#### HOMEWORK #7#

#### Course 314, Introduction In Quantum Mechanics, Professor K. Griffioen

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### **1 Problem** 7.17

The fundamental problem in harnessing nuclear fusion is getting the two particles (say, two deuterons) close enough together for the attractive (but short-range) nuclear force to overcame the Coulomb repulsion. The "brute force" method is to heat the particles to fantastic temperatures and allow teh random collisions to bring them together. A more exotic proposal is **muoncatalysis**, in which we construct a "hydrogen molecule ion", onlywith deuterons in place of protons, and a *muon* in place of the electron. Predict the equilibrium separation distance between the deuterons in such a structure, and explain why muons are superior to electrons for this purpose.

#### 2 Solution

The raport between atomic masses of moun and electron is :

$$\frac{m_{\mu}}{m_e} \approx 207\tag{1}$$

The point of minimum energy has x = 2.4 and is given by  $x = \frac{R}{a}$  where R is the separation of the two nuclei, and the Bohr radius is given by :

$$a = \frac{4\pi\epsilon k^2}{me^2} \tag{2}$$

We simly replace the mass of the electron with the mass of the muon :

$$R = a_{\mu}x = \frac{a_{\mu}}{a_{e}}(a_{e}x) = \frac{m_{e}}{m_{\mu}}(1.27 \times 10^{-10})m = 6.14 \times 10^{-13}m$$
(3)

## 3 Problem 8.7

Use the WKB approximation to find the allowed energies of the harmonic oscillator.

## 4 Solution

The potential (perturbation) of the harmonic oscillator is  $E = \frac{1}{2}m\omega^2 x_0^2$ , and using the WKB approximation we have :

$$p(x) = \sqrt{2m(E - \frac{1}{2}m\omega^2 x_0^2)}, \qquad (4)$$

$$\int_{-x_0}^{x_0} p(x)dx = \left(n - \frac{1}{2}\right)\pi\hbar, n = 1, 2, 3...$$

$$\int_{-x_0}^{x_0} \sqrt{2mx_0^2} \sqrt{1 - \left(\frac{x}{x_0}\right)^2} \sqrt{\frac{1}{2}m\omega^2} dx = \int_{-1}^{1} m\omega x_0^2 \sqrt{1 - u^2} du, \qquad \frac{1}{2}m\omega x_0^2 \pi = \left(n - \frac{1}{2}\right)\pi\hbar, \qquad E = \left(n - \frac{1}{2}\right)\hbar\omega$$

where we changed variables  $u = \frac{x}{x_0}$ . The result is exactly the harmonic oscillator result, since n starts with 1 in this case, instead of with 0.

## **5 Problem** 8.13

For spherically symmetrical potentials, we apply the WKB approximation to the radial equation, (Equation 4.37). In the case l = 0, it is reasonable to use Equation 8.47 in the form :

$$\int_{0}^{r_{0}} p(r)dr = (n - 1/4)\pi\hbar$$
(5)

where  $r_0$  is the turning point (in effect, we treat r = 0 as an infinite wall). Apply this formula to estimate the allowed energies of a particle in the logarithmic potential :

$$V(r) = V_0 ln(r/a) \tag{6}$$

(for constants  $V_0$  and a ). Treat only the case l = 0. Show that the spacing between the levels is independent of mass.

# 6 Solution

From WKB approximation we have :

$$\int_{0}^{r_{n}} = (n - 1/4)\pi\hbar$$
(7)

The n and (n + 1) energy levels will be given by :

$$E_{n+1} = V_0 ln \frac{r_{n+1}}{a},$$

$$E_n = V_0 ln \frac{r_n}{a},$$

$$E_{n+1} - E_n = V_0 \left( ln \frac{r_{n+1}}{a} - ln \frac{r_n}{a} \right),$$

$$E_{n+1} - E_n = V_0 ln \frac{r_{n+1}}{r_n}$$
(9)

