

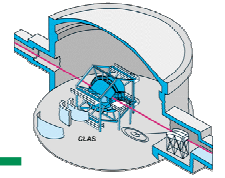
Spin Structure Functions at CLAS

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for the CLAS Collaboration

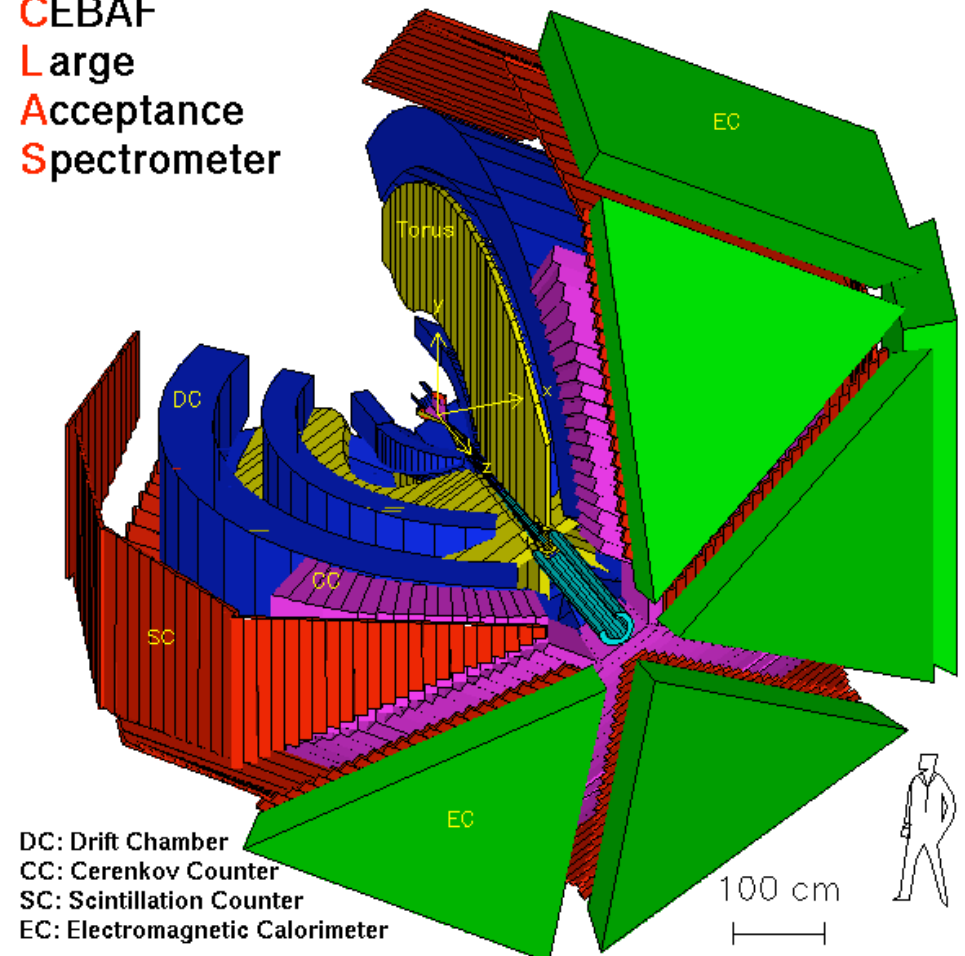
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Presented by J. Pierce
University of Virginia



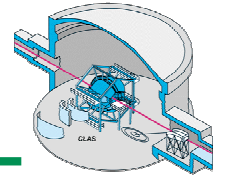
- Long-standing program in Hall-B at JLab to measure longitudinal double spin asymmetries $A_{||}$ on $^{15}\text{NH}_3$ and $^{15}\text{ND}_3$
- EG1: $0.05 < Q^2 < 3.5 \text{ GeV}^2$
– data (2001); anal (2007)
- EG4: $0.01 < Q^2 < 1 \text{ GeV}^2$
– data (2006); anal (2008)
- EG12: $0.5 < Q^2 < 7 \text{ GeV}^2$
– data (2012?); anal (2014)

CEBAF
Large
Acceptance
Spectrometer

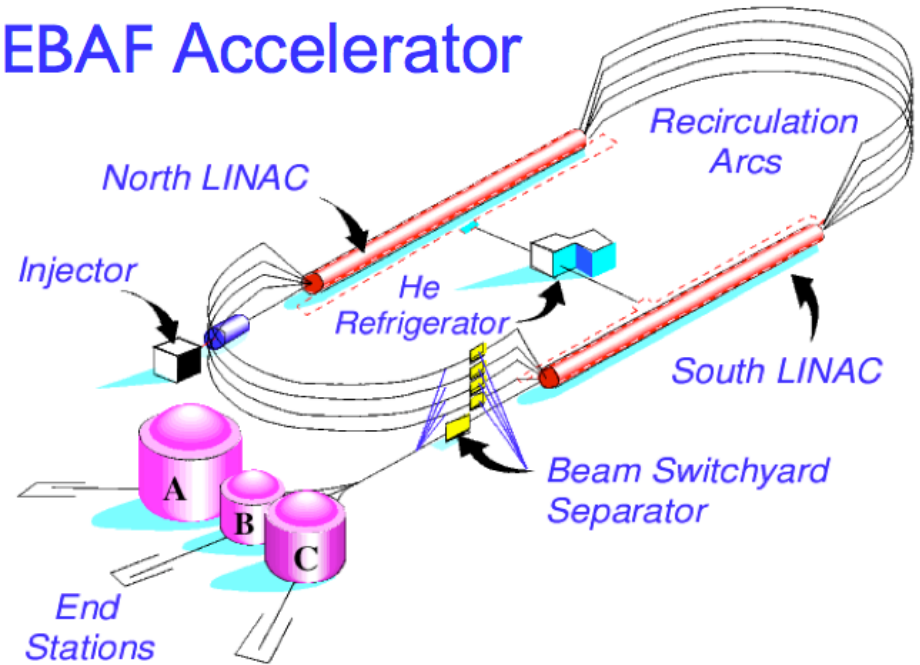




CEBAF Accelerator



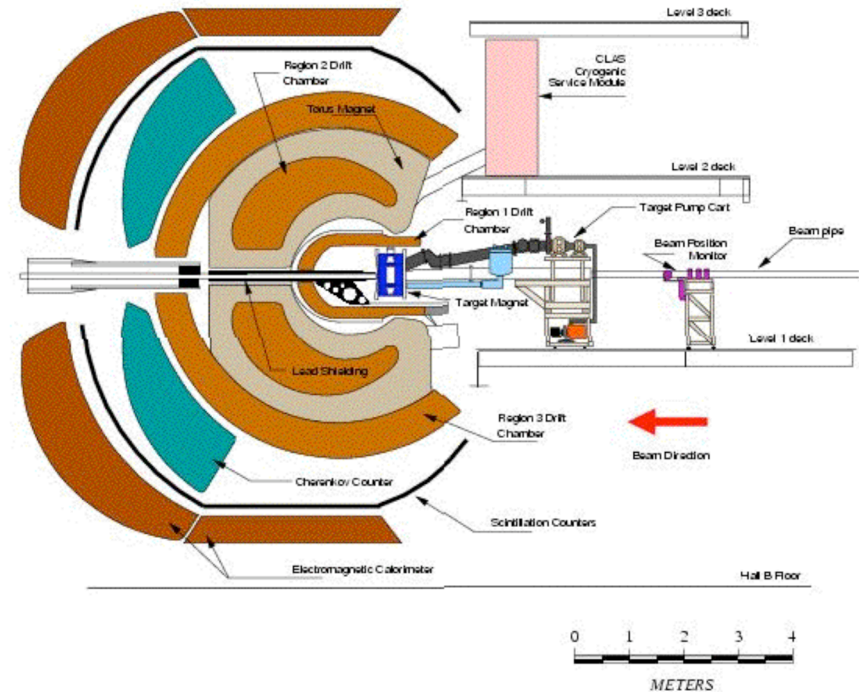
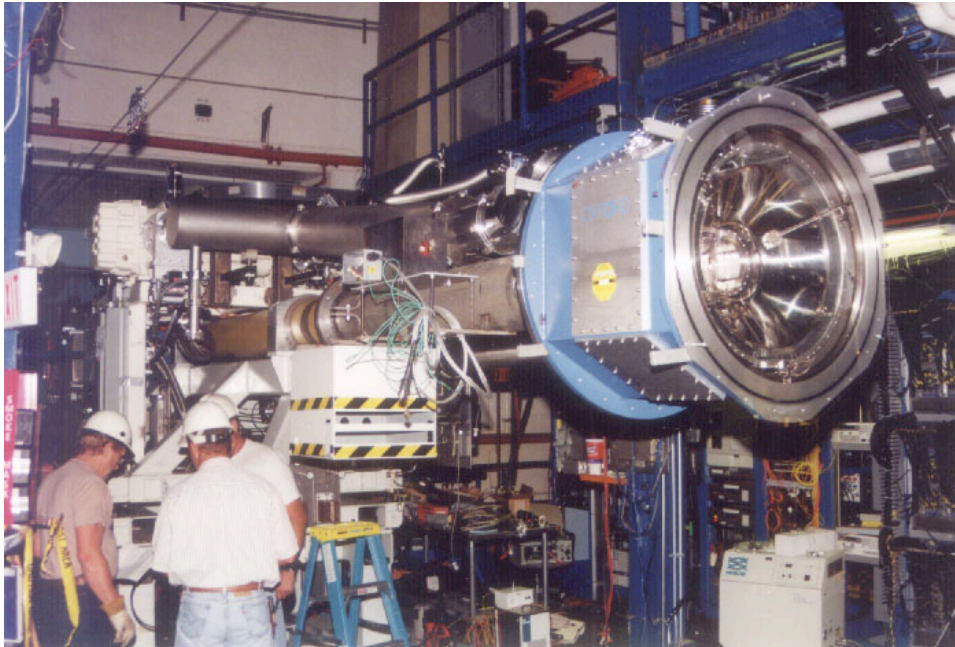
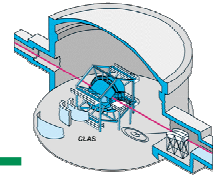
The CEBAF Accelerator



