

Precision measurements assisted by Rb vapor

Eugeniy E. Mikhailov²

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WILLIAM & MARY
CHARTERED 1693

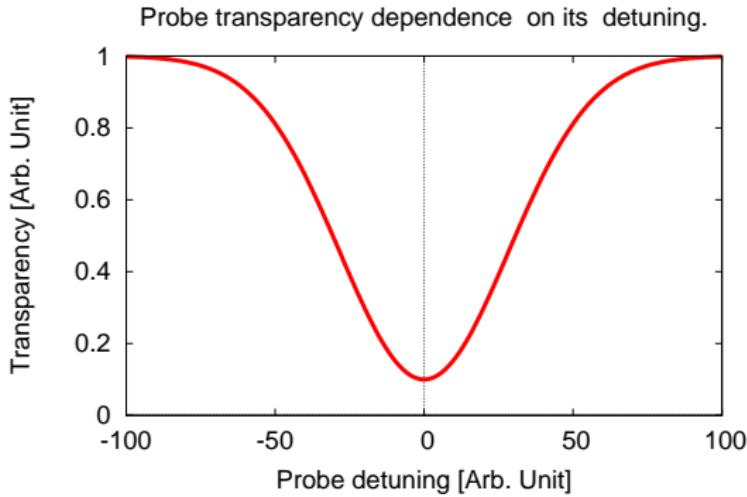
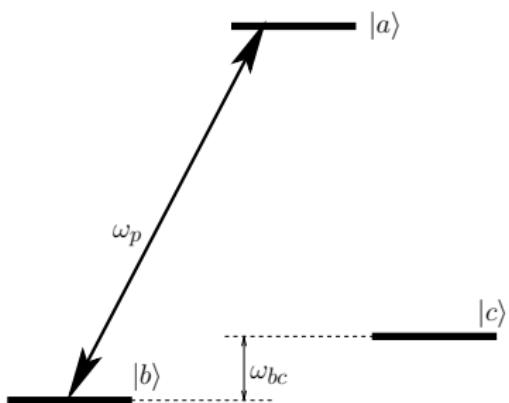
About William & Mary



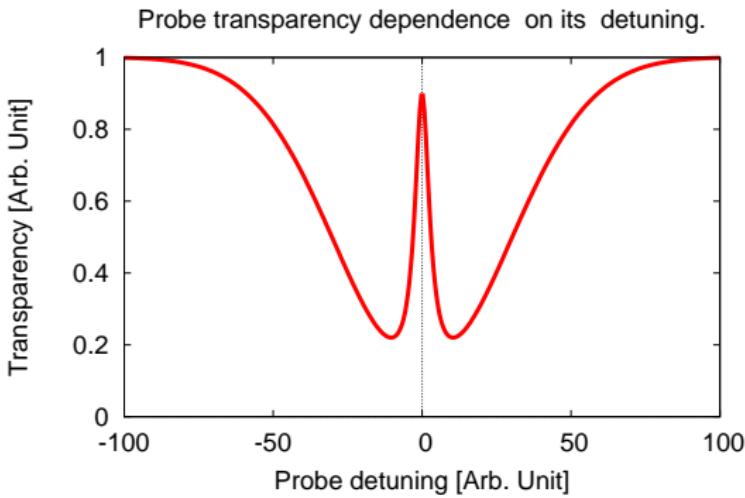
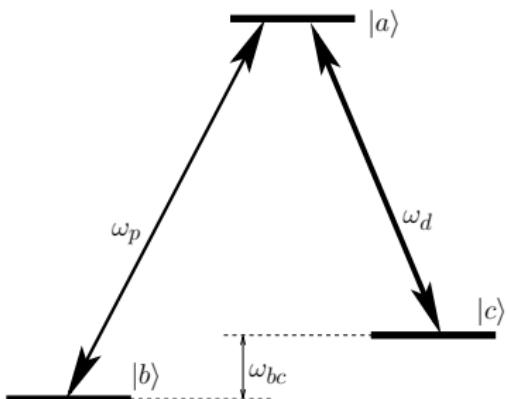
Outline

- Magnetic compass
- Quantum noise imaging

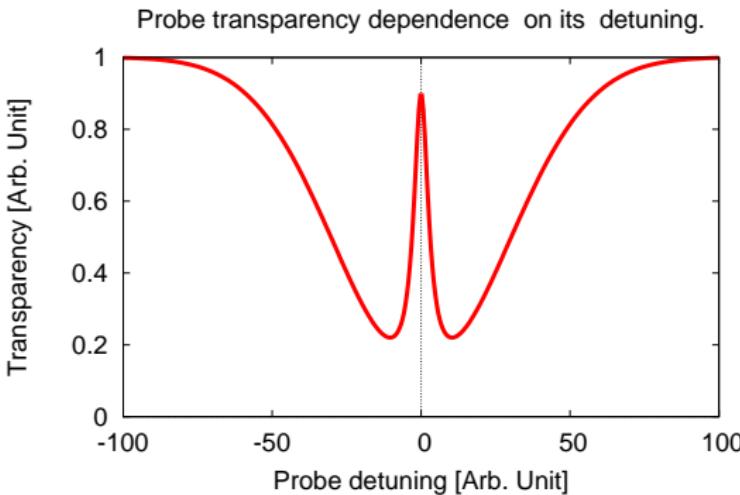
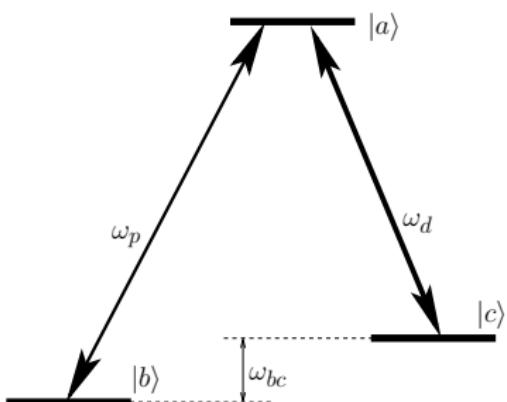
Electromagnetically Induced Transparency (EIT)



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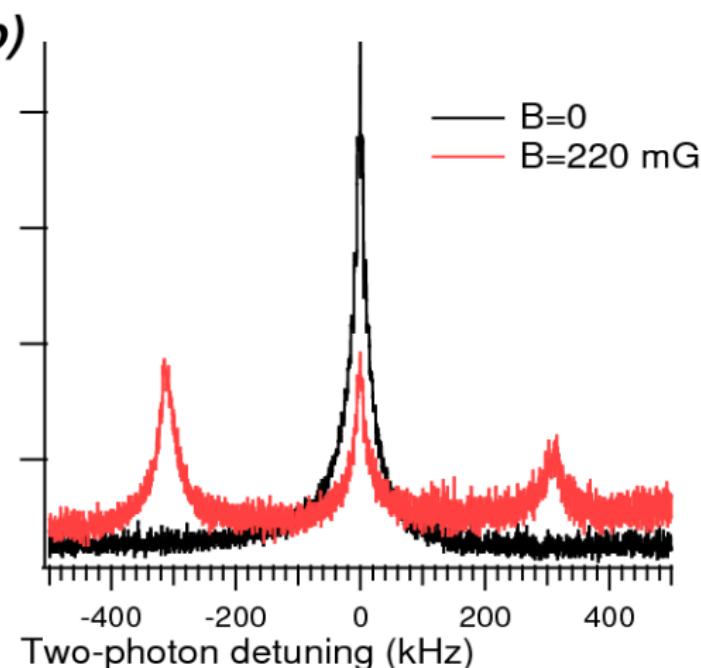
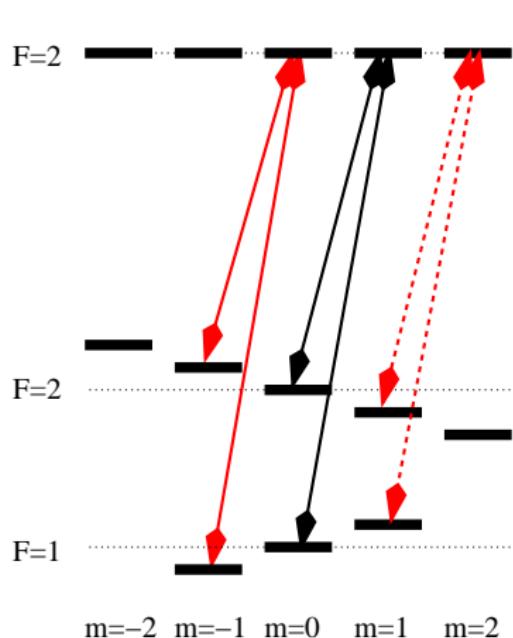


Coherent Population Trapping

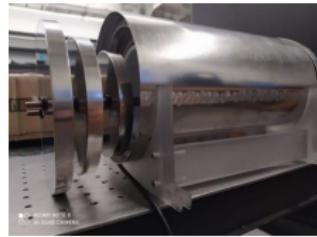
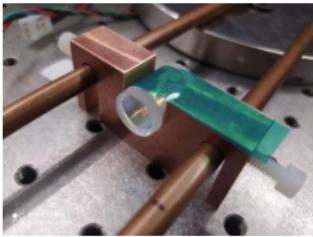
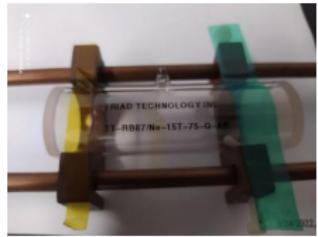
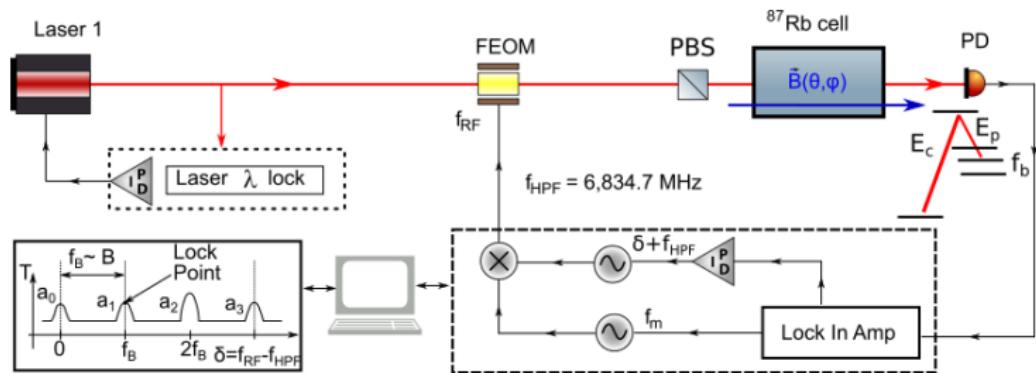
- Dark $|D\rangle = \Omega_d|b\rangle - \Omega_p|c\rangle$ and Bright $|B\rangle = \Omega_d|c\rangle + \Omega_p|b\rangle$ states
- resonance width ($\sim 10\text{kHz}$) much smaller than natural line width

