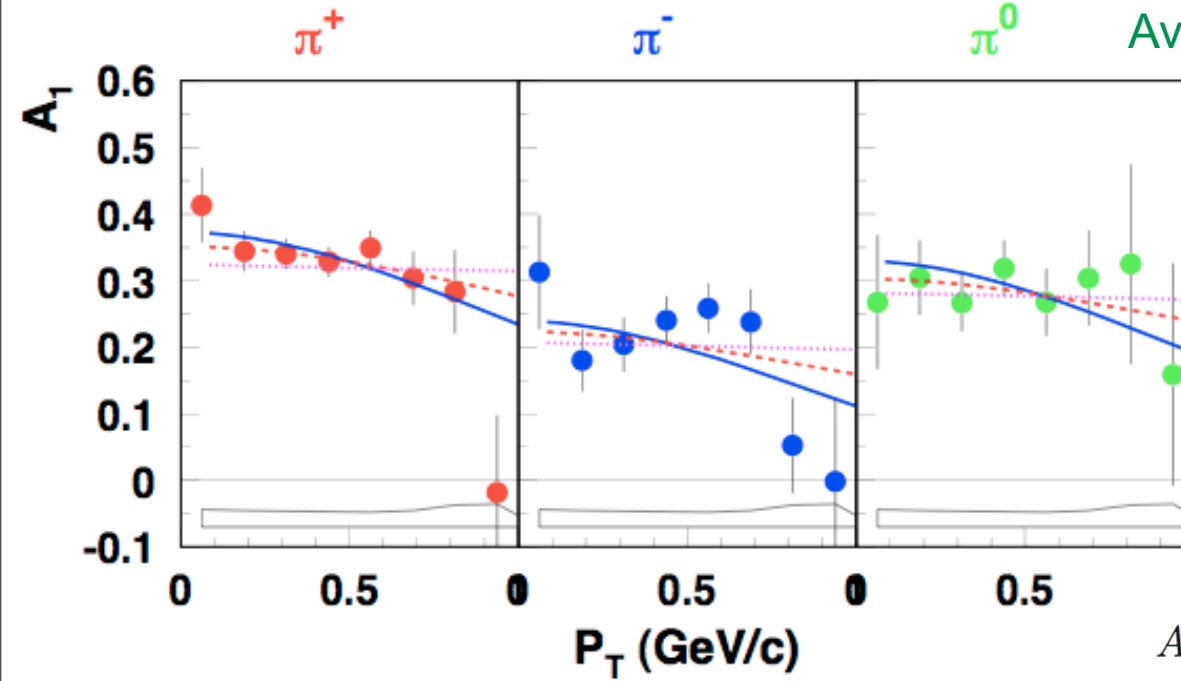


Preliminary





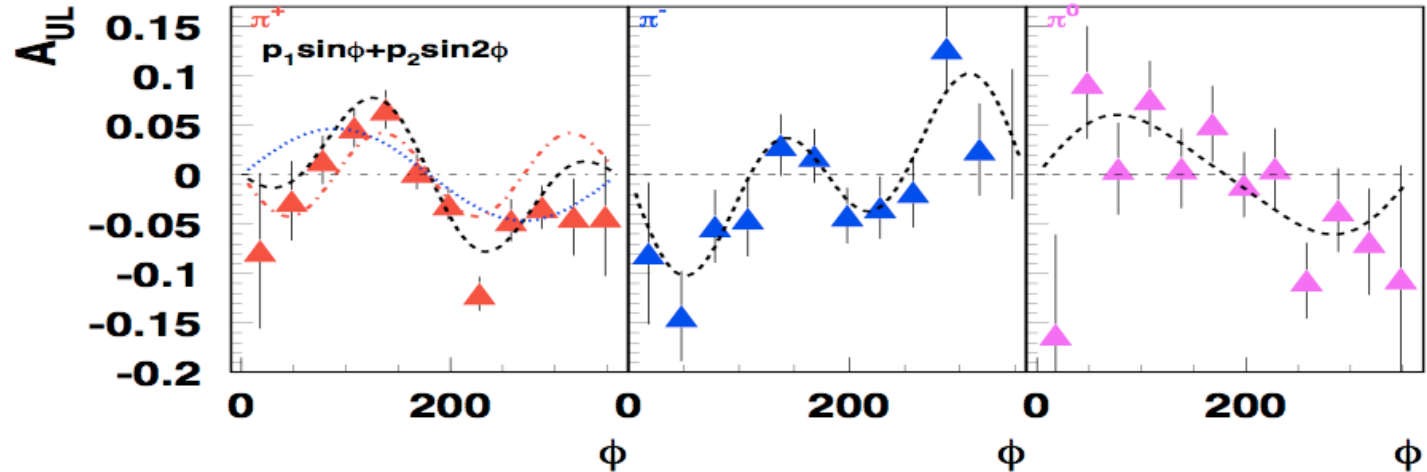
Avakian et al., PRL105(10)262002



$$A_1 \equiv \frac{F_{LL}}{F_{UU,T}}$$

CLAS

$$A_{UL}^{\sin 2\phi}(x) = \frac{\int dy [\cos \theta_\gamma (1-y)/Q^4] F_{UL}^{\sin 2\phi}}{\int dy [(1-y + \frac{1}{2}y^2)/Q^4] F_{UU,T}}$$





★WG6PST3: Sucheta Jawalkar *Spin azimuthal asymmetries on longitudinally polarized proton*