- Electron beams up to 5.7 GeV with >80% longitudinal polarization
- Beam currents of 1-50 nA in Hall B



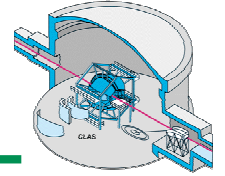
Polarized Target



- Dynamic nuclear polarization of NH_3 and ND_3
- Polarizations of 70-80% for p and 20-30% for d
- Luminosity $10^{35} \text{ cm}^{-2}\text{s}^{-1}$



Formalism



$$A_{\parallel} = \frac{\sigma^{\downarrow\uparrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\downarrow\uparrow} + \sigma^{\uparrow\uparrow}}$$

$$A_{\parallel} = D(A_1 + \eta A_2)$$

We can extract A_1 using a model for A_2 (small), or g_1 using a model for g_2 (small)

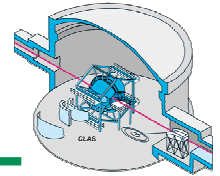
We can extract A_1 and A_2 from A_{\parallel} at multiple values of $\eta(E_{\text{beam}})$

$$A_1 = \frac{\sigma_{1/2}^T - \sigma_{3/2}^T}{\sigma_{1/2}^T + \sigma_{3/2}^T}$$

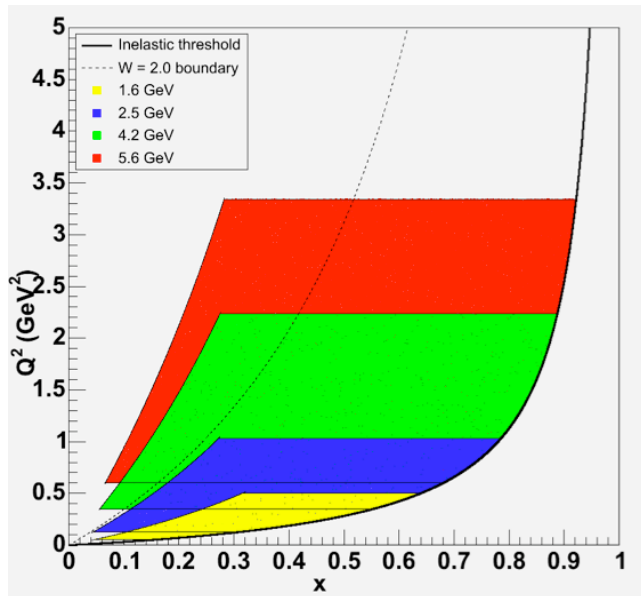
$$= \frac{g_1(x, Q^2) - \gamma^2 g_2(x, Q^2)}{F_1(x, Q^2)}$$

$$A_2 = \frac{2\sigma_{LT}}{\sigma_{1/2}^T + \sigma_{3/2}^T}$$

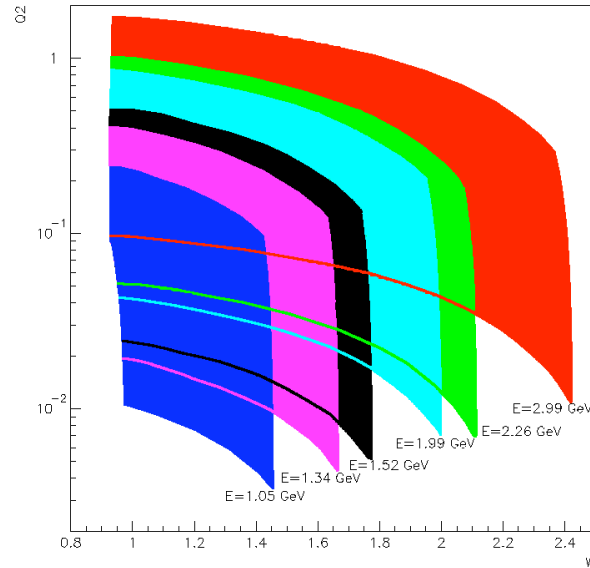
$$= \frac{\gamma[g_1(x, Q^2) + g_2(x, Q^2)]}{F_1(x, Q^2)}$$



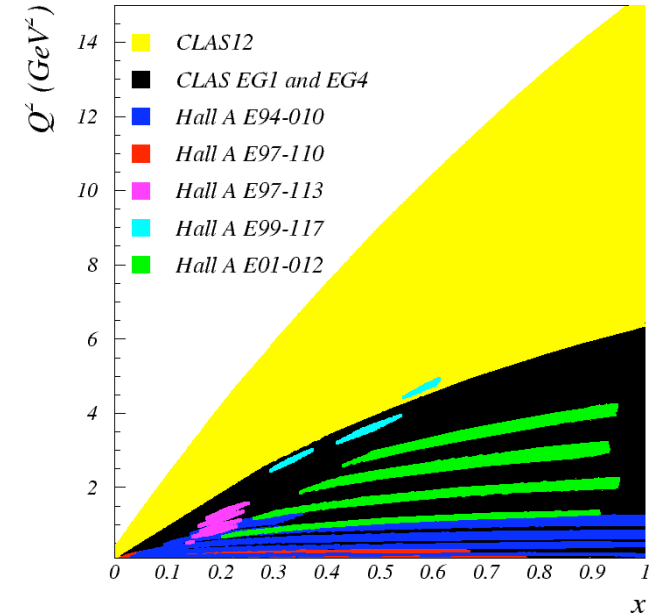
EG1



EG4



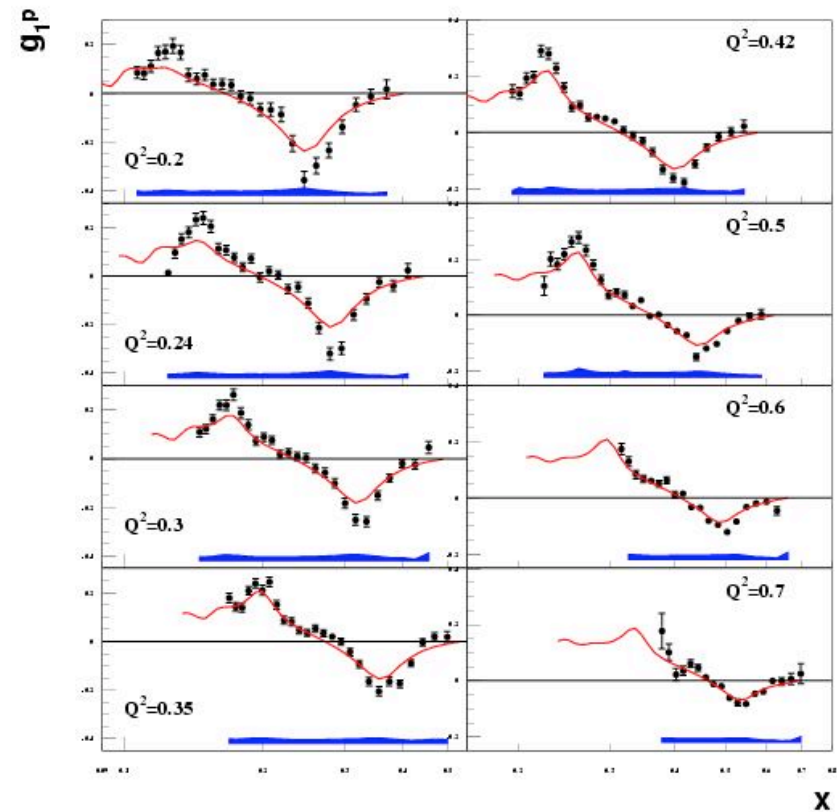
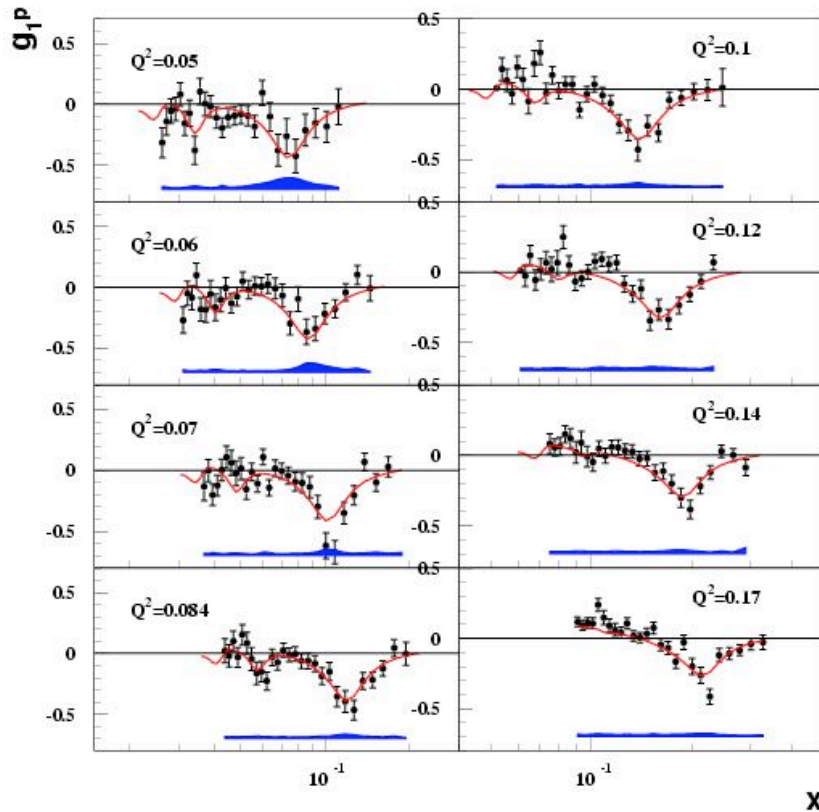
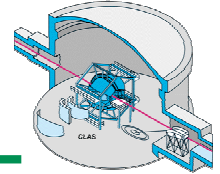
EG12