Simple EIT magnetometer

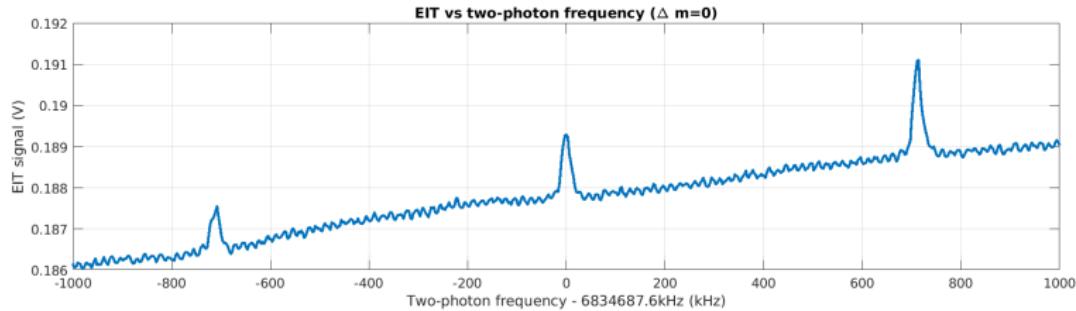
EIT with circularly polarized light



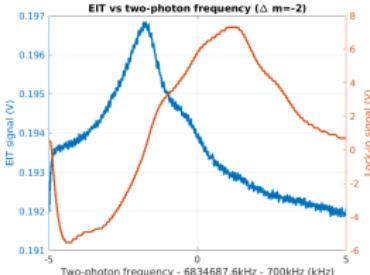
Conceptual design



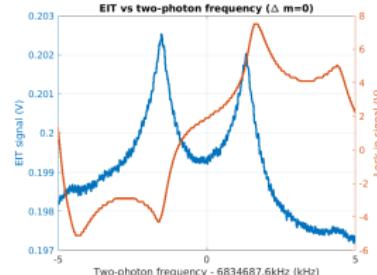
EIT signals vs two-photon detuning. $B=50\mu\text{T}$, $f_{center} = 6'834'687.6\text{kHz}$



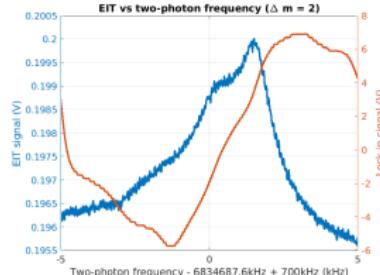
$\Delta m = -2$,
 $f_{center} - 700\text{kHz}$



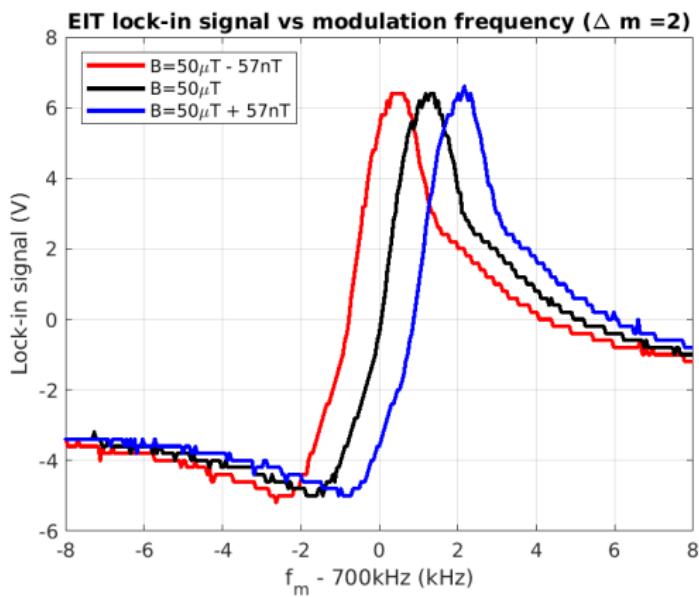
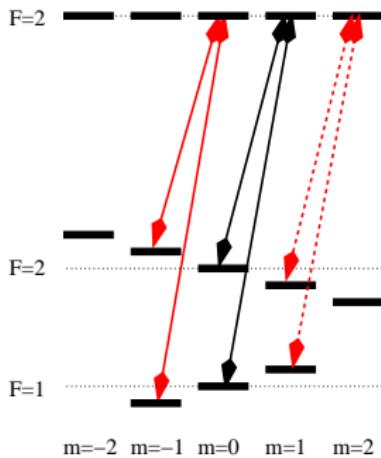
$\Delta m = 0$,
 f_{center}



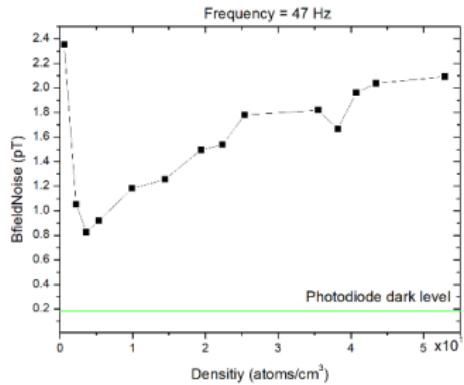
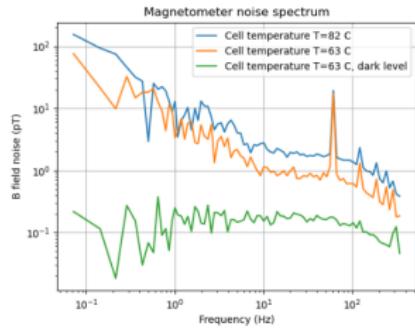
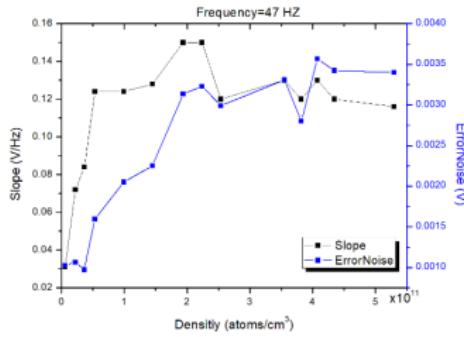
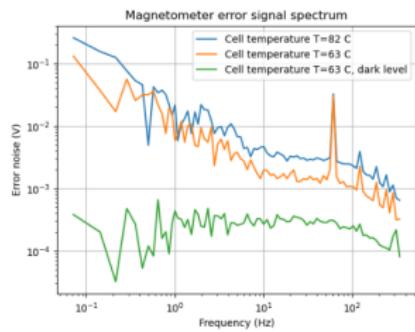
$\Delta m = 2$,
 $f_{center} + 700\text{kHz}$



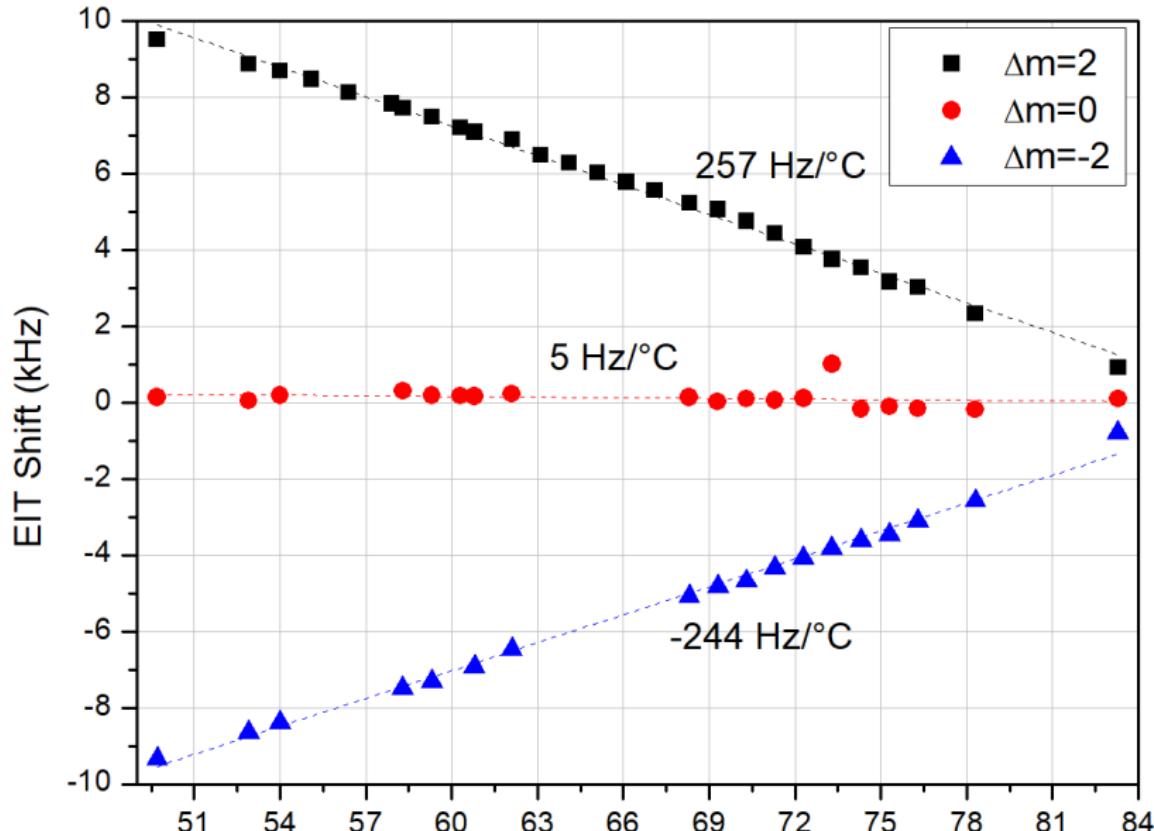
Shift by $\pm 57\text{nT}$ magnetic field change