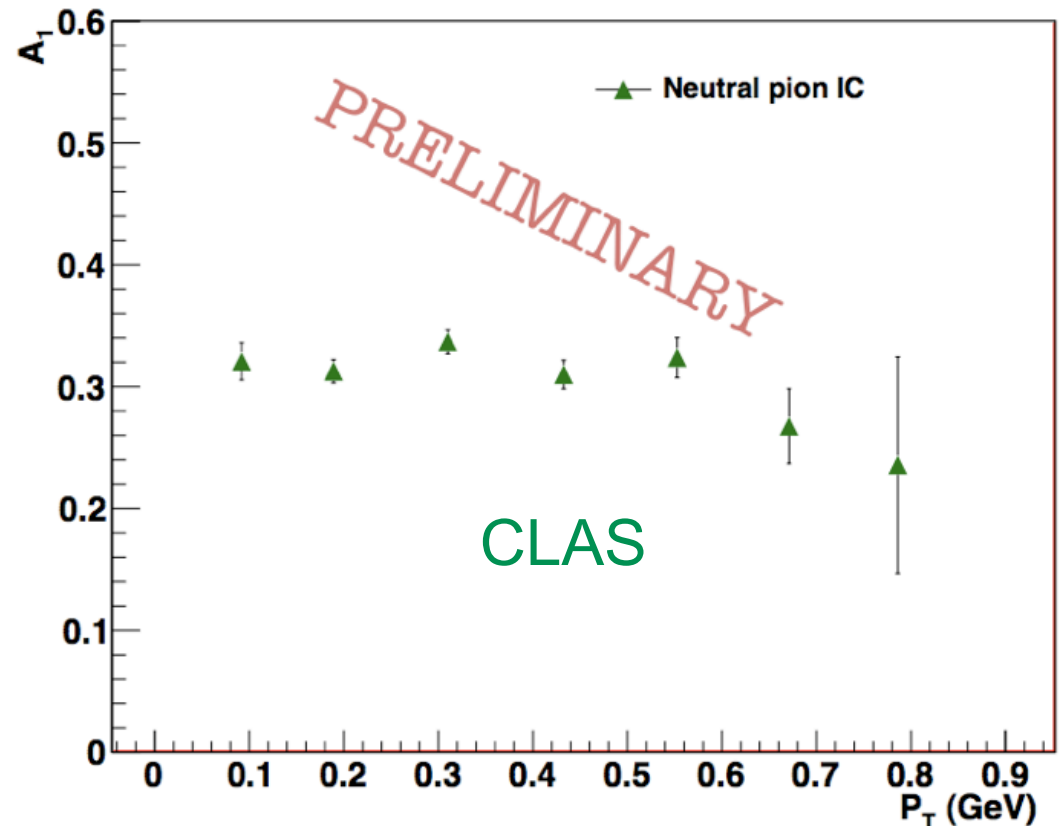
$$f_1^q(x, k_T) = f_1(x) \frac{1}{\pi \mu_0^2} \exp\left(-\frac{k_T^2}{\mu_0^2}\right)$$

$$g_1^q(x, k_T) = g_1(x) \frac{1}{\pi \mu_2^2} \exp\left(-\frac{k_T^2}{\mu_2^2}\right)$$

$$D_1^q(z, p_T) = D_1(z) \frac{1}{\pi \mu_D^2} \exp\left(-\frac{p_T^2}{\mu_D^2}\right),$$

$$\frac{g_1}{F_1} \propto \frac{\sum_q e_q^2 g_1^q(x) D_1^{q \rightarrow \pi}(z)}{\sum_q e_q^2 f_1^q(x) D_1^{q \rightarrow \pi}(z)} e^{-z^2 P_T^2 \frac{(\mu_0^2 - \mu_2^2)}{(\mu_D^2 + z^2 \mu_0^2)(\mu_D^2 + z^2 \mu_2^2)}}$$

- eg1-dvcs data
- P_T dependence $\rightarrow \mu_0 \neq \mu_2$
- For π^+ , π^- and π^0





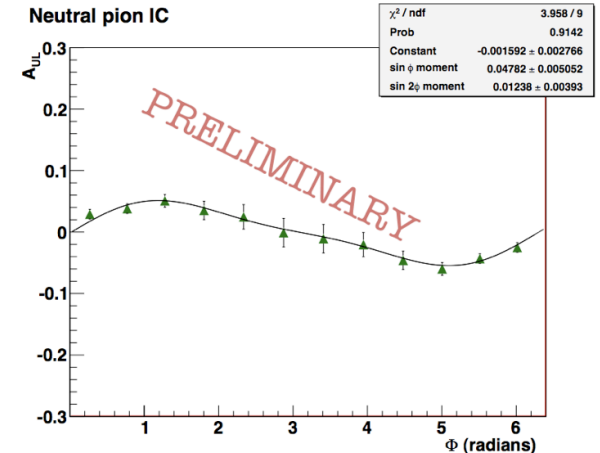
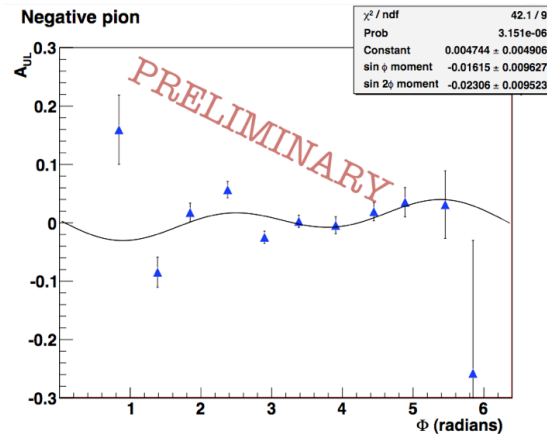
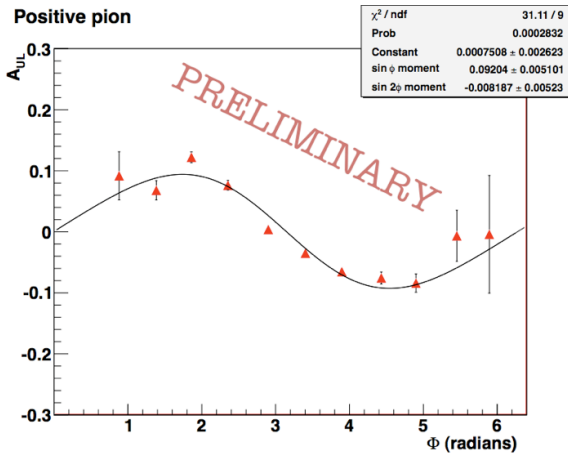
★WG6PST3: Sucheta Jawalkar *Spin azimuthal asymmetries on longitudinally polarized proton*

The target spin asymmetries as a function of ϕ have both $\sin\phi$ and $\sin 2\phi$ components.