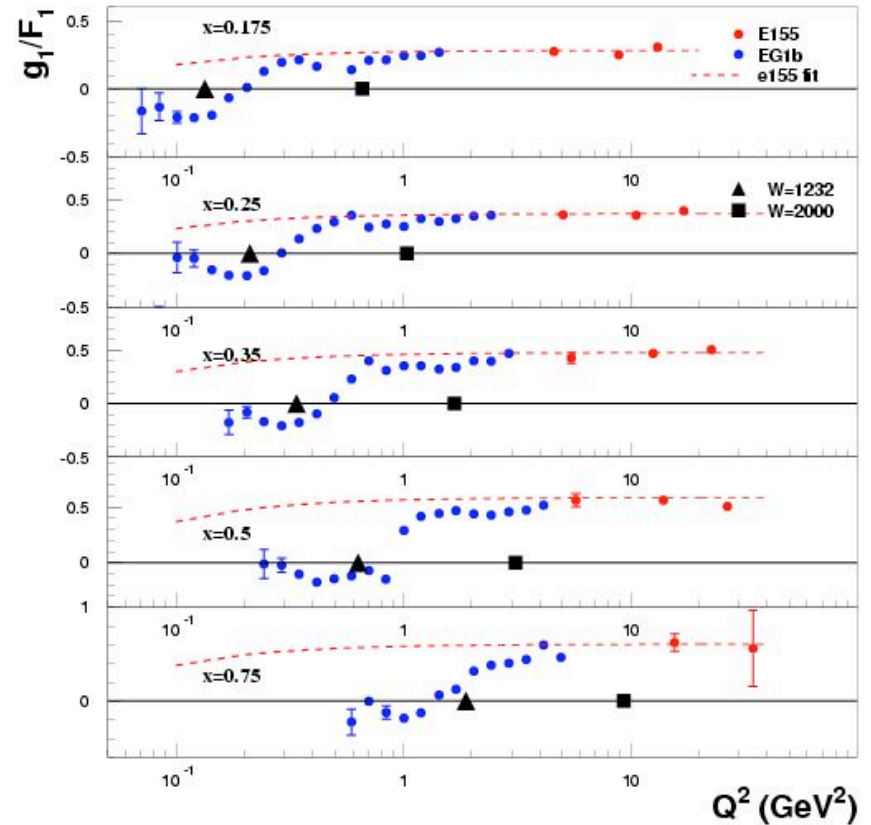
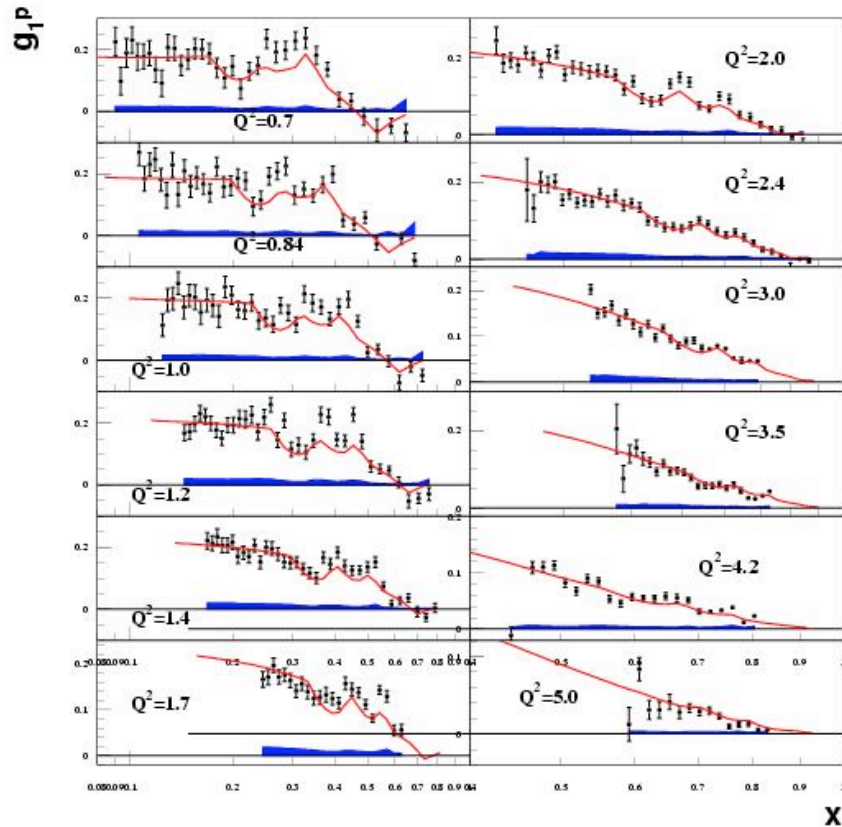
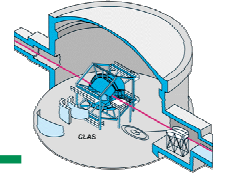
- Overlapping colors correspond to different beam energies
- CLAS measures a large range in x at each fixed Q^2
- Different E_{beam} for fixed (x, Q^2) allows separation of A_1 & A_2



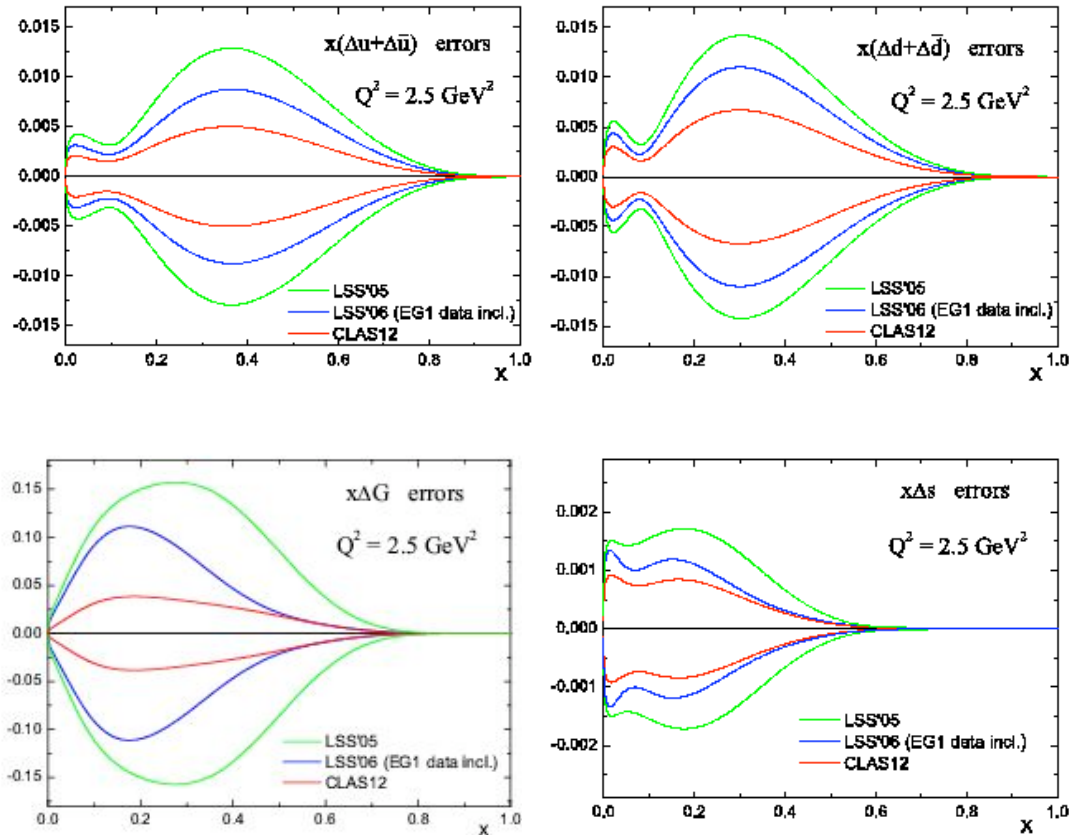
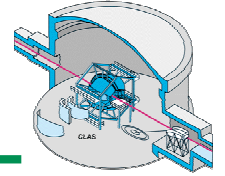
EG1 $g_1^p(Q^2 < 0.7)$



- At low Q^2 the Δ resonance drives g_1 negative
- Extensive x -range at fixed Q^2 allows integration over x
- Red curve is the EG1 model used for radiative corrections



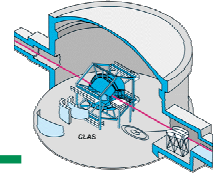
- At higher Q^2 , g_1 becomes positive everywhere
- g_1/F_1 falls far below the DIS extrapolation at low Q^2
- Red curve is the EG1 model (dashed: DIS extrapolation)



- Error envelopes for PDFs from LSS05 global analysis (green)
- CLAS EG1 data significantly improve errors on Δu , Δd , Δx and ΔG (blue)
- CLAS EG12 (12 GeV upgrade) will especially improve ΔG (red)