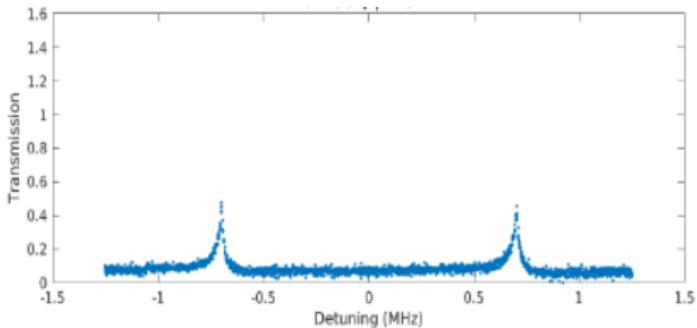
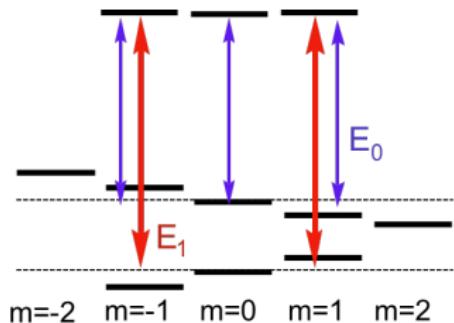
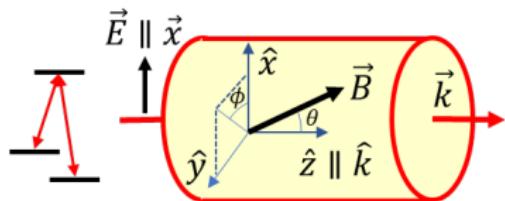
Signal to noise optimization



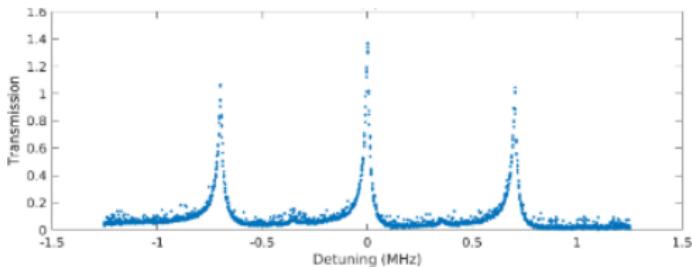
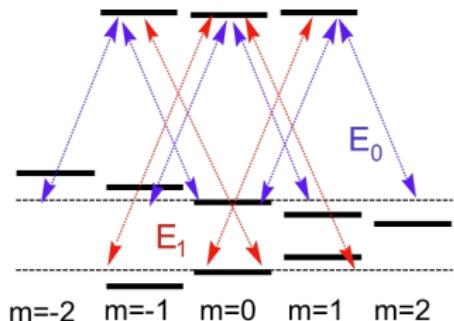
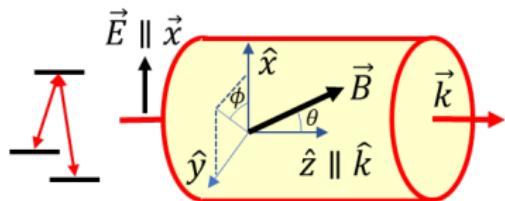
EIT shift vs temperature



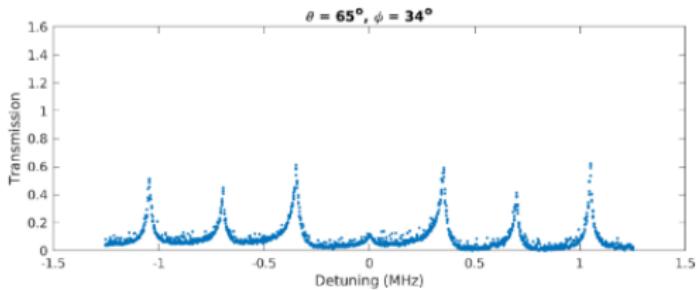
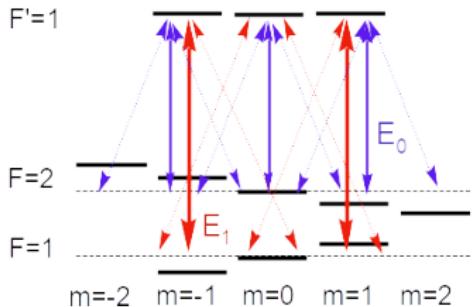
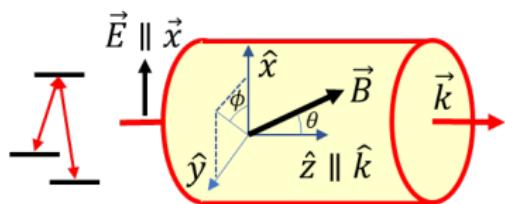
Compass idea



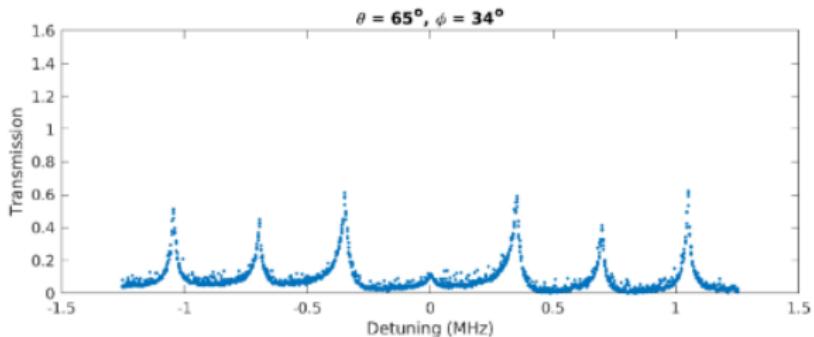
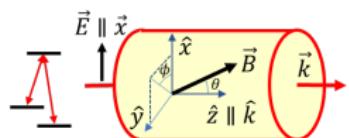
Compass idea



Compass idea

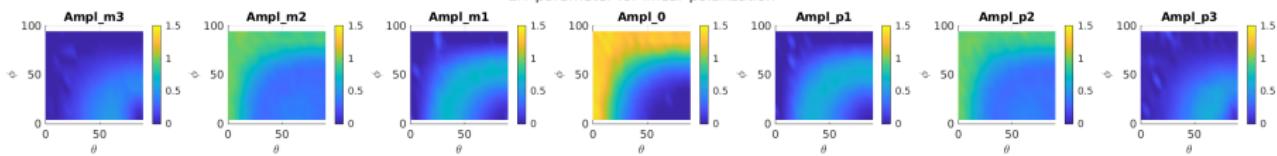


Linear polarization: angular dependence on θ and ϕ



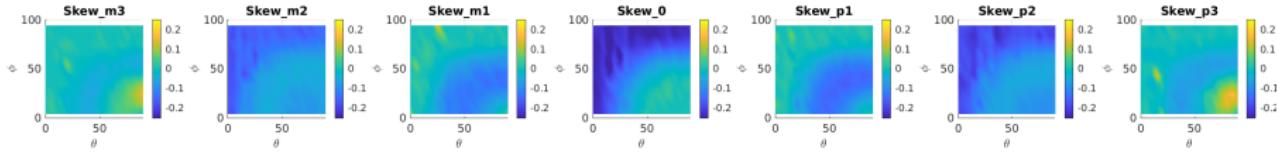
Amplitude

EIT parameter for linear polarization

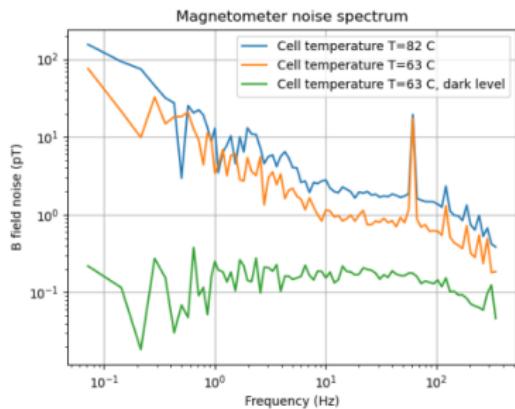
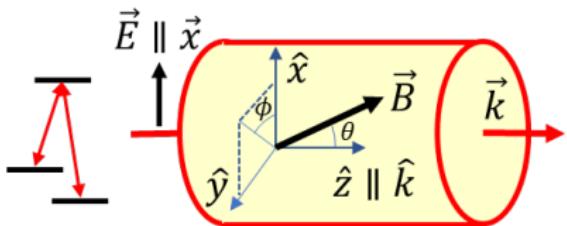


Skew

EIT parameter for linear polarization



Compass summary



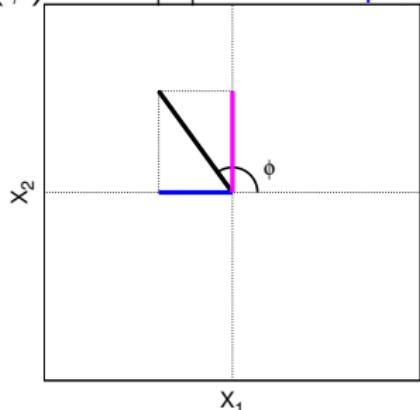
- measures B-field vector
- can operate at the Earth magnetic field
- sub-pT sensitivity

Transition from classical to quantum field

Classical analog

- Field amplitude a
- Field real part
 $X_1 = (a^* + a)/2$
- Field imaginary part
 $X_2 = i(a^* - a)/2$

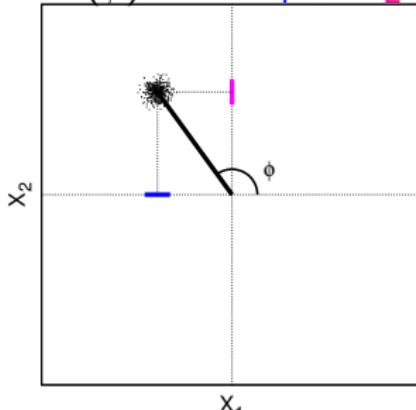
$$E(\phi) = |a|e^{-i\phi} = X_1 + iX_2$$