$A_{UL}^{\sin\phi}$ (higher twist) is significant for π^+ , π^0

$A_{UL}^{\sin 2\phi}$ (leading twist) is small suggesting, like for eg1b and HERMES, that the Collins favored and unfavored fragmentation functions are nearly equal and opposite.

CLAS



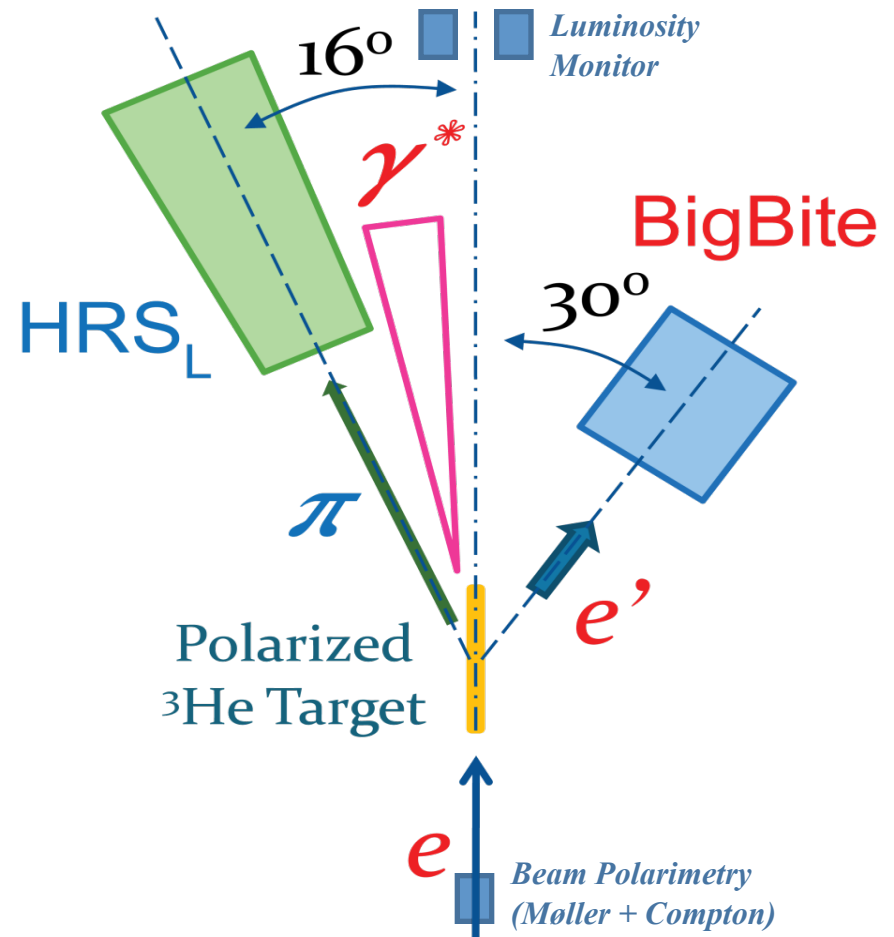


$${}^3\text{He}^\uparrow (\vec{e}, e'\pi^\pm) X$$

E06-010: Transversity

Spokespersons: J. P. Chen, E. Cisbani, H. Gao, X. Jiang,
J. C. Peng

- First measurement on n
- Polarized ${}^3\text{He}$ target

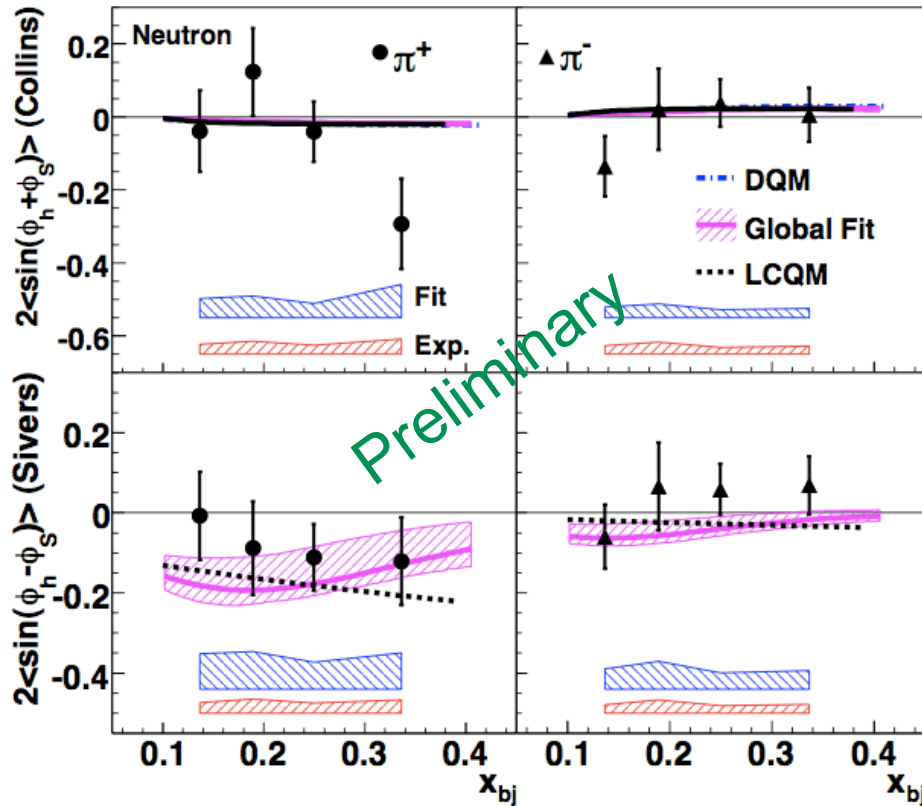




★WG6PST4: Kalyan Allada *Single spin asymmetry results from neutron*

X. Qian, et al., in preparation