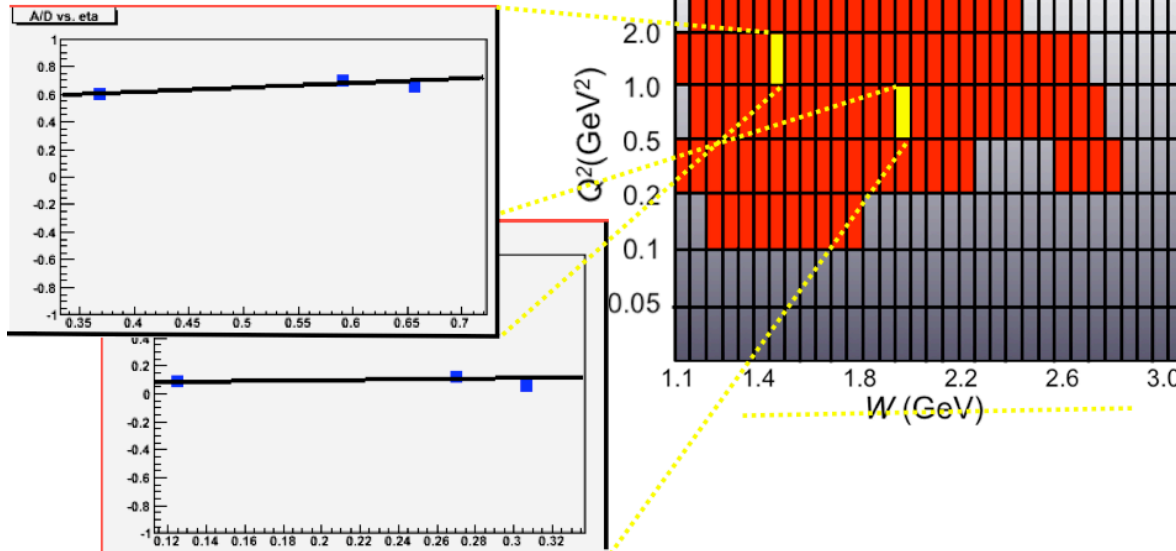


EG1 Extraction of A_2

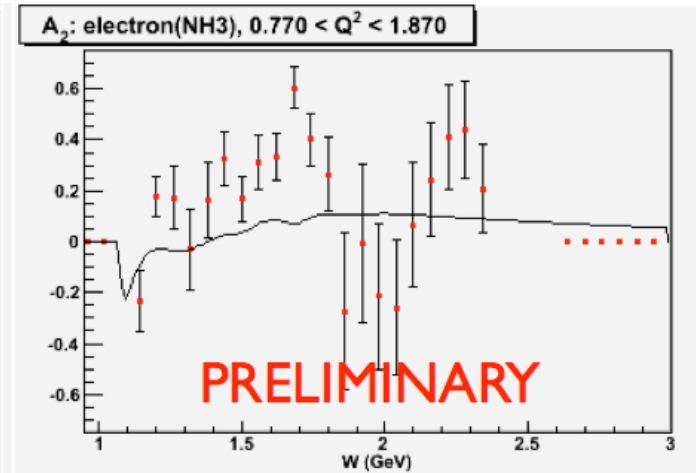
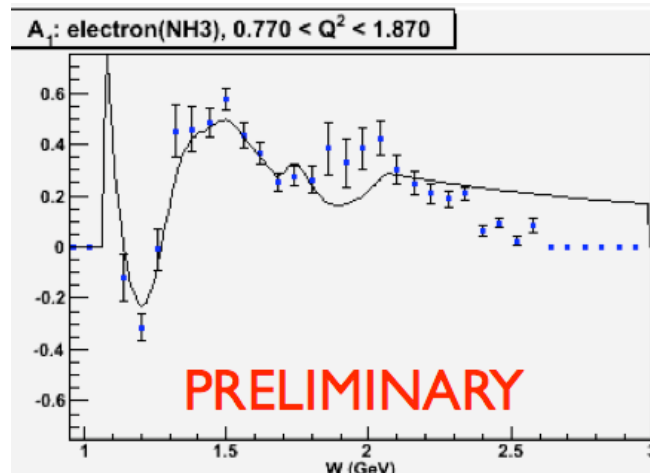


$$A_{\parallel} = D(A_1 + \eta A_2)$$

$$\eta = \frac{\epsilon \sqrt{Q^2/E}}{1 - \epsilon E'/E} \quad D = \frac{1 - \epsilon E'/E}{1 + \epsilon R}$$

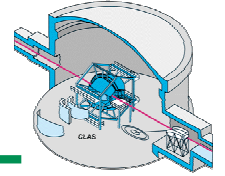


- Analysis is in progress to obtain both A_1 and A_2 from the EG1 data
- Intercept gives A_1
- Slope gives A_2
- A_2 is larger than EG1 model (MAID, AO) as is Hall C RSS experiment



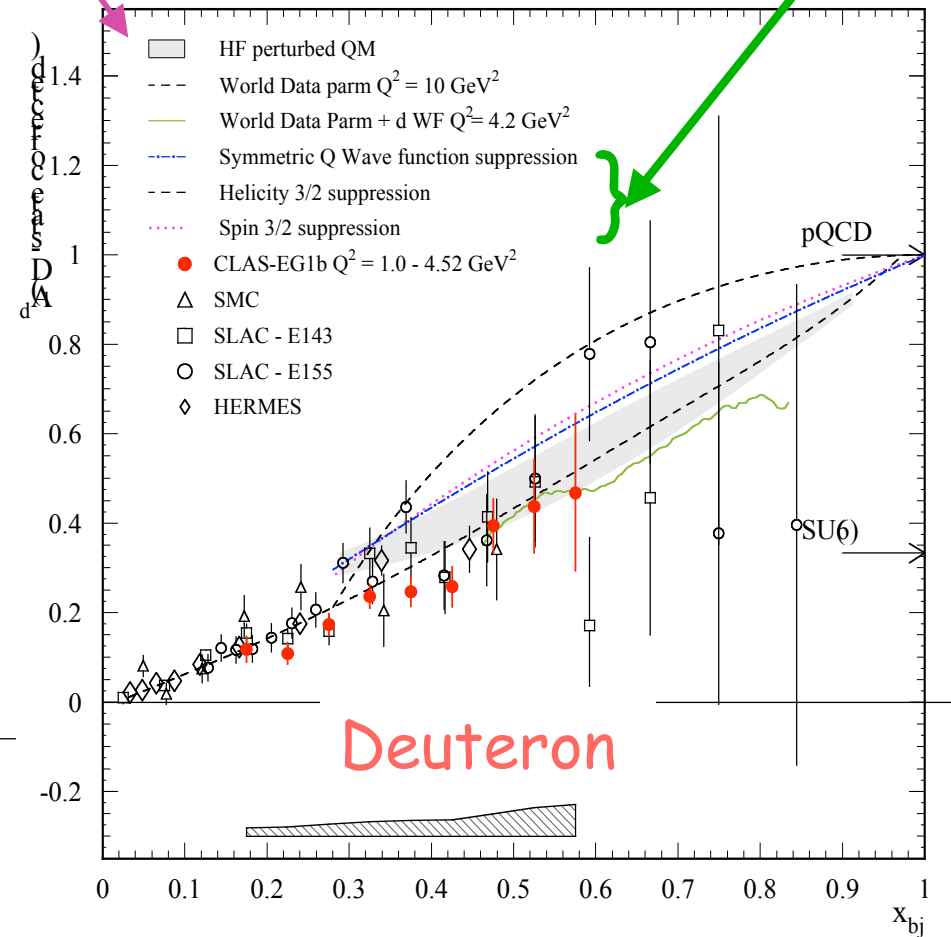
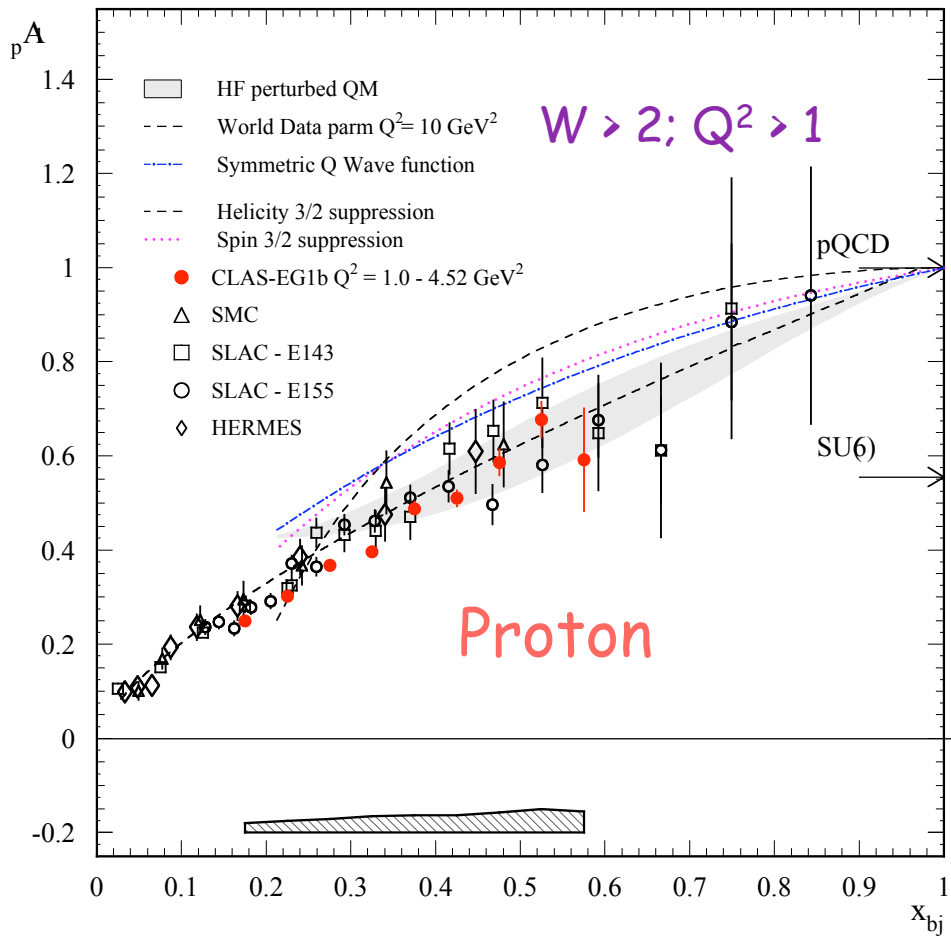


A₁ Data from EG1



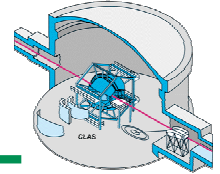
Isgur, PRD 59, 034013 (2003)

Close and Melnitchouk, PRC 68, 035210 (2003)