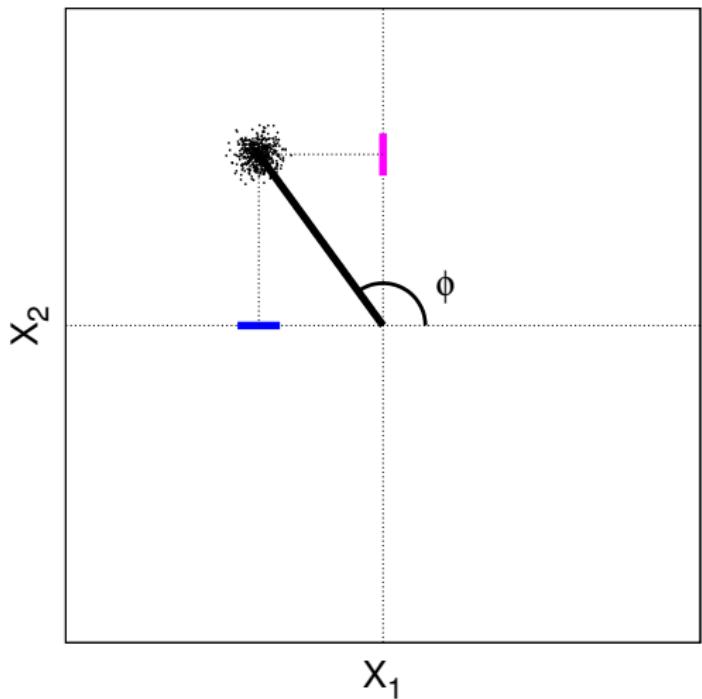
Quantum approach

- Field operator \hat{a}
- Amplitude quadrature
 $\hat{X}_1 = (\hat{a}^\dagger + \hat{a})/2$
- Phase quadrature
 $\hat{X}_2 = i(\hat{a}^\dagger - \hat{a})/2$

$$\hat{E}(\phi) = \hat{X}_1 + i\hat{X}_2$$



Quantum optics summary



Light consist of photons

- $\hat{N} = a^\dagger a$

Commutator relationship

- $[a, a^\dagger] = 1$

- $[X_1, X_2] = i/2$

Detectors measure

- number of photons \hat{N}

- Quadratures \hat{X}_1 and \hat{X}_2

Uncertainty relationship

- $\Delta X_1 \Delta X_2 \geq 1/4$

Heisenberg uncertainty principle and its optics equivalent



Heisenberg uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa

Heisenberg uncertainty principle and its optics equivalent



Heisenberg uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa

Optics equivalent

$$\Delta\phi \Delta N \geq 1$$

The more precisely the PHASE is determined, the less precisely the AMPLITUDE is known, and vice versa

Heisenberg uncertainty principle and its optics equivalent



Heisenberg uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa

Optics equivalent

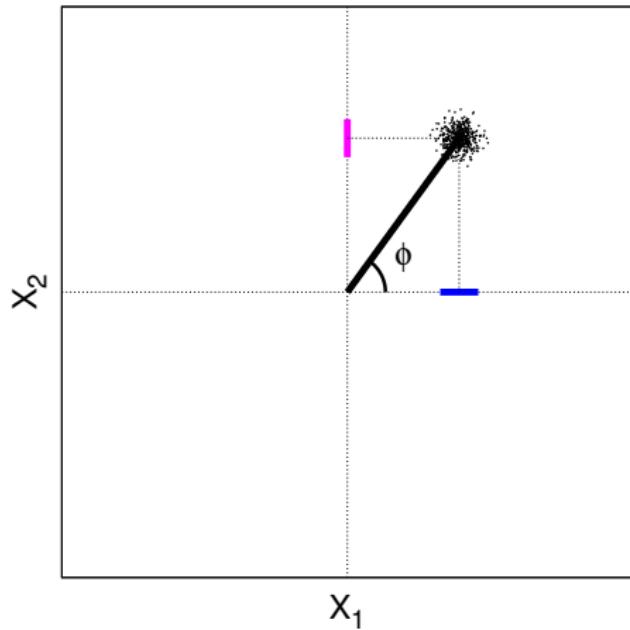
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The more precisely the PHASE is determined, the less precisely the AMPLITUDE is known, and vice versa

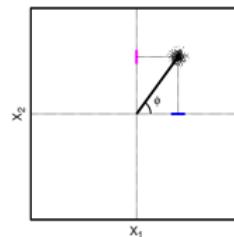
Optics equivalent strict definition

$$\Delta X_1 \Delta X_2 \geq 1/4$$

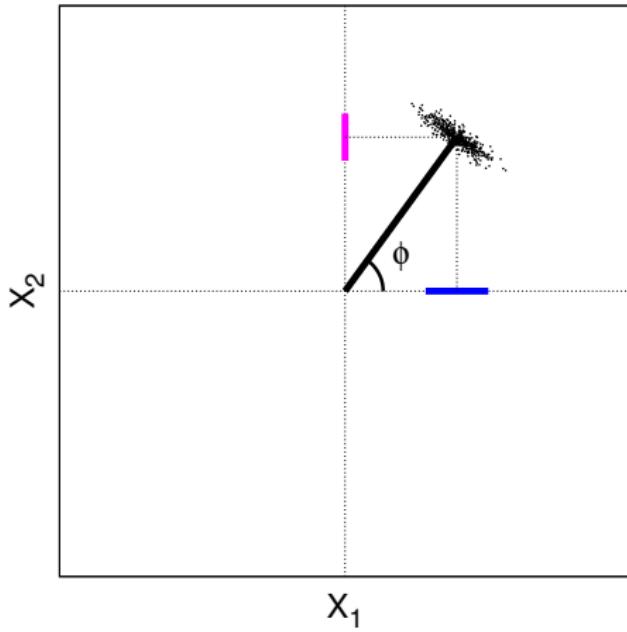
Squeezed quantum states zoo



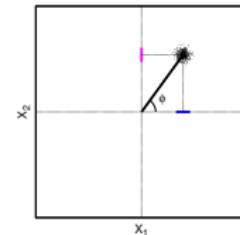
Unsqueezed
coherent



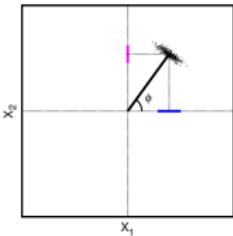
Squeezed quantum states zoo



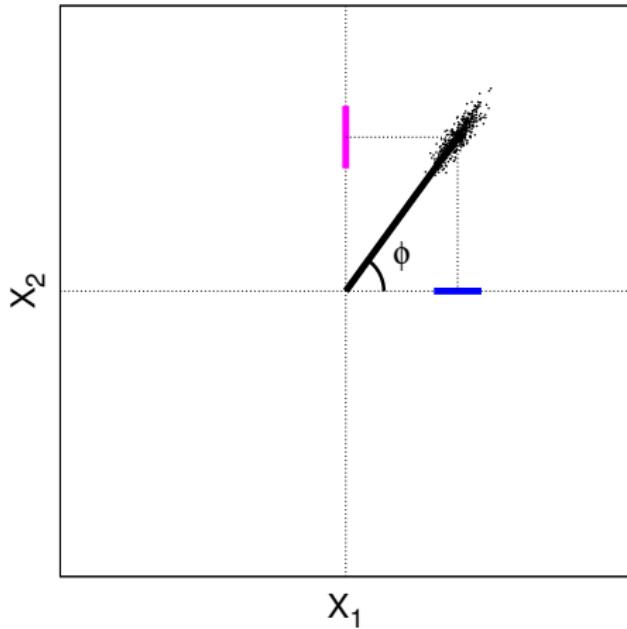
Unsqueezed
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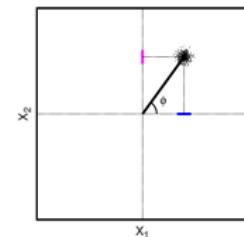
Amplitude
squeezed



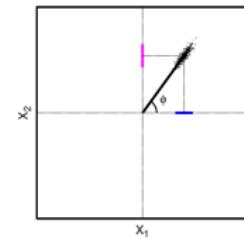
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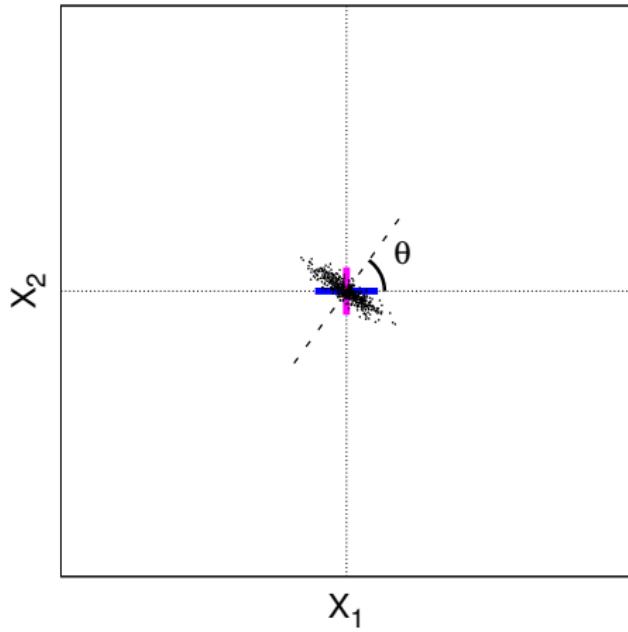
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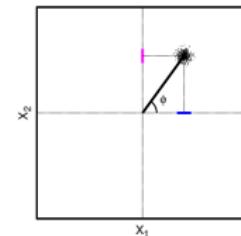
Amplitude
squeezed