A_{UT}



- Preliminary Collins/Sivers for n
- 5.9 GeV electron beam
- Polarized ³He target
- 0.14 < x < 0.35
- 1.3 < Q² < 2.7 GeV²
- Still working on systematic uncertainties
- Curves: diquark model (Ma), global fit (Anselmino), light-cone quark model (Pasquini)

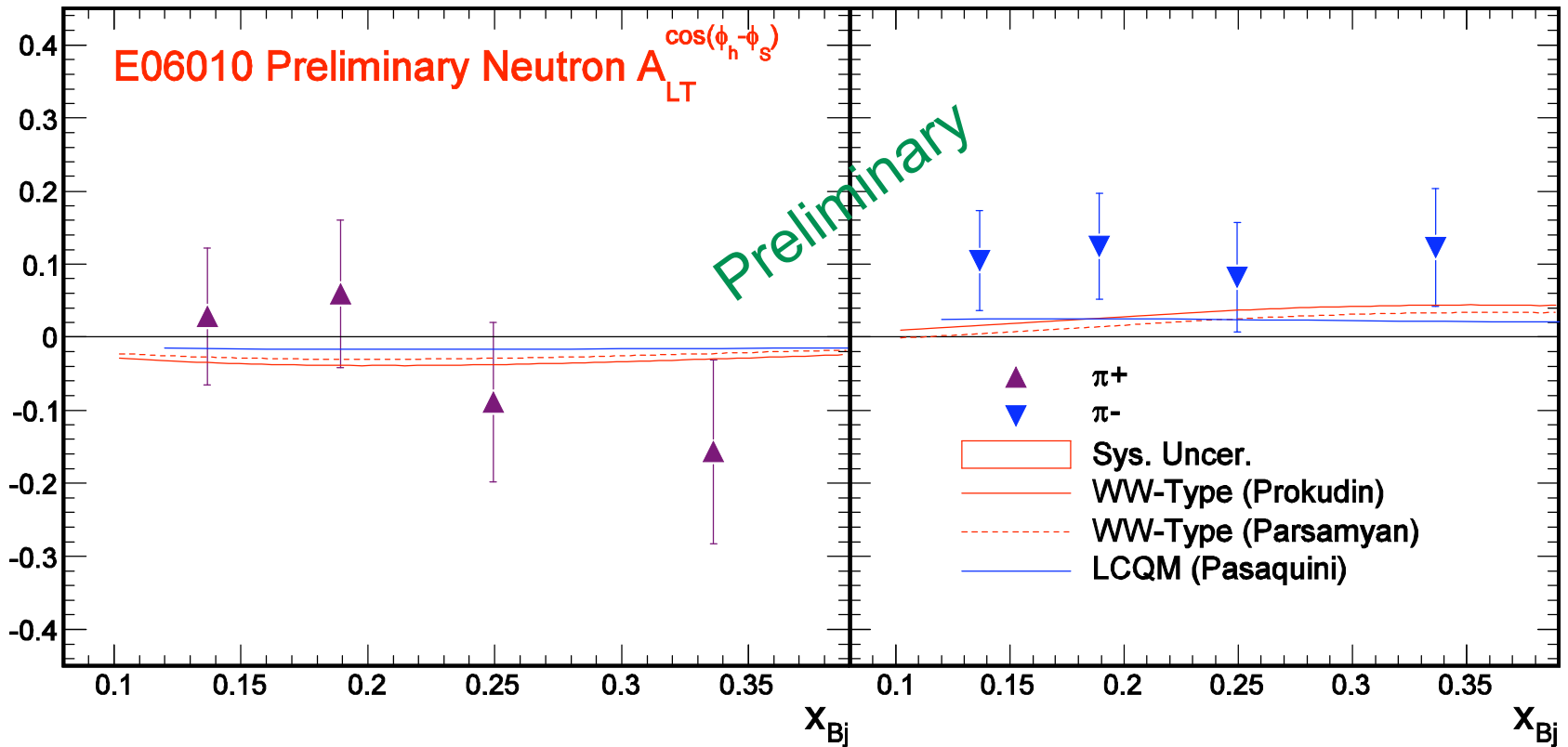
$$2 \langle \cos(\phi_h + \phi_s) \rangle \propto h_{1T}^q \otimes H_{1q}^h$$

$$2 \langle \cos(\phi_h - \phi_s) \rangle \propto f_{1T}^q \otimes D_{1q}^h$$



★WG6PST2: Jin Huang *Measurement of double spin asymmetry A_{LT}*

At leading twist: $A_{LT}^{\cos(\phi_h - \phi_s)} \propto g_{1T}^q \otimes D_{1q}^h$