We'll plug into (7) and we'll obtain :

$$ln \int_{0}^{r_{n}} \sqrt{2m \left(V_{0} ln \frac{r_{n}}{a} - V_{0} ln \frac{r}{a}\right)} dr = ln \left[\left(n - \frac{1}{4}\right) \pi \hbar\right],$$
(10)  
$$ln \int_{0}^{r_{n+1}} \sqrt{2m \left(V_{0} ln \frac{r_{n+1}}{a} - V_{0} ln \frac{r}{a}\right)} dr = ln \left[\left(n + \frac{3}{4}\right) \pi \hbar\right],$$
$$ln \sqrt{2mV_{0}} + ln \int_{0}^{r_{n}} \sqrt{ln(r_{n}/r)} dr = ln \pi \hbar + ln(n - 1/4),$$
$$ln \sqrt{2mV_{0}} + ln \int_{0}^{r_{n+1}} \sqrt{ln(r_{n+1}/r)} dr = ln \pi \hbar + ln(n + 3/4)$$

We'll substract the above relations one from each other and we'll make the following change of variable:

$$u = \frac{r}{r_n}, du = \frac{dr}{r_n},\tag{11}$$

$$u = \frac{r}{r_{n+1}}, du = \frac{dr}{r_{n+1}}$$

We'll obtain the relations :

$$ln\left(\frac{n+3/4}{n-1/4}\right) = lnr_{n+1} - lnr_n,$$
(12)  
$$ln\frac{r_{n+1}}{r_n} = \frac{1}{V_0}(E_{n+1} - E_n),$$
  
$$E_{n+1} - E_n = V_0 ln\left(\frac{n+3/4}{n-1/4}\right)$$

## 7 Problem 9.1

A hidrogen atom is placed in a (time-dependent) electric field  $E = E(t)\hat{k}$ . Calculate all four matrix elements  $H'_{ij}$  of the perturbation H' = -eEz between the ground state (n = 1) and the (quadruply degenerate) first excited states (n = 2). Also show that  $H'_{ij} = 0$  for all five states.*Note*: there is only one integral to be done here, if you exploit oddness with respect to z. As a result, only one of the n = 2 states is "accessible" from the ground state by a perturbation of this form, and therefore the system functions as a two-level configuration - assuming transitions to higher excited states can be ignored.

### 8 Solution

$$H' = -eEz, z = r\cos\theta = \sqrt{\frac{4\pi}{3}}rY_1^0,$$
(13)

where  $Y_l^m$  are the angular functions of the wave functions and :

$$\int Y_{l'm'}^* Y_{lm} d\omega = \delta_{ll'} \delta_{mm'}, \qquad (14)$$
$$Y_{00} = 0$$

The wave functions are given by :

$$\psi_{100} = Y_{00}R_{10}, \tag{15}$$

$$\psi_{200} = Y_{00}R_{20},$$
  

$$\psi_{210} = Y_{10}R_{21},$$
  

$$\psi_{21,-1} = Y_{1,-1}R_{21},$$
  

$$\psi_{211} = Y_{11}R_{21}$$

The matrix elements of H'are given by :

$$H'_{ij} = \langle n'l'm'|H'|nlm \rangle = -eE \langle 100|z|2lm \rangle = -eEY_{00}Y_{10}Y_{lm}$$
(16)

But from the definition of  $<\psi|H|\psi>$  we get that :

$$\int d\Omega Y_{00} Y_{10} Y_{lm} = Y_{00} \int d\Omega Y_{10}^* Y_{lm} = Y_{00} \delta_{l1} \delta_{m0}$$
(17)

And we can observe that only  $\psi_{210}$  is possible. We plug  $\psi_{210}$  into (16) and the nonvanishing term of H' will be :

$$<\psi_{210}|H'|\psi_{210} = -eE < R_{10}|r|R_{21} > \frac{1}{\sqrt{12\pi}},$$

$$\int_{0}^{\infty} r^{3} \frac{2}{a^{3/2}} e^{-r/a} \frac{1}{\sqrt{24}} \frac{1}{a^{3/2}} \frac{r}{a} e^{-r/2a} dr = \frac{2}{a^{4}} \frac{1}{\sqrt{24}} \left(\frac{2a}{3}\right)^{5} \int_{0}^{\infty} u^{4} e^{-u} du,$$

$$< H' > = -eEa \frac{2^{15/2}}{3^{5}} = -0.745 eEa$$
(18)

# References

[1] D.J.Griffiths, Introduction To Quantum Mechanics, 1995