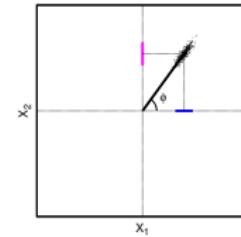
Squeezed quantum states zoo



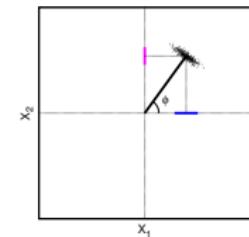
Unsqueezed
coherent



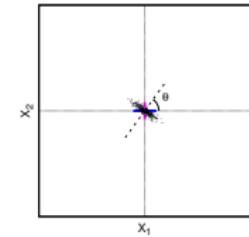
Phase
squeezed



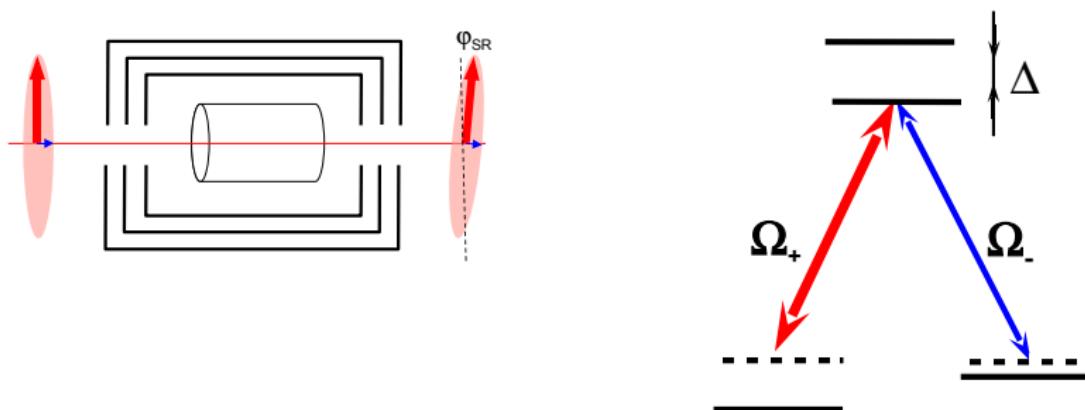
Amplitude
squeezed



Vacuum
squeezed



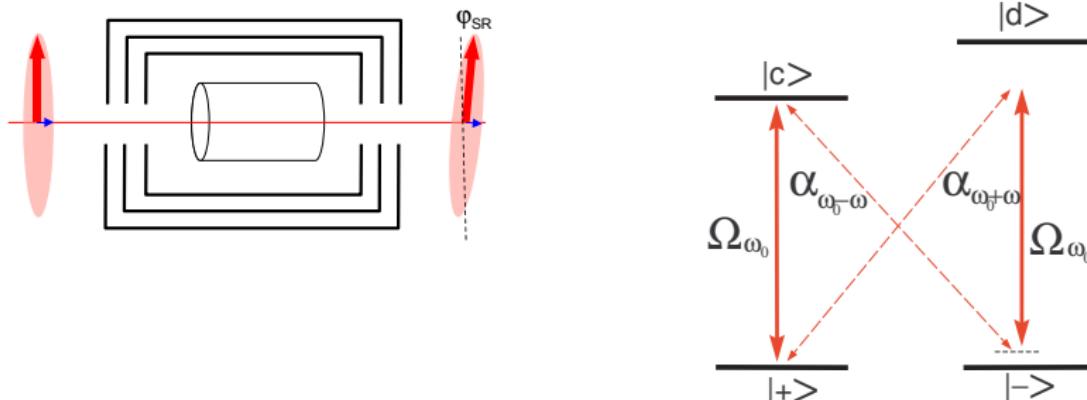
Self-rotation of elliptical polarization in atomic medium



A. B. Matsko, I. Novikova, G. R. Welch, D. Budker, D. F. Kimball, and S. M. Rochester, PRA 66, 043815 (2002):
theoretical prediction of 4–6 dB noise suppression

$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^\dagger - a_{in})$$

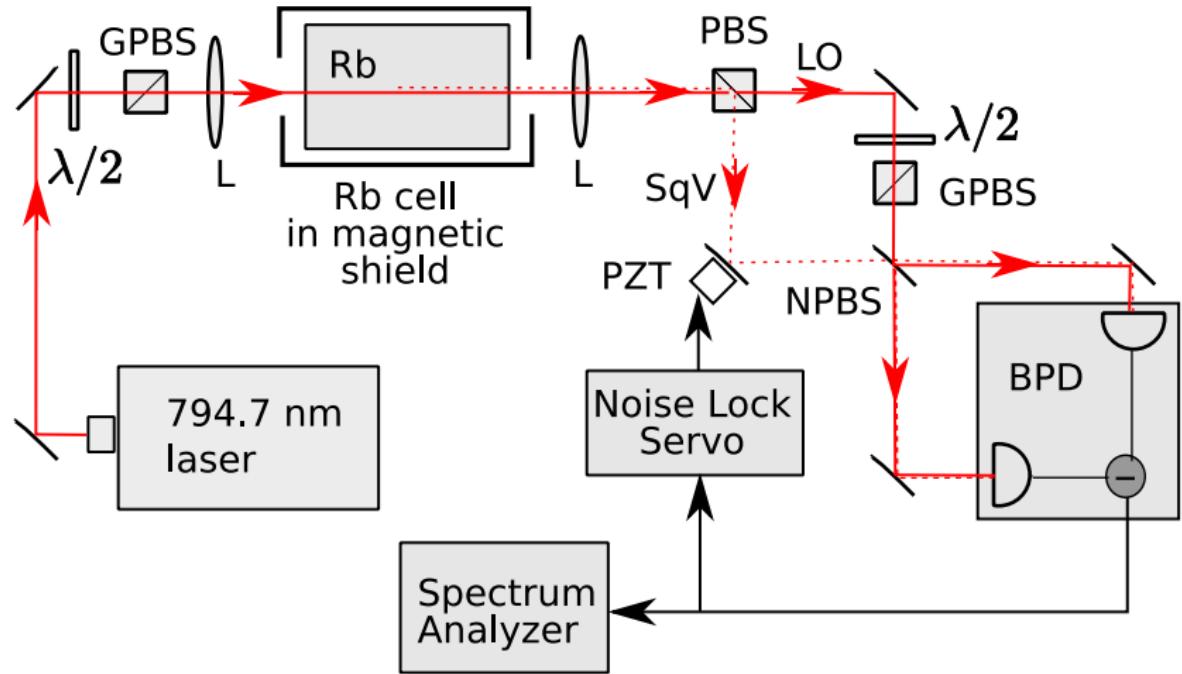
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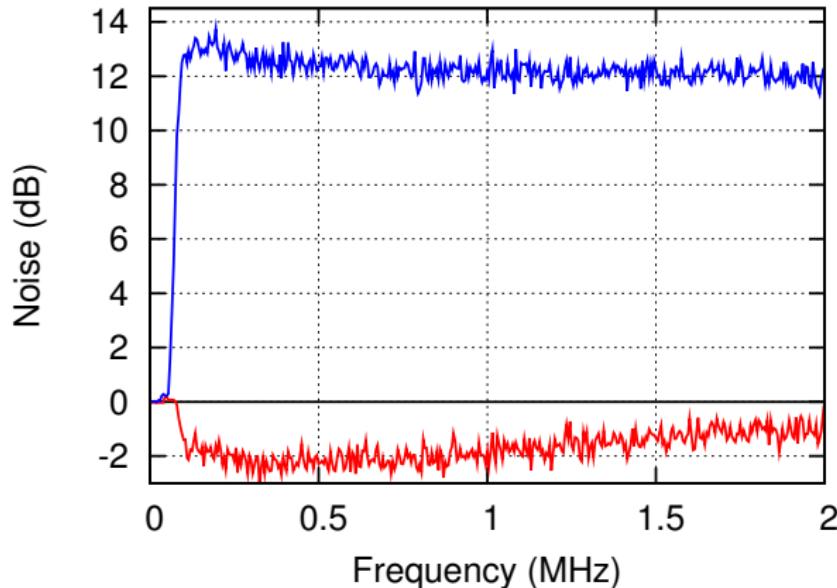
$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^\dagger - a_{in})$$

Setup



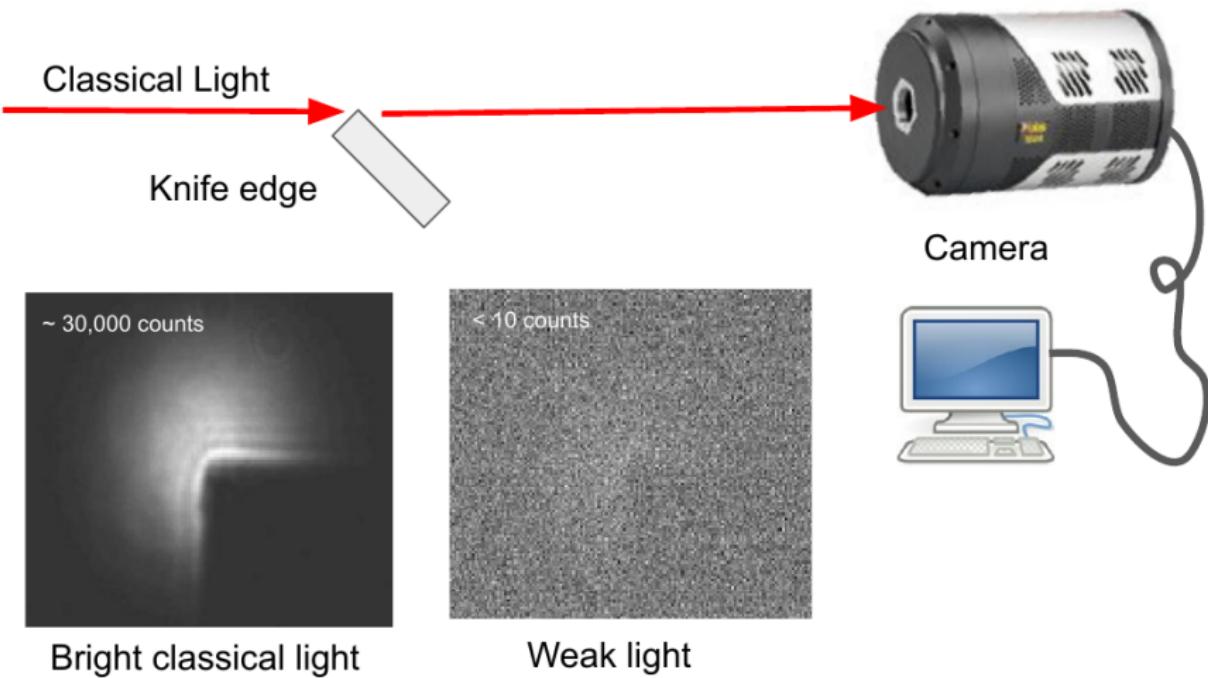
Maximally squeezed spectrum with ^{87}Rb