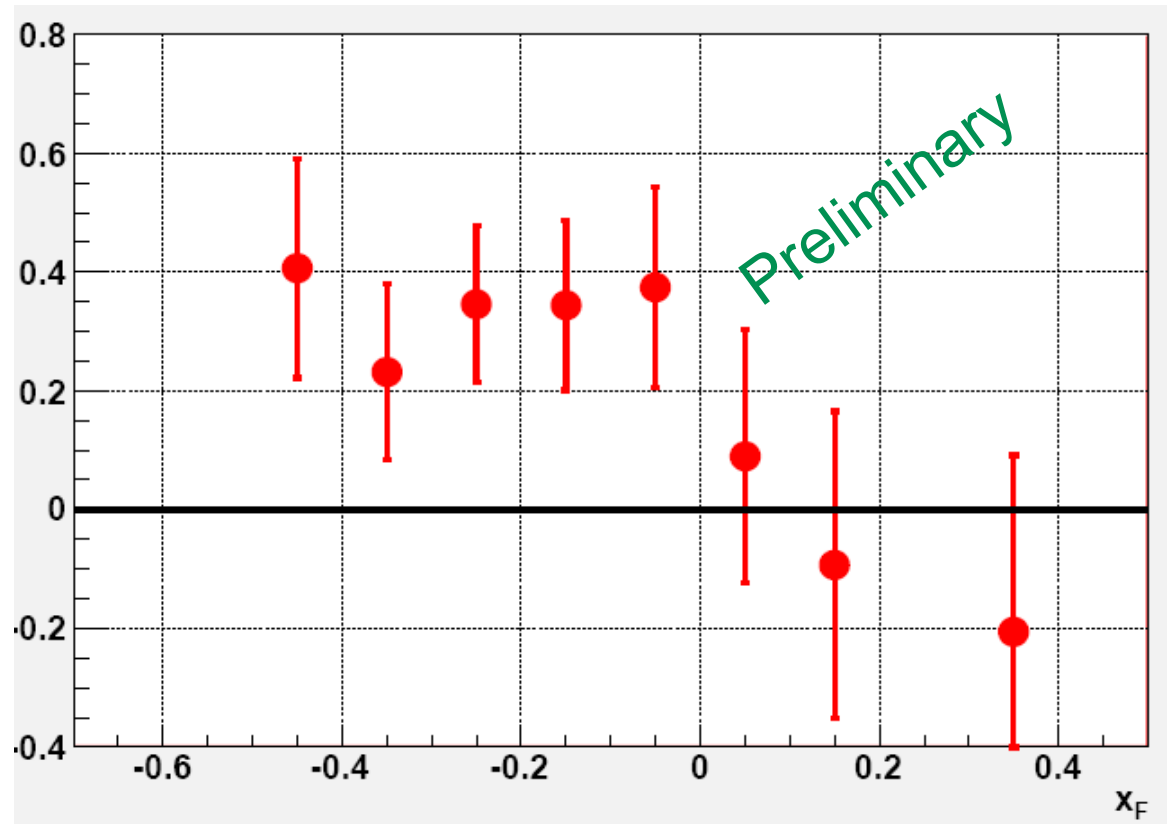


★WG6PSTV: Marco Mirazita *Lambda polarization in electroproduction at CLAS*

$$P_{\Lambda}^{measured} = P_{\Lambda}^0 + P_{\Lambda}^{\cos\phi} \langle \cos(\phi) \rangle \approx P_{\Lambda}^0 - 0.85 P_{\Lambda}^{\cos\phi}$$

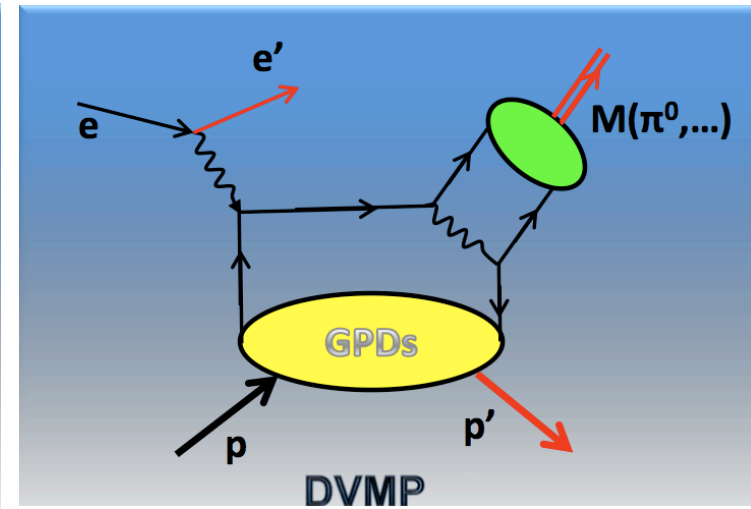
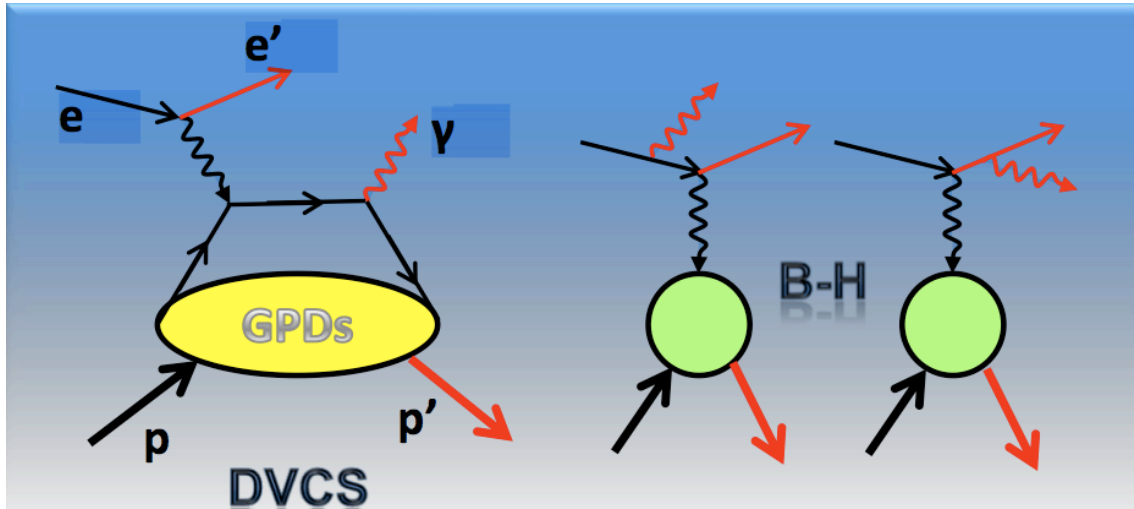
$x_F < 0$
 $P_{\Lambda}^{measured} \sim 0.3$
 $\Rightarrow P_{\Lambda}^0 \neq P_{\Lambda}^{\cos\phi}$