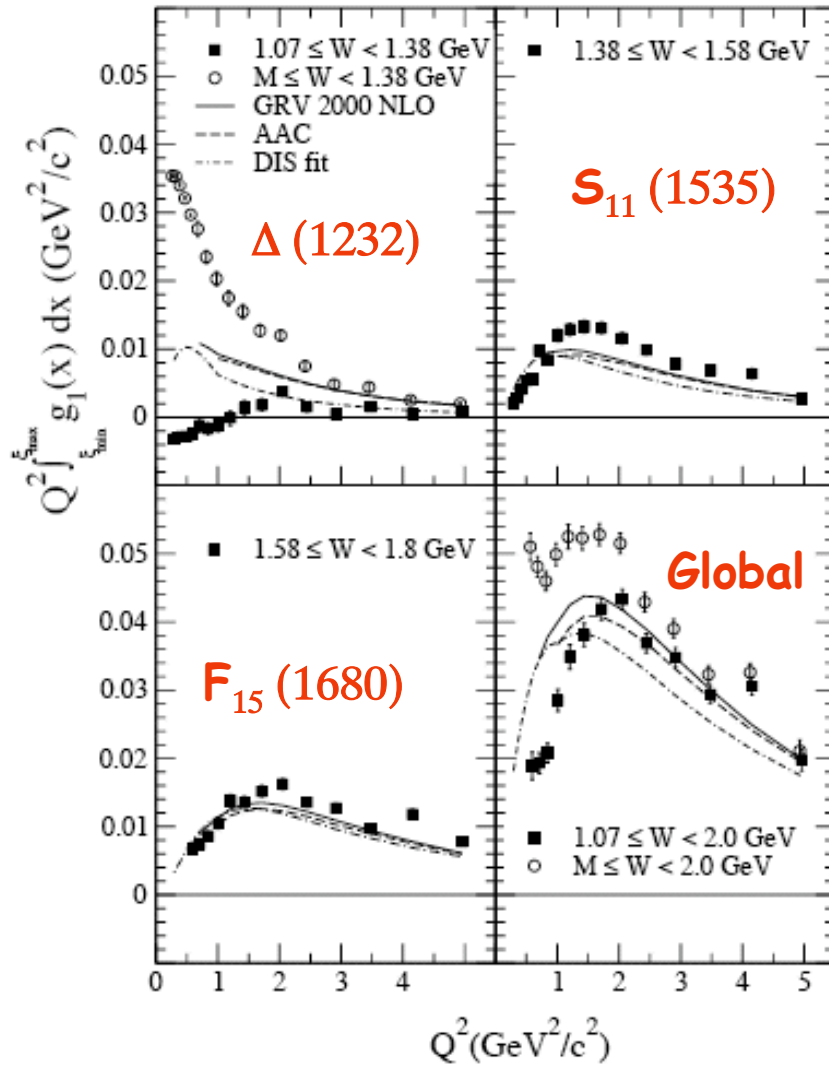




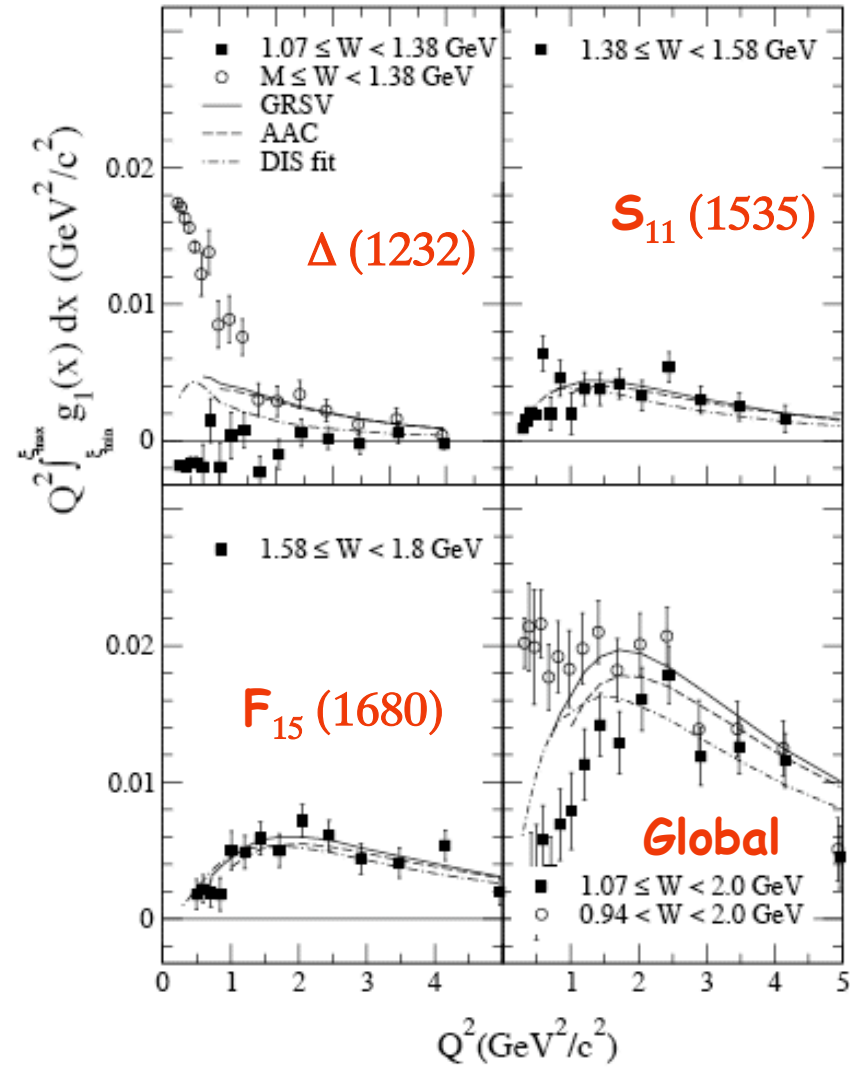
Duality



Proton

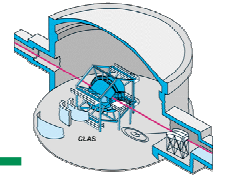


Deuteron

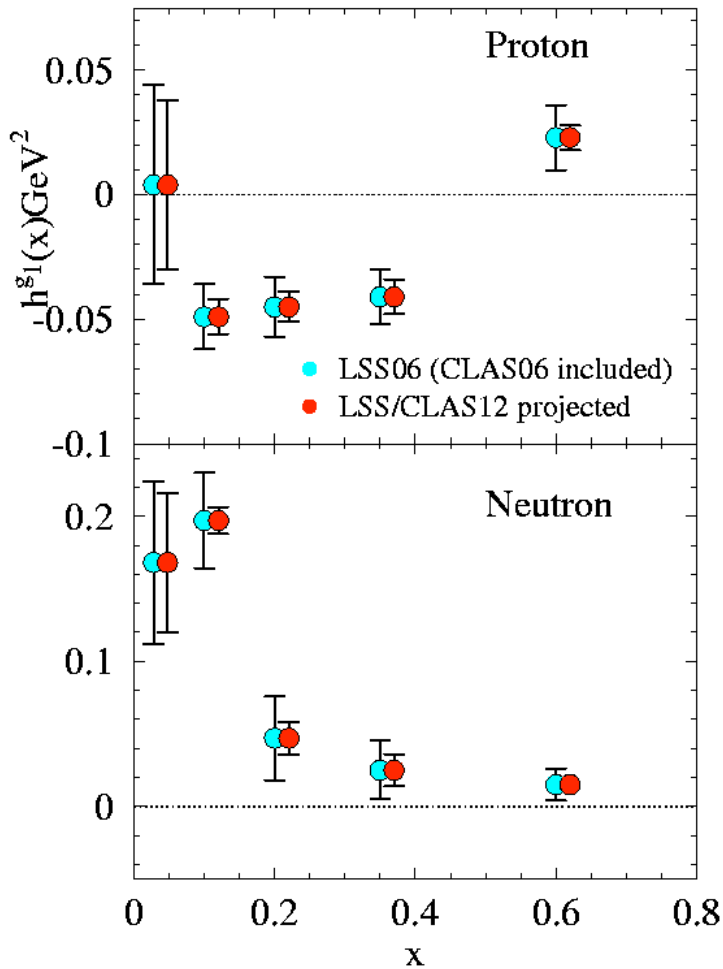




Higher Twist from g_1



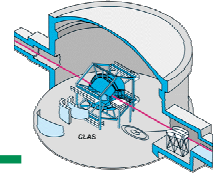
$$\left[\frac{g_1(x, Q^2)}{F_1(x, Q^2)} \right]_{\text{exp}} F_1(x, Q^2)_{\text{exp}} = g_1(x, Q^2)_{\text{exp}} = g_1(x, Q^2)_{LT} + h^{g_1}(x)/Q^2$$



$$1 < Q^2 < 5 \text{ GeV}^2, \quad 2 < W < 3.5 \text{ GeV}$$

$$\int_0^1 dx h^{g_1}(x) = \frac{4}{9} M^2 (d_2 + f_2)$$

- F_1 from NMC fit to F_2 and 1998 SLAC fit to R
- g_1 (leading twist) from NLO fit at high Q^2
- h from fit to all data, especially CLAS in the pre-asymptotic region
- d_2 : twist-3, f_2 : twist-4



$$\Gamma_1^{(n)} = \int_0^1 x^n g_1(x, Q^2) dx = \frac{a_n}{2}, \quad n=0,2,4,\dots,$$

$$\Gamma_2^{(n)} = \int_0^1 x^n g_2(x, Q^2) dx = \frac{1}{2} \frac{n}{n+1} (d_n - a_n), \quad n=2,4,\dots,$$

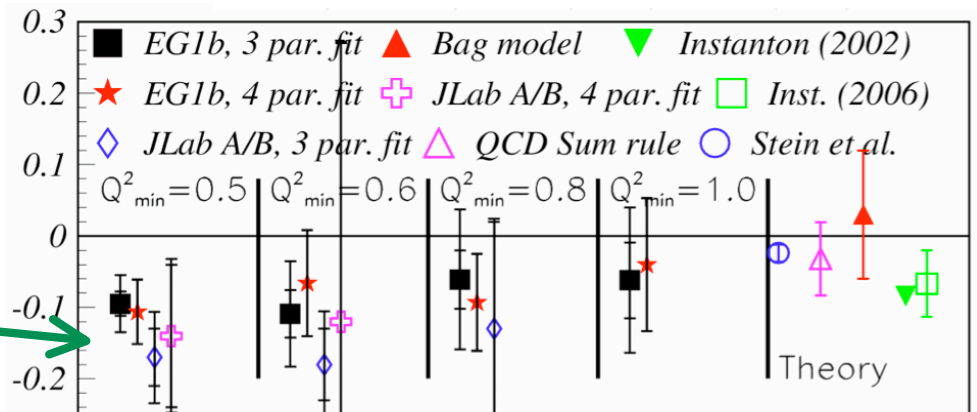
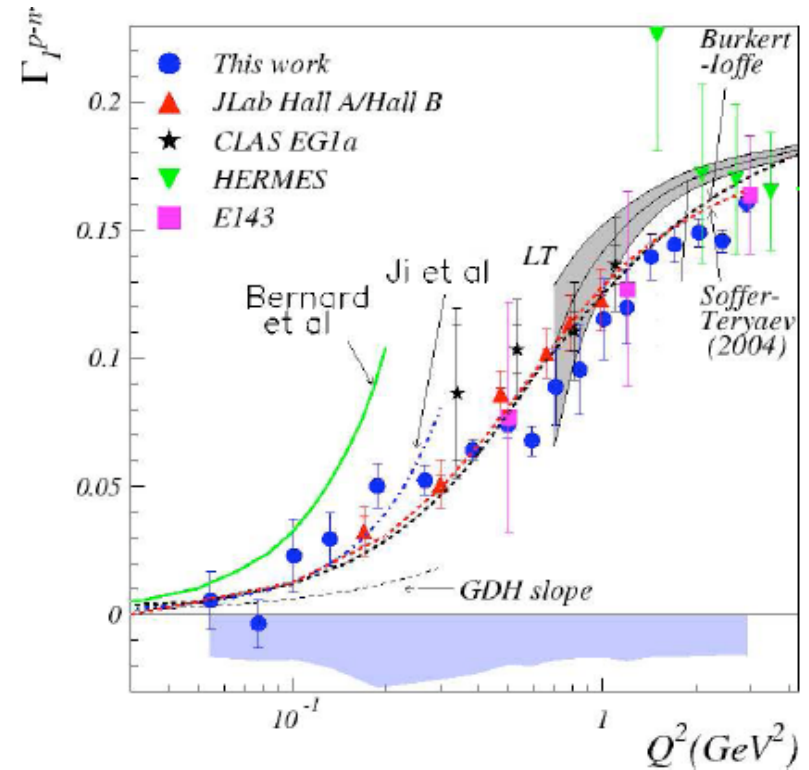
Bjorken Sum Rule:

$$\Gamma_1^{p-n} = \frac{g_A}{6} \left[1 - \frac{\alpha_s}{\pi} - 3.58 \left(\frac{\alpha_s}{\pi} \right)^2 - 20.21 \left(\frac{\alpha_s}{\pi} \right)^3 \right] + \frac{\mu_4^{p-n}}{Q^2} + \dots$$

$$\mu_4^{p-n} = \frac{M^2}{9} (a_2^{p-n} + 4d_2^{p-n} + 4f_2^{p-n})$$

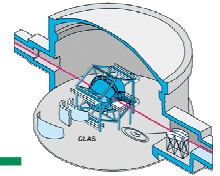
$$d_2^{p-n} = \int_0^1 dx x^2 (2g_1^{p-n} + 3g_2^{p-n})$$

Fit Γ_1^{p-n} to powers of $1/Q^2$ and extract f_2^{p-n}





Moments $\Gamma_1^{p,d}$



$$\Gamma_1^{p,d}(Q^2) = \int_0^1 g_1^{p,d}(x, Q^2) dx \quad \Gamma_1(Q^2) = aQ^2 + bQ^4 + cQ^6 + dQ^8$$

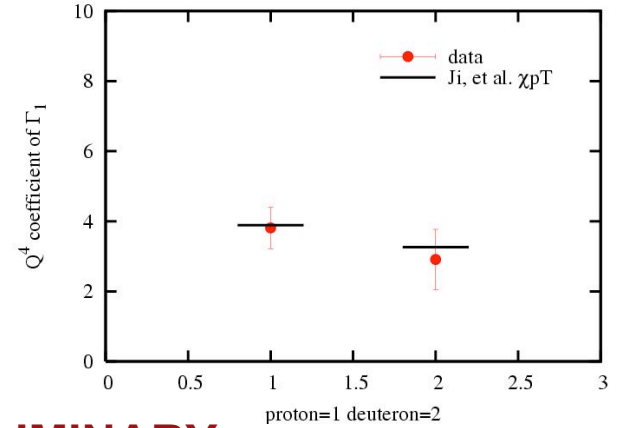
$$\Gamma_1^d(Q^2) = (1 - 1.5\omega_D) \{ \Gamma_1^p(Q^2) + \Gamma_1^n(Q^2) \}$$

$$\Gamma_1^p(Q^2) = -\frac{\kappa_p^z}{8M^2} Q^2 + 3.89Q^4 + \dots$$

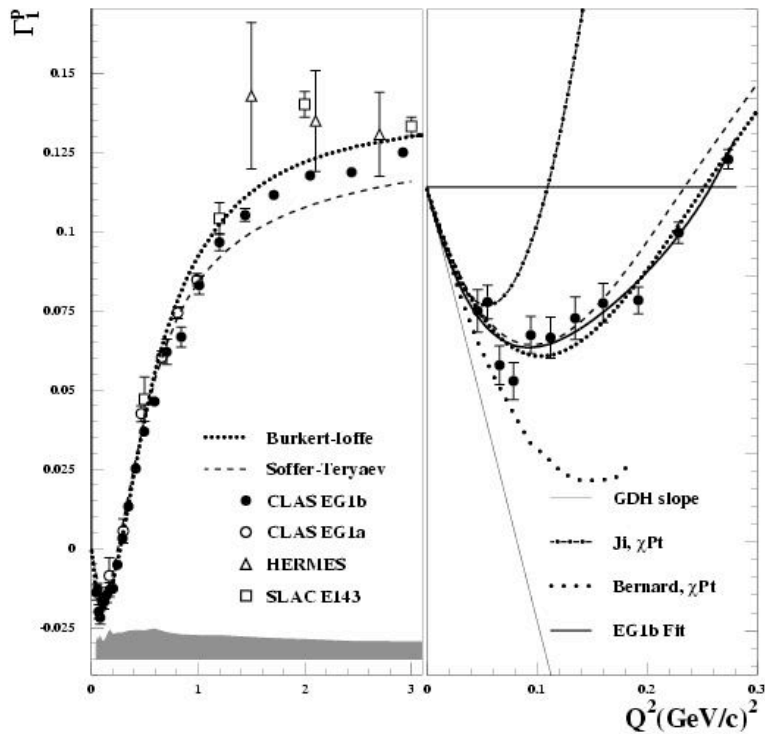
$$\Gamma_1^n(Q^2) = -\frac{\kappa_n^z}{8M^2} Q^2 + 3.15Q^4 + \dots$$