W&M team. $^{87}\text{Rb } F_g = 2 \rightarrow F_e = 2$, laser power 7 mW, $T=65^\circ\text{ C}$

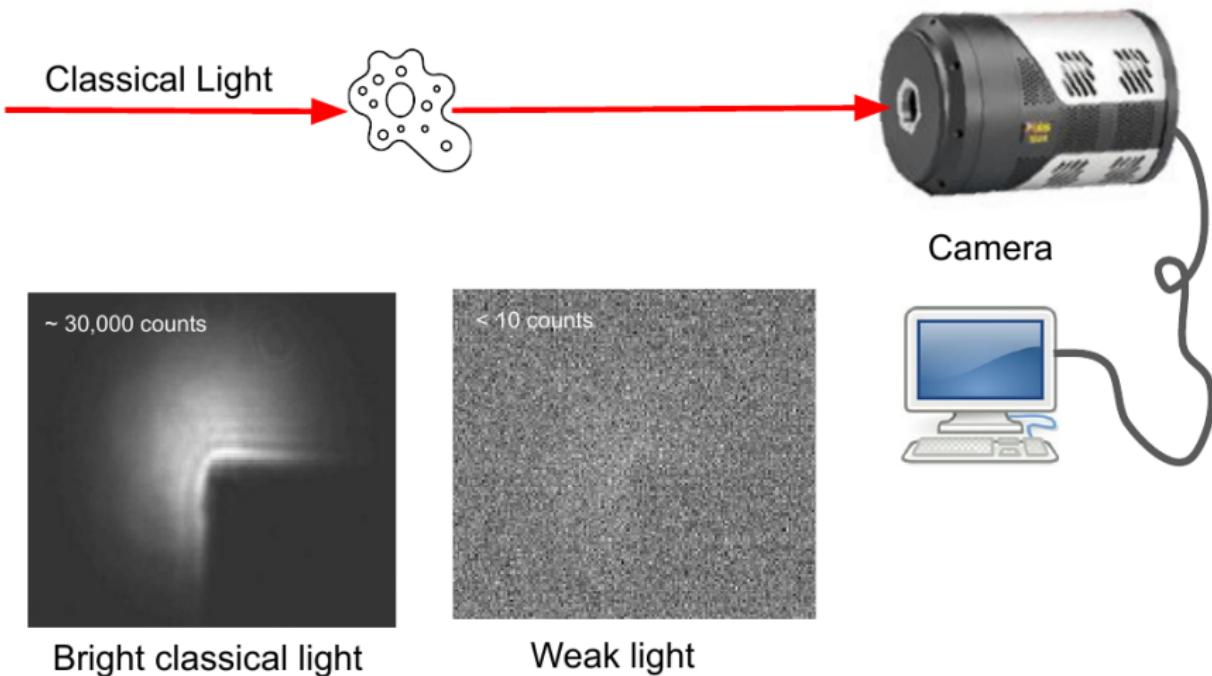


Lezama et.al report 3 dB squeezing in similar setup
Phys. Rev. A 84, 033851 (2011)

From bright to low light imaging

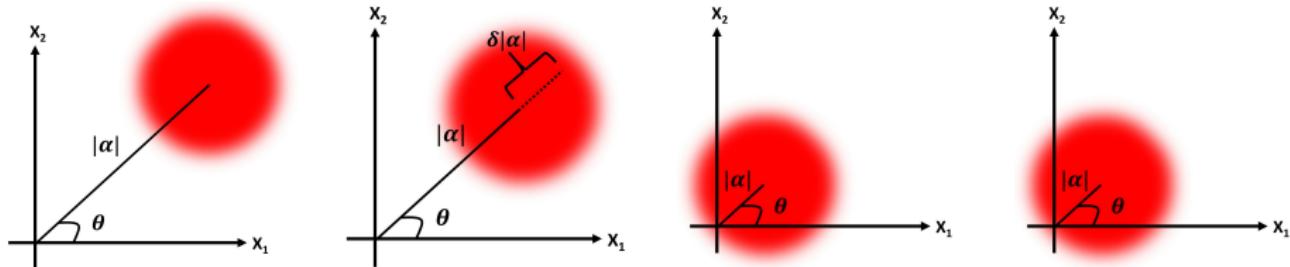


From bright to low light imaging

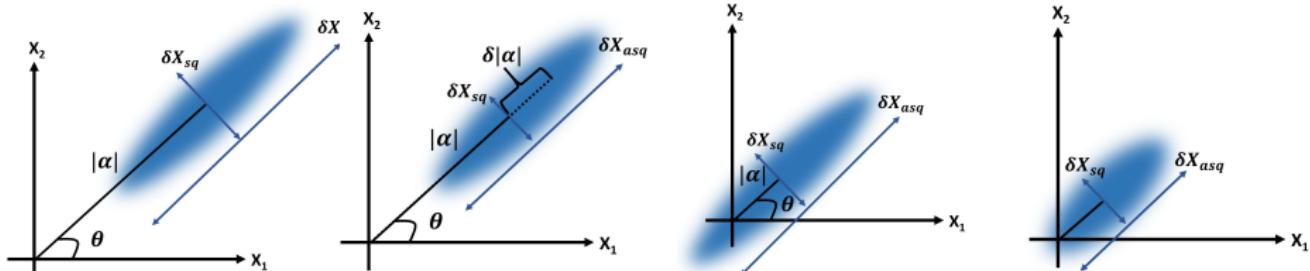


Let's look at quantum picture

Bright state in Bright state out Low-light state in Low-light out

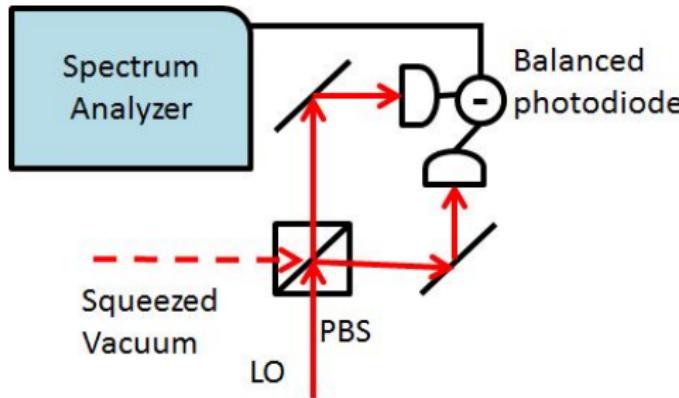
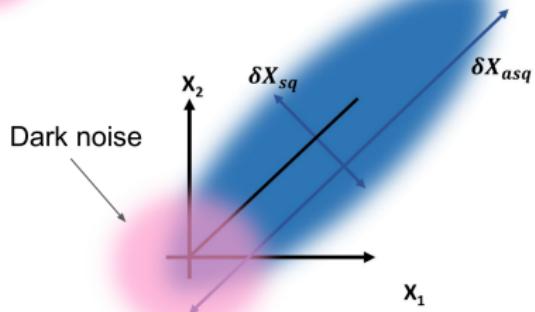
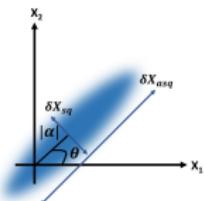
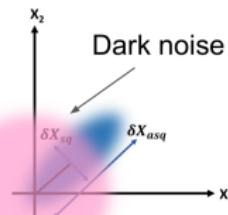


$$\alpha_{out}^2 = \alpha_{in}^2 T$$

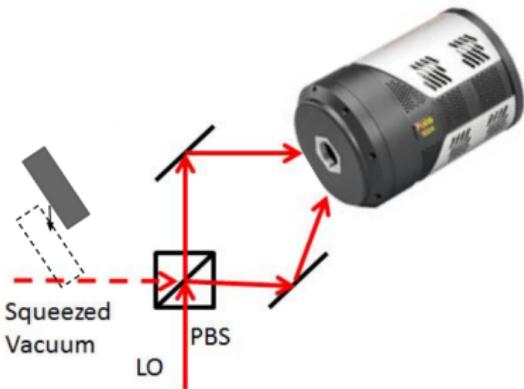
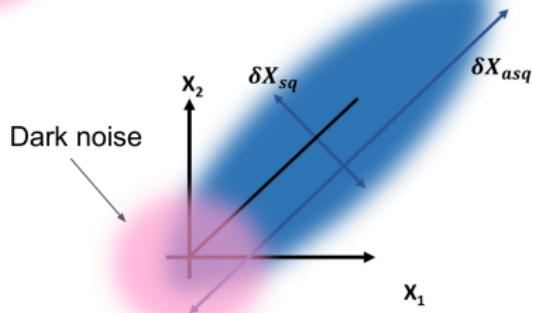
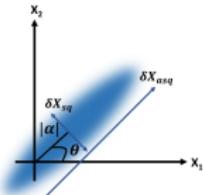
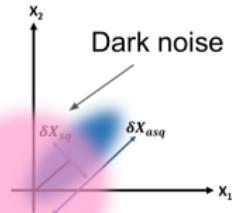


$$V = 1 + (\delta X_{sq/asq}^2 - 1) |\mathcal{O}|^2 T$$

Detector dark noise

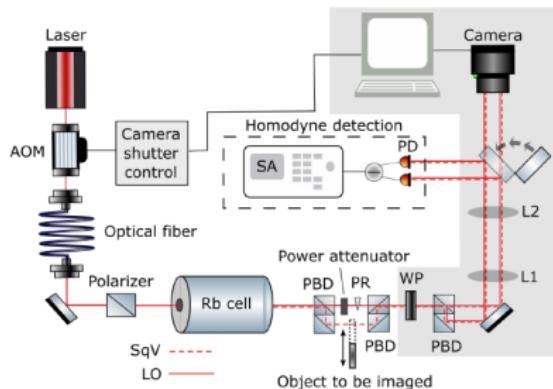
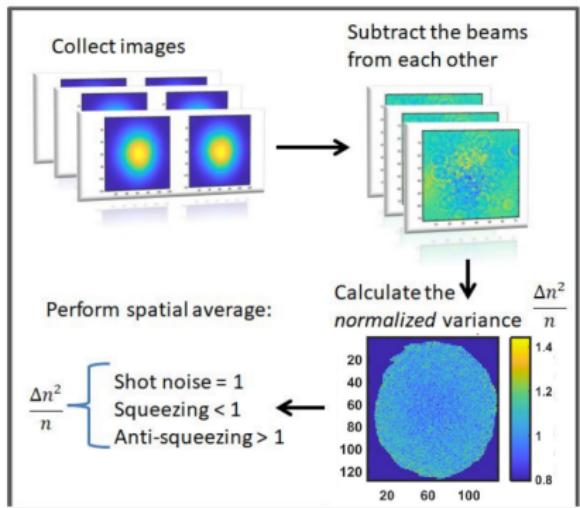