$x_F > 0$
 $P_{\Lambda}^{measured} \sim 0.$





★WG2PSVM: Valery Kubarovsky *Vector-mesons production and DVCS at JLab*



$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} \sim |\mathbf{T}^{\text{DVCS}} + \mathbf{T}^{\text{BH}}|^2$$

$$A_{LU} \sim \Im m \left\{ F_1 \mathcal{H} + \xi (F_1 + F_2) \tilde{\mathcal{H}} - \frac{\Delta^2}{4M^2} F_2 \mathcal{E} \right\} \sin \phi$$

suppressed

$$A_{UL} = \Im m \left\{ F_1 \tilde{\mathcal{H}} + \xi (F_1 + F_2) \left(\mathcal{H} - \frac{x_B}{2} \mathcal{E} \right) - \xi \left(\frac{x_B}{2} F_1 + \frac{\Delta^2}{4M^2} F_2 \right) \tilde{\mathcal{E}} \right\} \sin \phi$$

$\sin\phi$ moments of A_{LU} and A_{UL}
are related to linear combinations
of generalized parton distributions

$$A = \alpha \sin\phi + \beta \sin 2\phi$$

higher twist



Morrow et al, EPJA39(09)5

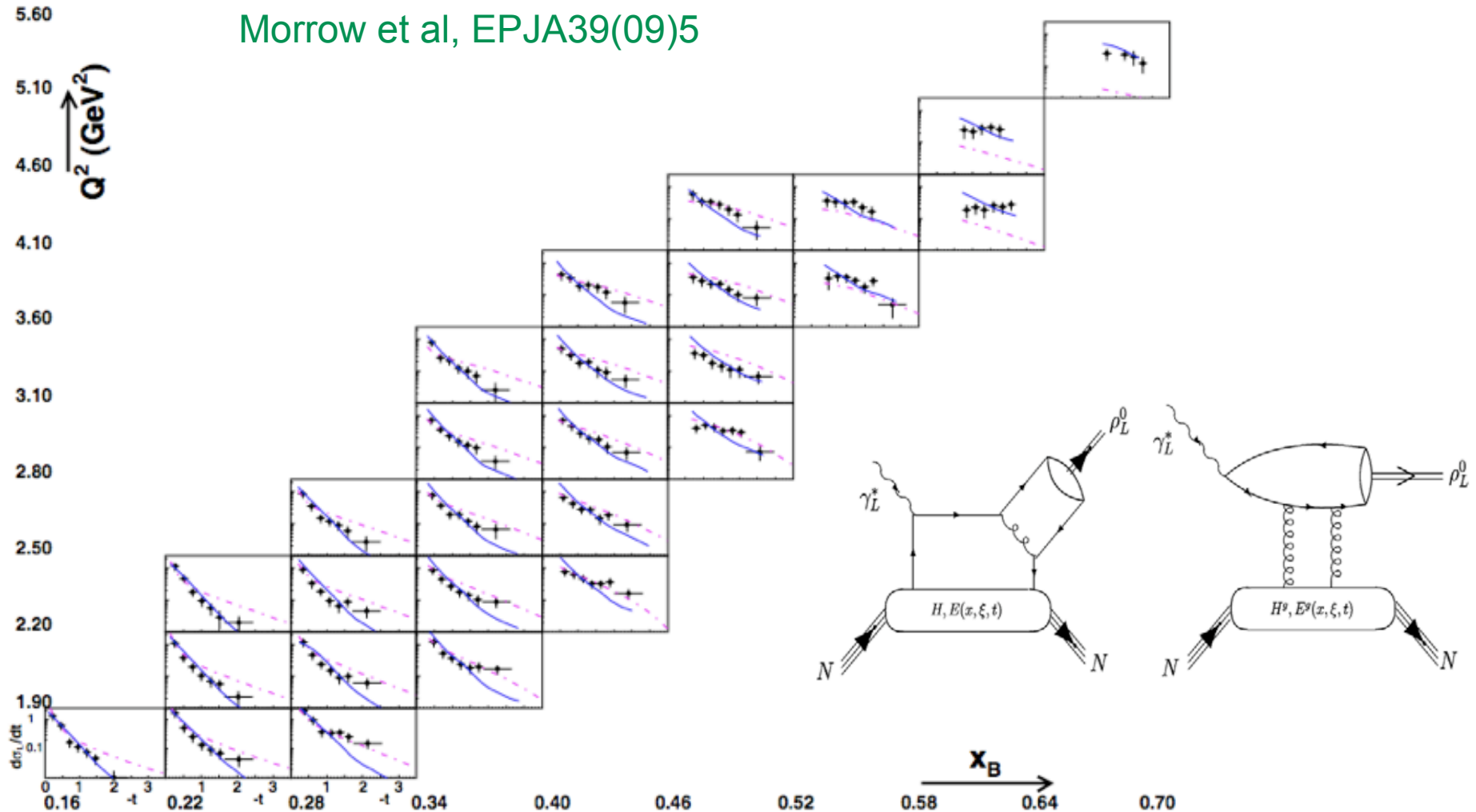


Fig. 26. Longitudinal cross-section $d\sigma_L/dt$ (in $\mu\text{b}/\text{GeV}^2$) for all bins in (Q^2, x_B) as a function of t (in GeV^2). The thick solid curve represents the result of the VGG calculation with the addition of the generalized D -term. The dash-dotted curve is the result of the JML model.