low Q^2 fit

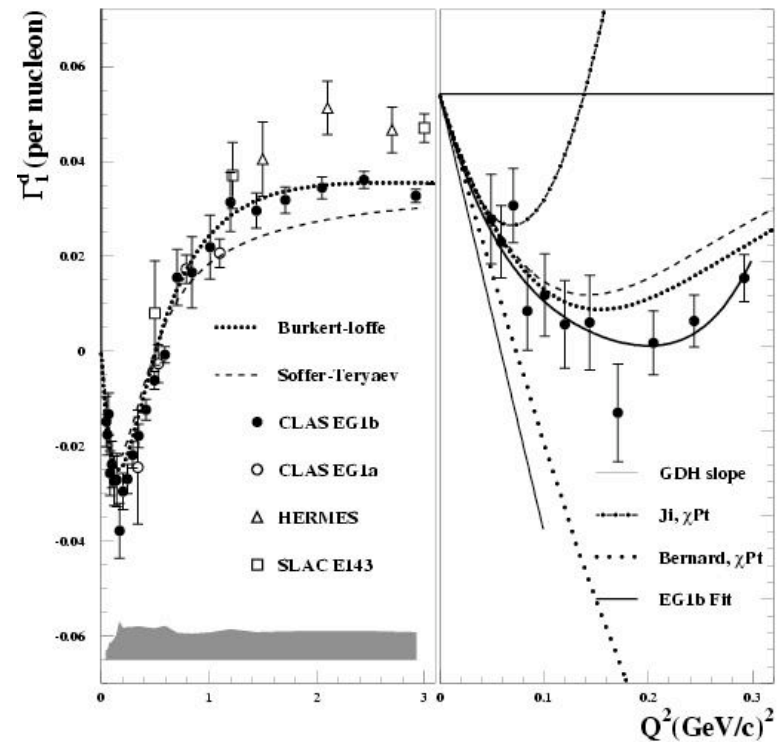
GDH + χpT

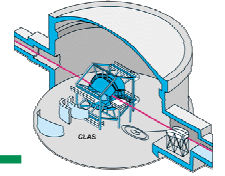


PRELIMINARY

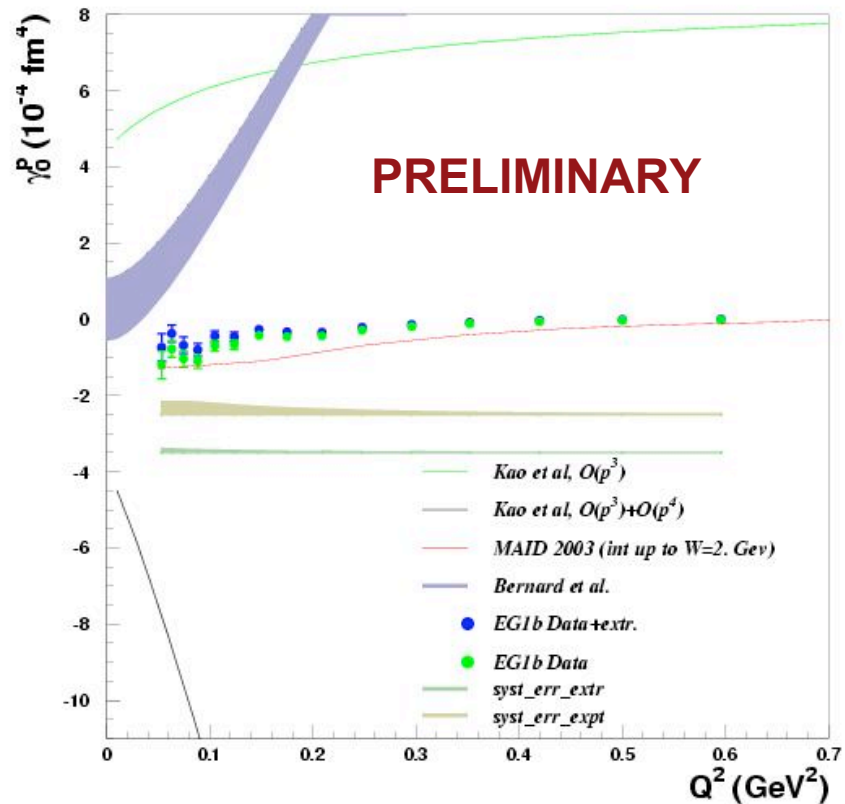
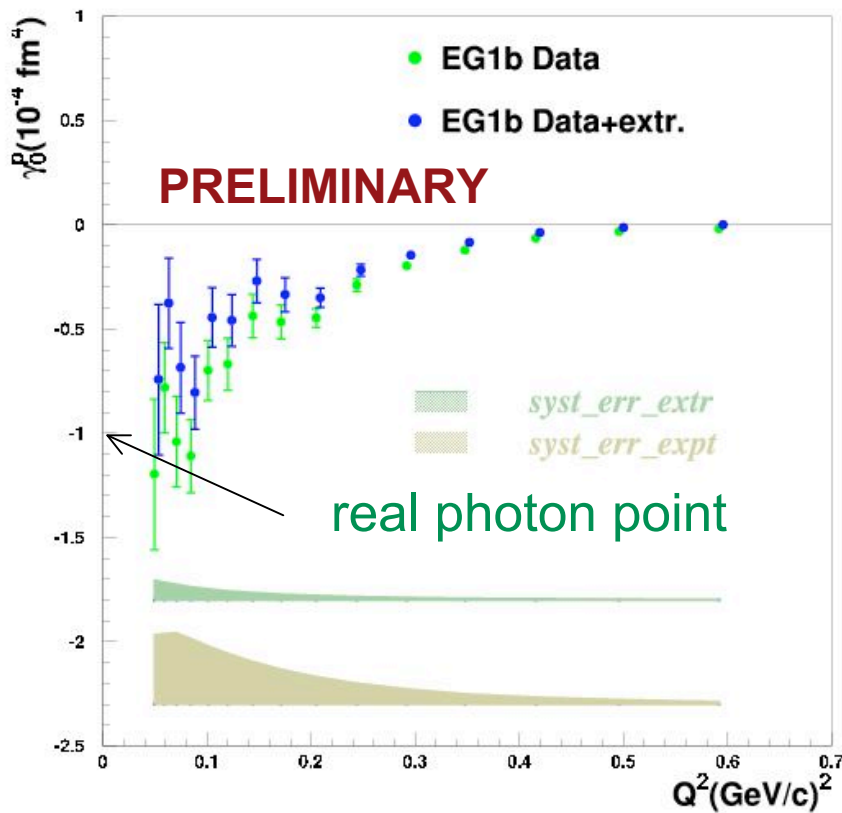


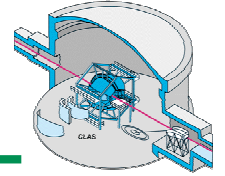
PRELIMINARY



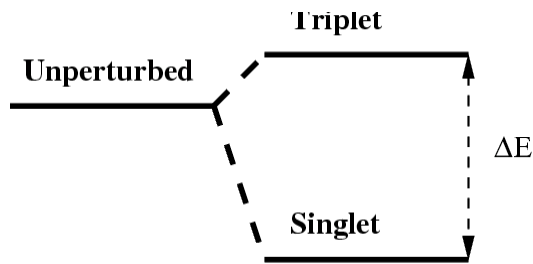


$$\gamma_0(Q^2) = \frac{4e^2 M^2}{\pi Q^6} \int_0^{x_0} dx x^2 \{g_1(x, Q^2) - \gamma^2 g_2(x, Q^2)\}$$





$$E_{\text{HFS}}(e^- p) = 1.4204057517667(9) \text{ GHz} = (1 + \Delta_{\text{QED}} + \Delta_R^p + \Delta_S) E_F^p$$



$$\Delta_S = \Delta_Z + \Delta_{\text{pol}} \quad \delta_Z^{\text{rad}} = \frac{\alpha}{3\pi} \left[2 \ln \frac{\Lambda^2}{m^2} - \frac{4111}{420} \right]$$

$$\text{Zemach: } \Delta_Z = -2\alpha m_e \langle r \rangle_Z (1 + \delta_Z^{\text{rad}})$$

$$\langle r \rangle_Z = -\frac{4}{\pi} \int_0^\infty \frac{dQ}{Q^2} \left[G_E(Q^2) \frac{G_M(Q^2)}{1+\kappa} - 1 \right]$$

$$\Delta_S = -38.62(16) \text{ ppm} \quad \Delta_Z = -41.0(5) \text{ ppm} \quad \Delta_{\text{pol}} = 2.38(58) \text{ ppm}$$

$$\Delta_{\text{pol}} = \frac{\alpha m_e}{2\pi(1+\kappa)M} (\Delta_1 + \Delta_2) = (0.2264798 \text{ ppm}) (\Delta_1 + \Delta_2)$$

$$\Delta_1 = \frac{9}{4} \int_0^\infty \frac{dQ^2}{Q^2} \left\{ F_2^2(Q^2) + \frac{8m_p^2}{Q^2} B_1(Q^2) \right\}$$

$$\Delta_2 = -24m_p^2 \int_0^\infty \frac{dQ^2}{Q^4} B_2(Q^2).$$

$$\tau = \nu^2 / Q^2$$

$$B_1 = \int_0^{x_{\text{th}}} dx \beta(\tau) g_1(x, Q^2),$$

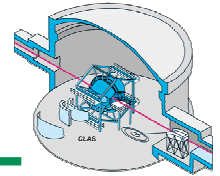
$$B_2 = \int_0^{x_{\text{th}}} dx \beta_2(\tau) g_2(x, Q^2),$$

$$\beta(\tau) = \frac{4}{9} \left(-3\tau + 2\tau^2 + 2(2-\tau)\sqrt{\tau(\tau+1)} \right)$$

$$\beta_2(\tau) = 1 + 2\tau - 2\sqrt{\tau(\tau+1)},$$



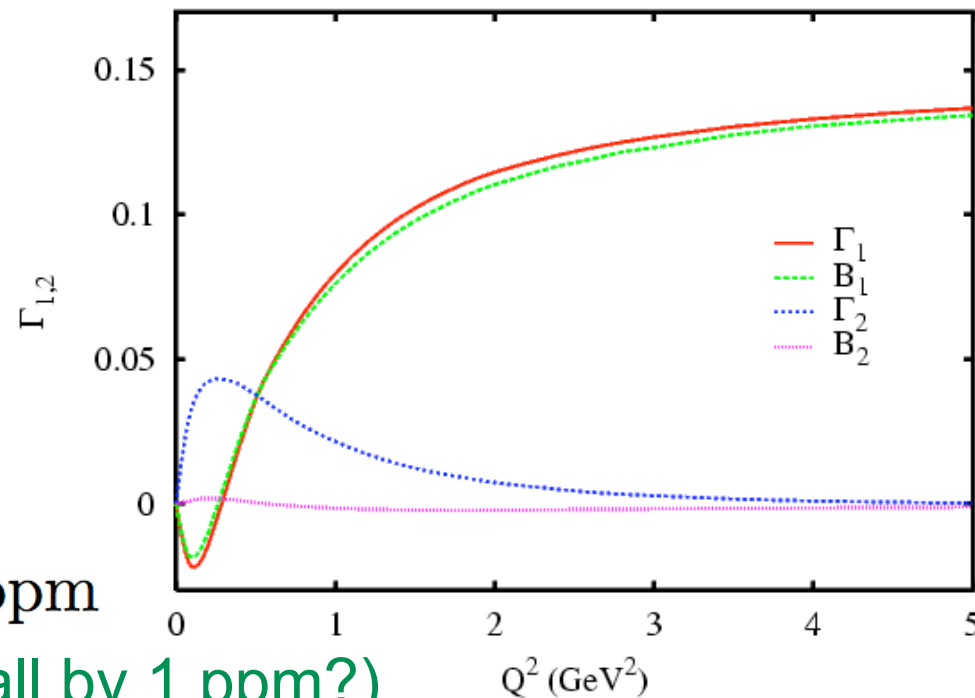
$\Delta_{1,2}$ from $g_{1,2}$



Comparisons between $\Gamma_1 = \int g_1 dx$ and $B_1 = \int \beta_1 g_1 dx$
and between $\Gamma_2 = \int g_2 dx$ and $B_2 = \int \beta_2 g_2 dx$

PRL96,163001

- $B_1 \approx \Gamma_1$
- $B_2 \approx 0$
- Experimentally, errors on Γ_1 are understood; we exploit this fact.
- $\Gamma_2 = \int g_2 dx \neq 0$ at low Q^2 .