Detector dark noise

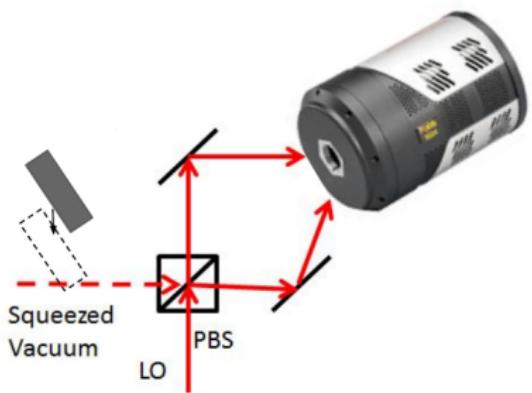


"Quantum-Limited Squeezed Light Detection with a Camera", Phys. Rev. Lett. **125**, 113602

Imaging quantum noise

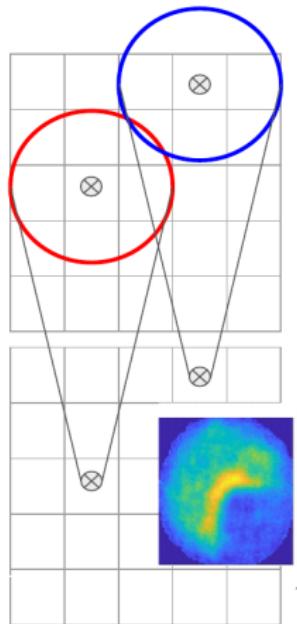
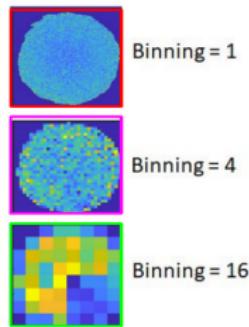


Imaging quantum noise with binning

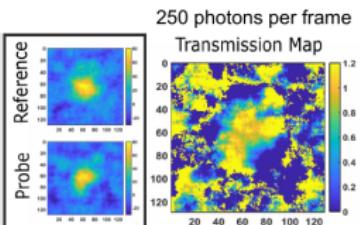
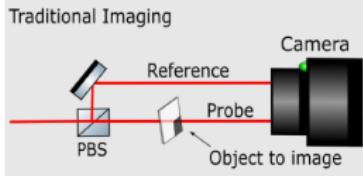
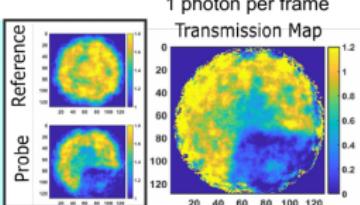
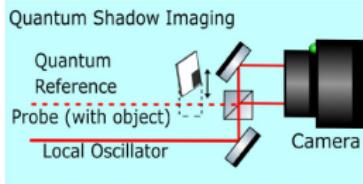


$$V = 1 + (\delta X_{sq/asq}^2 - 1)|\mathcal{O}|^2 T$$

- Single pixel analysis = shot noise limited
- Binning pixels reveals non-classical statistics

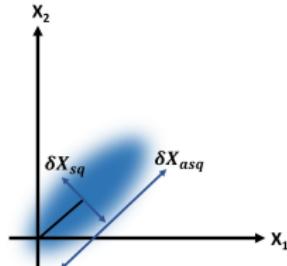


Shadow imaging



$$V_{pr} = 1 + (\delta X_{ref}^2 - 1)|\mathcal{O}|^2 T$$

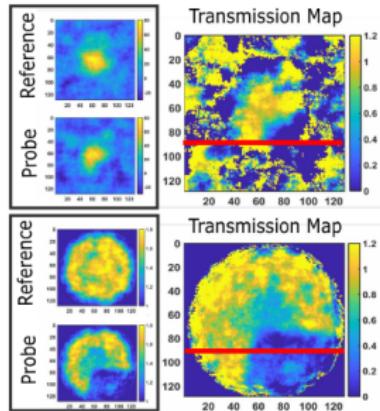
$$T = \frac{V_{pr} - 1}{V_{ref} - 1}$$



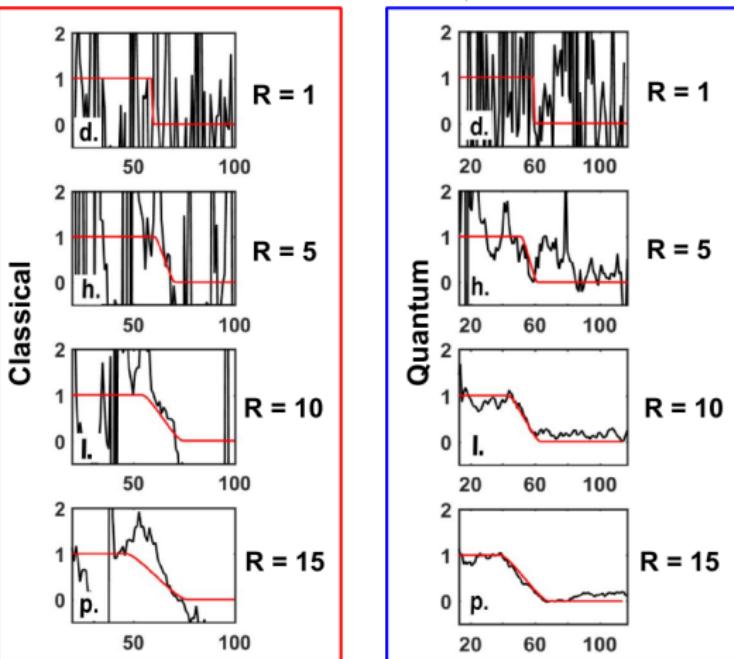
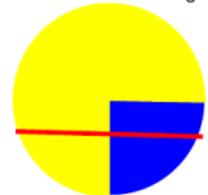
$$T = \frac{|\alpha_{pr}|^2}{|\alpha_{ref}|^2} = \frac{N_{pr}}{N_{ref}}$$

Similarity Parameter

Transmission Map Cross-section



Ideal case: T_o



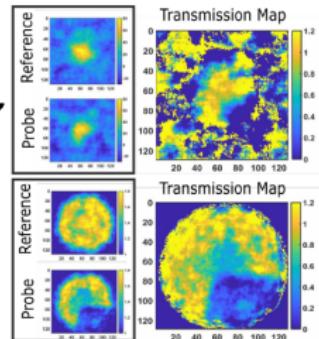
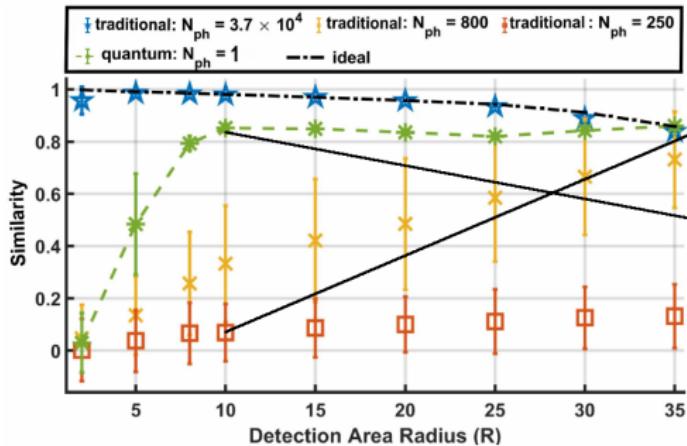
$$S = \frac{\sum T_{exp} T_o}{\sqrt{\sum T_{exp}^2 \sum T_o^2}}$$

Similarity Parameter

$$S = \frac{\sum T_{exp} T_o}{\sqrt{\sum T_{exp}^2 \sum T_o^2}}$$



Savannah Couzzo

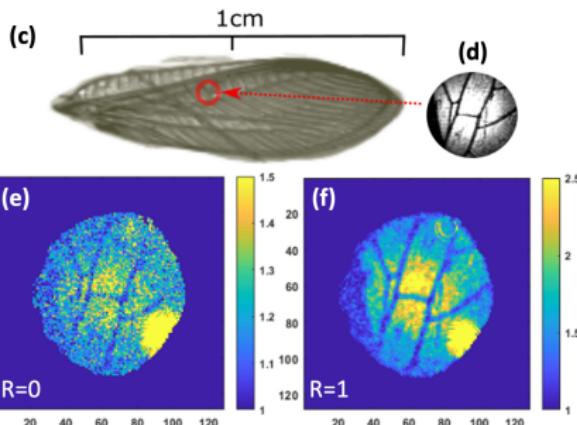
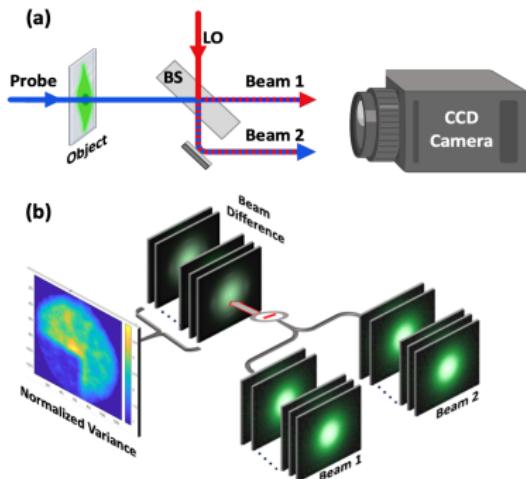


"Low-Light Shadow Imaging using Quantum-Noise Detection with a Camera"
<https://arxiv.org/abs/2106.00785>