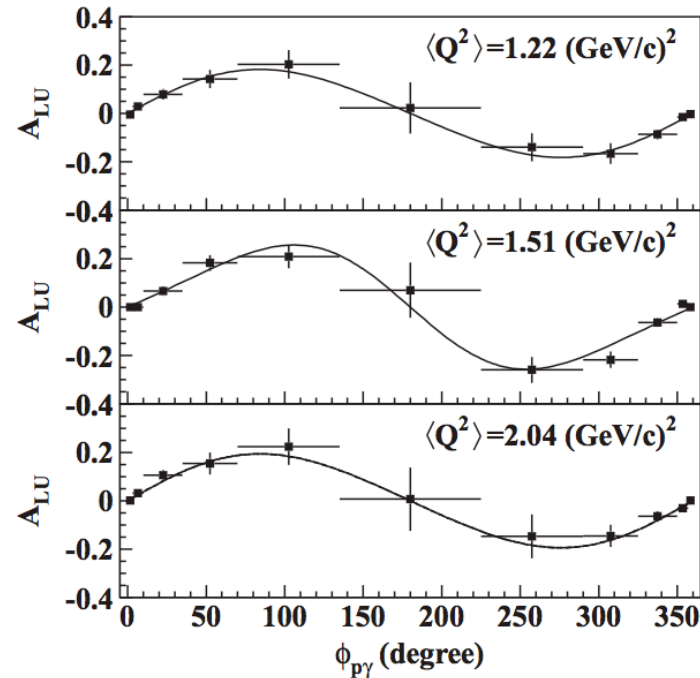
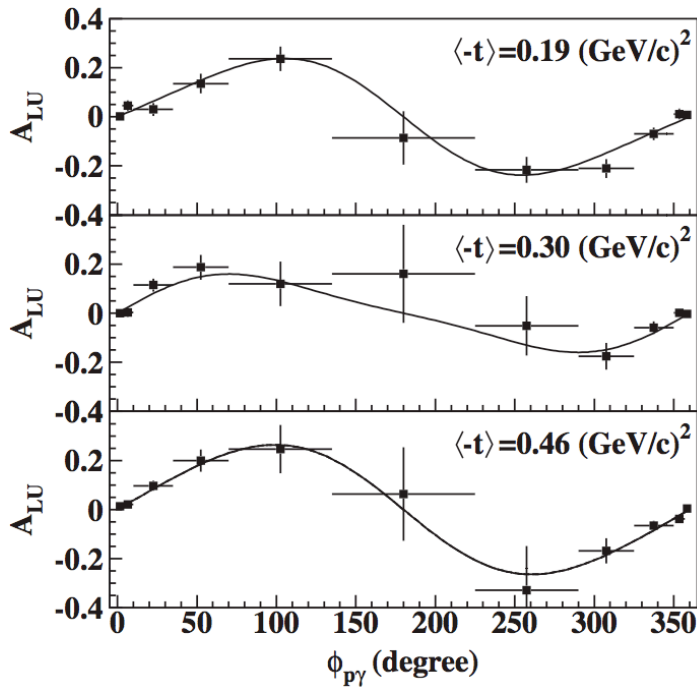
G. GAVALIAN *et al.*

$$A = \alpha \sin\phi + \beta \sin 2\phi$$

PHYSICAL REVIEW C **80**, 035206 (2009)

TABLE III. Results from the fits to the ϕ dependences of A_{LU} with the functions presented in Eqs. (19) and (22). Only statistical uncertainties are presented.

$\langle Q^2 \rangle$ [(GeV/c) ²]	$\langle x_B \rangle$	$\langle -t \rangle$ [(GeV/c) ²]	α	β	α'	γ
1.22	0.17	0.23	0.181 ± 0.032	0.099 ± 0.023	0.181 ± 0.032	-0.098 ± 0.228
1.51	0.20	0.26	0.245 ± 0.028	-0.040 ± 0.021	0.234 ± 0.024	0.319 ± 0.195
2.04	0.28	0.38	0.192 ± 0.044	0.010 ± 0.030	0.191 ± 0.045	-0.107 ± 0.288