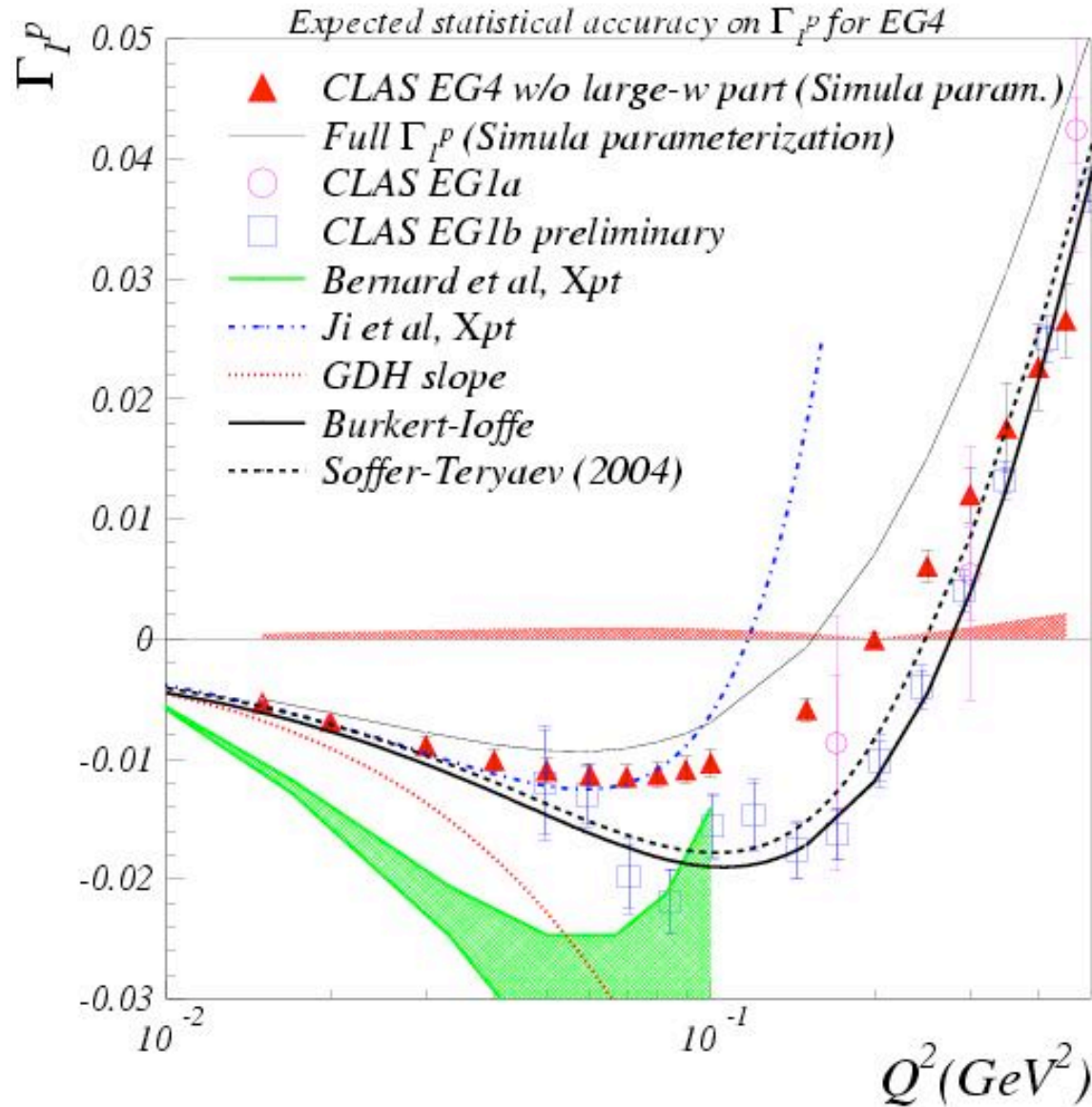
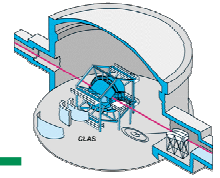


$\Delta_{\text{pol}} = (1.3 \pm 0.3) \text{ ppm}$
(from EG1: too small by 1 ppm?)

Nucleon structure is the largest uncertainty in calculating HFS.
Better g_1, g_2, G_M, G_E data at low Q^2 required to resolve discrepancy.

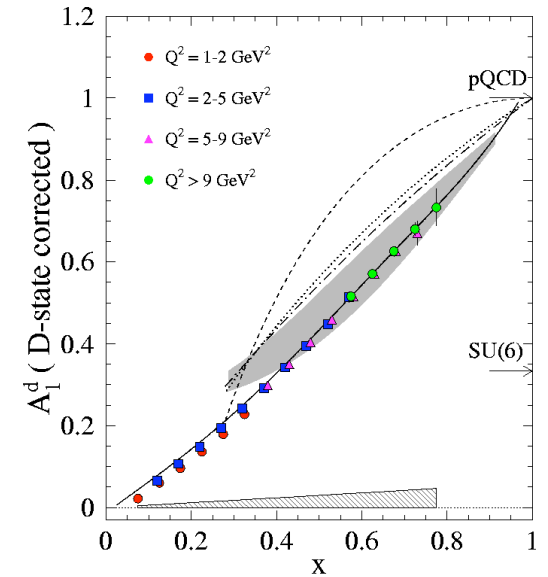
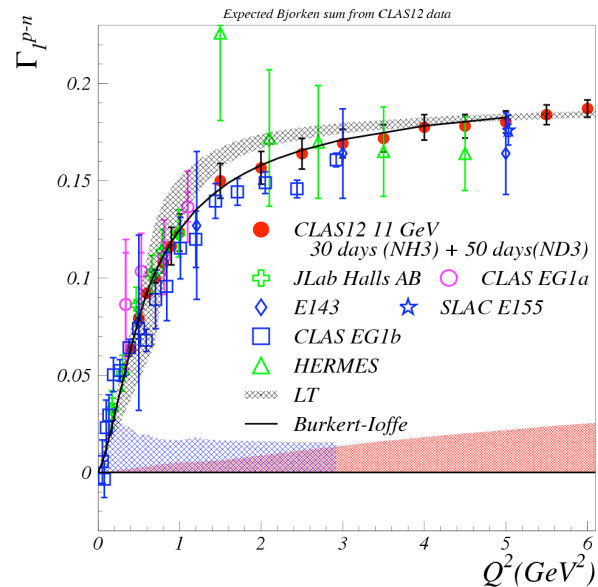
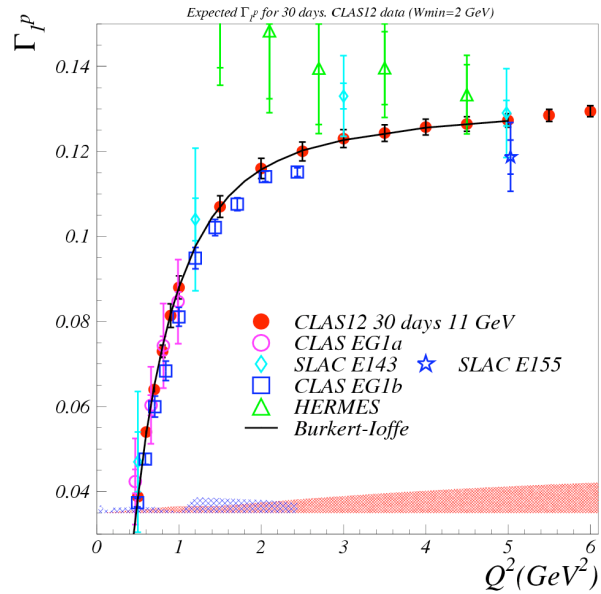
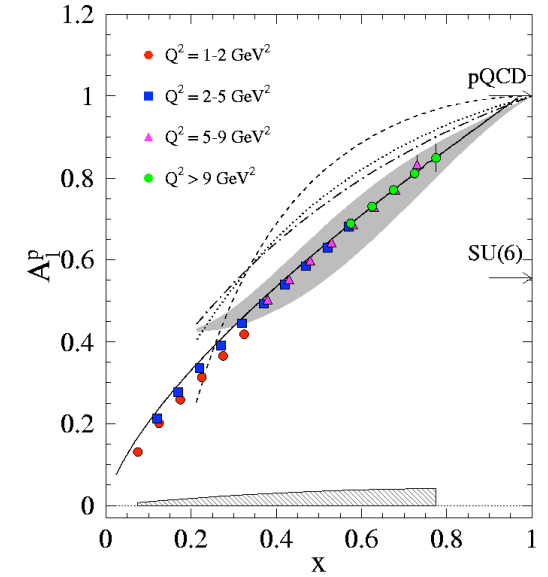
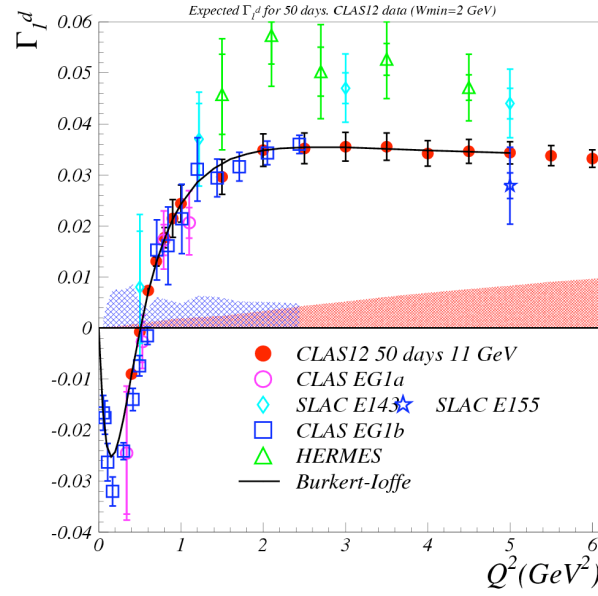
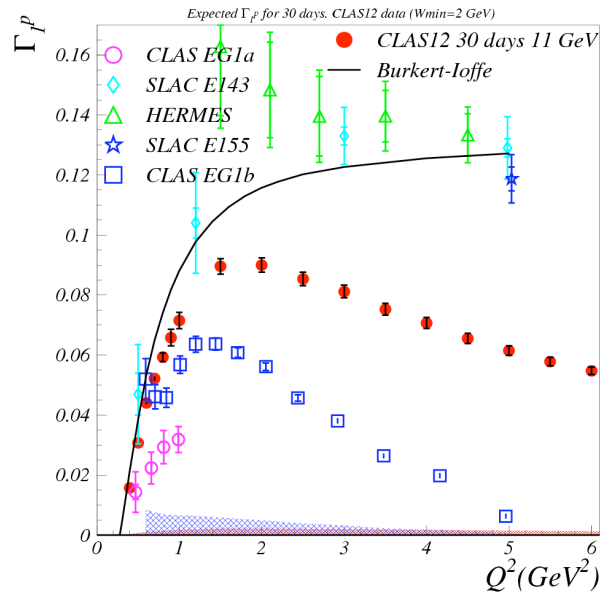
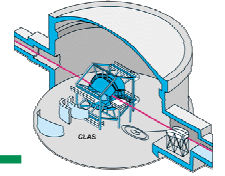


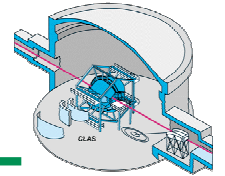
EG4 Expectations





EG12 Expectations





- CLAS, past, present and future, provides high-quality $A_{||}$ data over a large and continuous range in x and Q^2 that
- significantly improve global PDF fits to Δu , Δd , Δs and ΔG
 - precisely determine higher twists
 - rigorously probe duality over a wide Q^2 range
 - quantitatively test χ PT calculations at low Q^2
 - accurately yield the polarizability correction to hydrogen hyperfine splittings

NB The scope of this talk was limited to the inclusive measurements. A number of semi-inclusive and exclusive measurements are also in progress using the same data sets (see J. Pierce's talk in this conference).