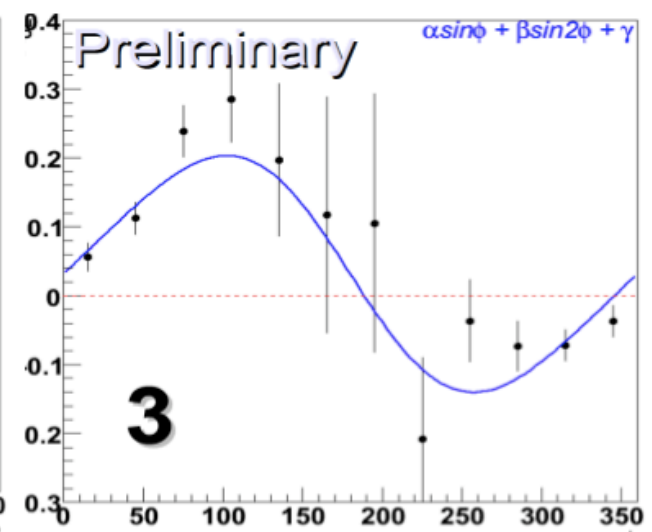
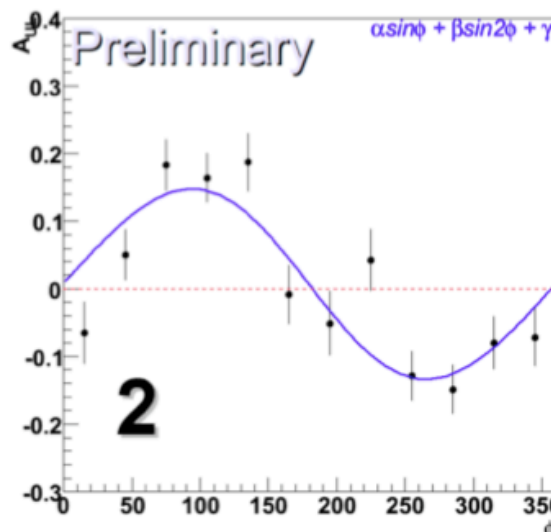
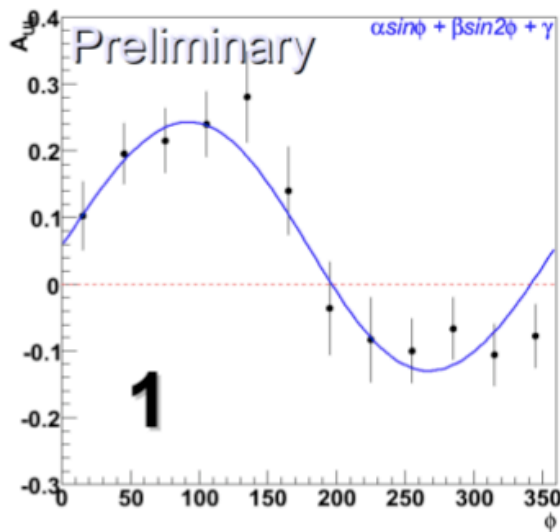




★WG6PSHP1: Andrey Kim *Studies of exclusive processes with a longitudinally polarized target*

$$A = \alpha \sin\phi + \beta \sin 2\phi$$

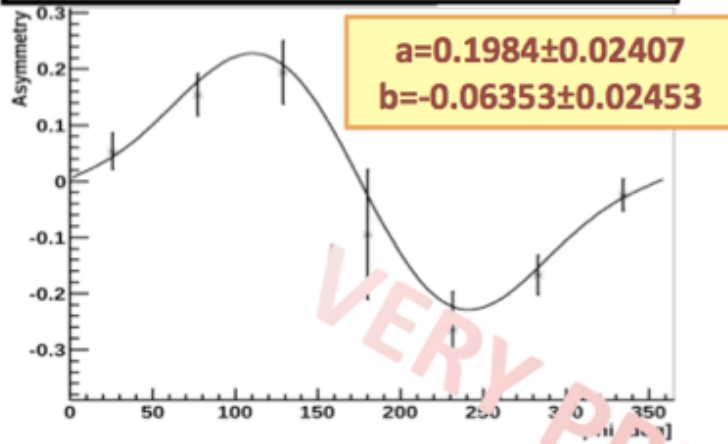
CLAS
eg1-dvcs data





★WG6PSHP1: Andrey Kim *Studies of exclusive processes with a longitudinally polarized target*

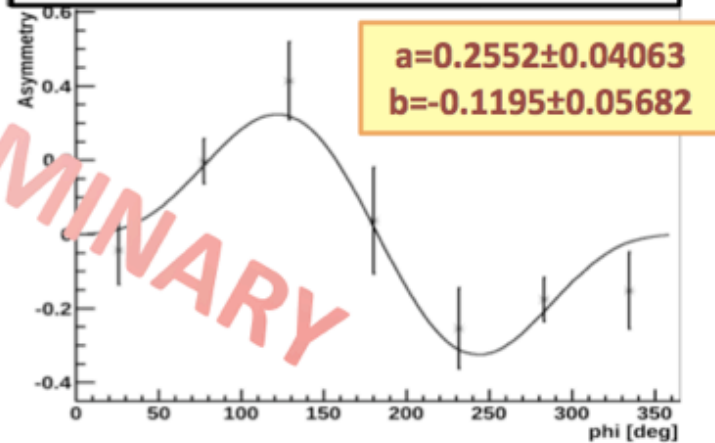
✓ Both photons in IC:



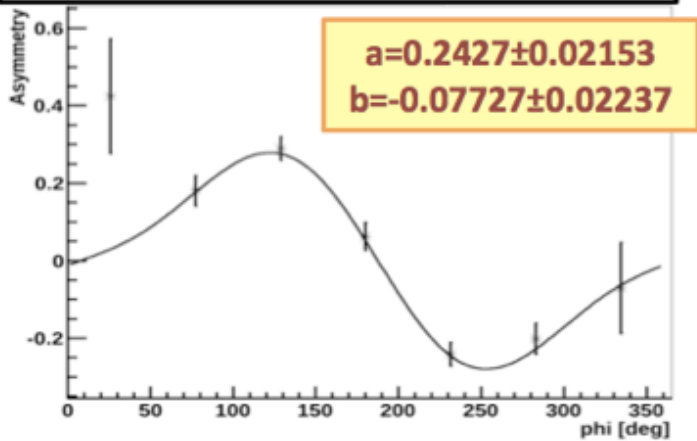
CLAS neutral pions

$$A=a\cdot\sin\phi + b\cdot\sin2\phi$$

✓ One photon in EC, second in IC:



✓ Both photons in EC:



VERY PRELIMINARY



- Jefferson Lab has an intense program on:
 - unpolarized and polarized inclusive DIS
 - semi-inclusive DIS with pions
 - DVCS
 - DVMS with pions and rhos
 - using proton, deuteron, ^3He , and nuclear targets
- Details of these topics can be found:
 - in the advertised talks at DIS2011