Imaging with thermal light



Ziqi Niu



Structural light imaging: single pixel camera

Single-pixel imaging 12 years on: a review

GRAHAM M. GIBSON,^{1,2}  STEVEN D. JOHNSON,^{1,3}  AND MILES J. PADGETT^{1,4} 

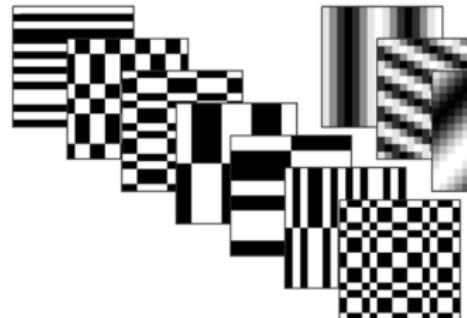
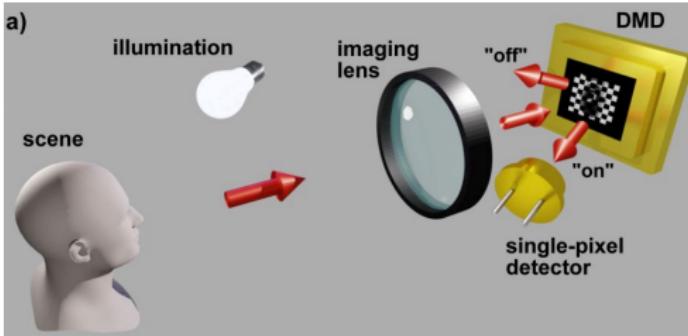
¹School of Physics and Astronomy, University of Glasgow, Glasgow G12 8QQ, UK

²graham.gibson@glasgow.ac.uk

³steven.johnson@glasgow.ac.uk

⁴miles.padgett@glasgow.ac.uk

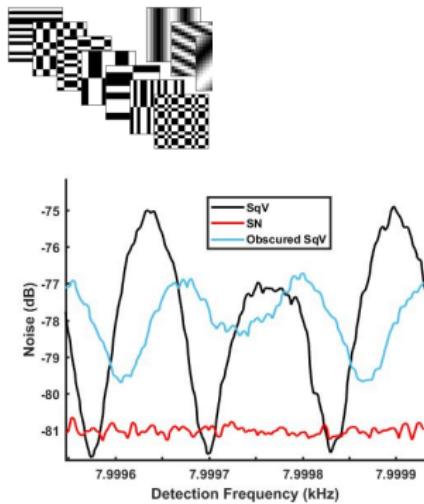
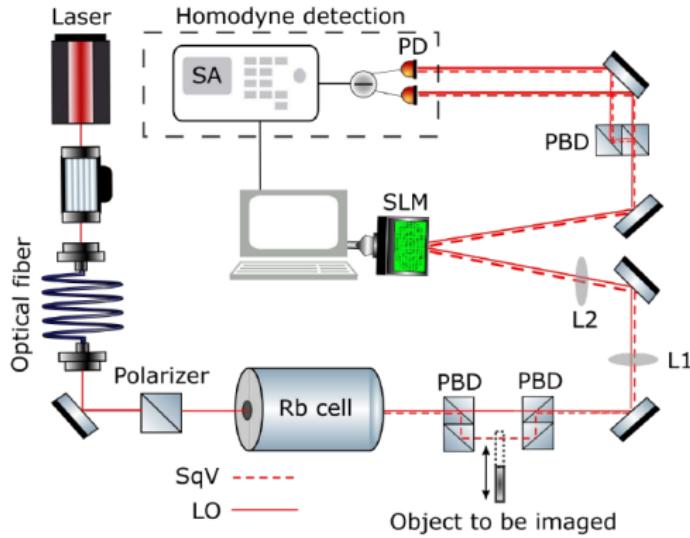
<https://www.gla.ac.uk/schools/physics/ourresearch/groups/optics/>



Reconstructed object

$$Ob(x, y) = \frac{1}{M} S_m P_m(x, y)$$

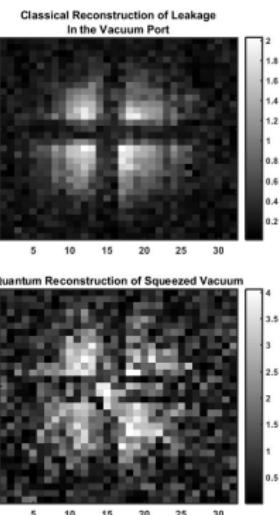
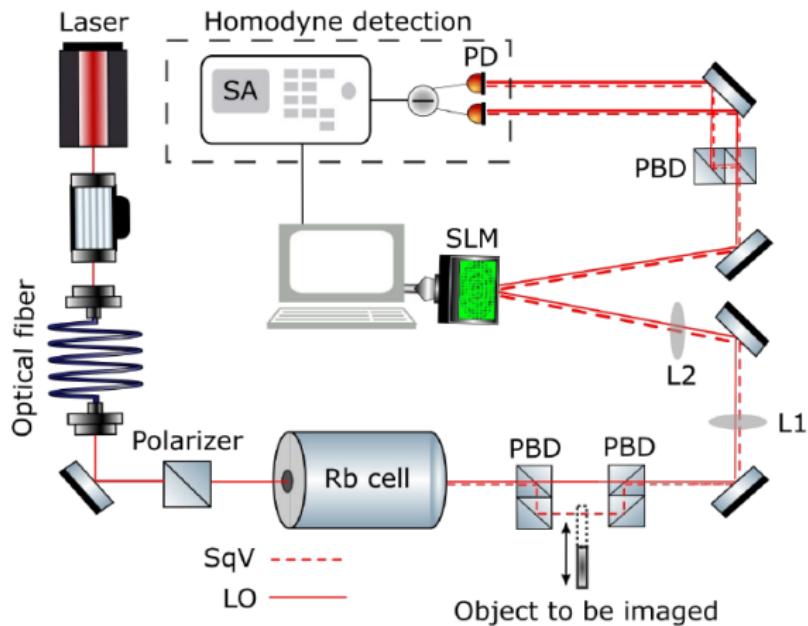
Structural light imaging with quantum noise



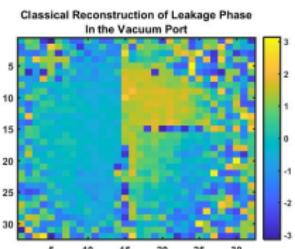
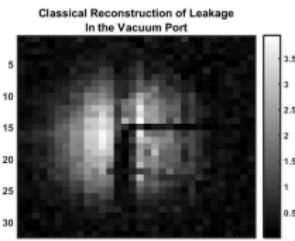
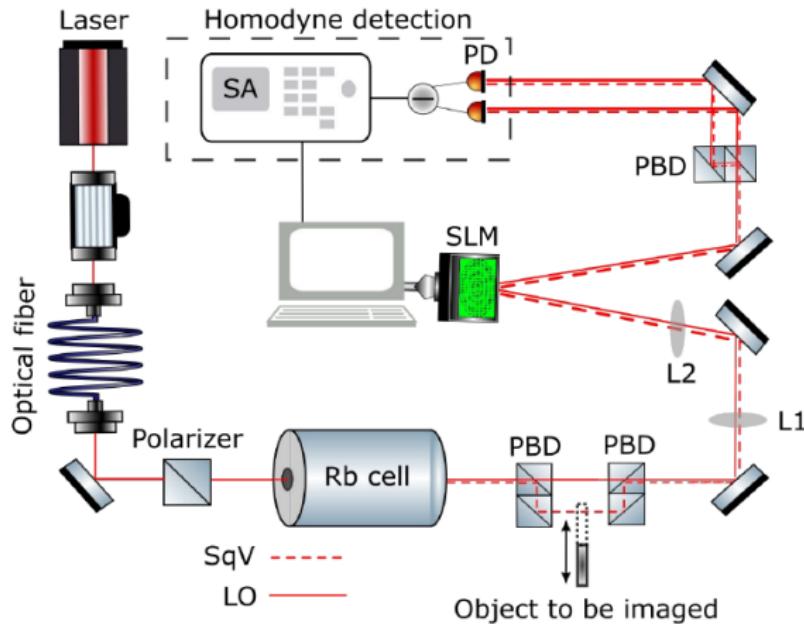
$$V_m = 1 + (\delta X_{sq/asq}^2 - 1) |\mathcal{O}_m|^2$$
$$\mathcal{O}_m = \int_A P_m u_{lo} u_q^* T dA$$

$$u_{lo} u_q^* T = \frac{1}{M} \sum \mathcal{O}_m P_m(x, y)$$

Structural light imaging with quantum noise



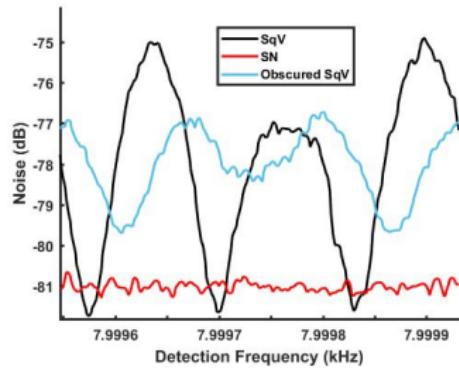
Structural light imaging with quantum noise



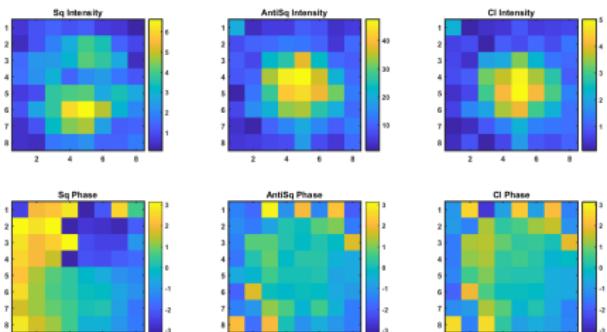
Structural light using different quadratures

$$V_m = 1 + (\delta \chi_{sq/asq}^2 - 1) |\mathcal{O}_m|^2$$

$$\mathcal{O}_m = \int_A P_m u_{lo} u_q^* T dA$$

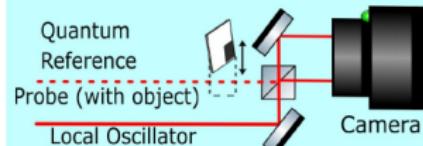


$$u_{lo} u_q^* T = \frac{1}{M} \sum \mathcal{O}_m P_m(x, y)$$



Summary

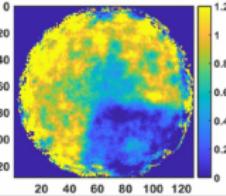
Quantum Shadow Imaging



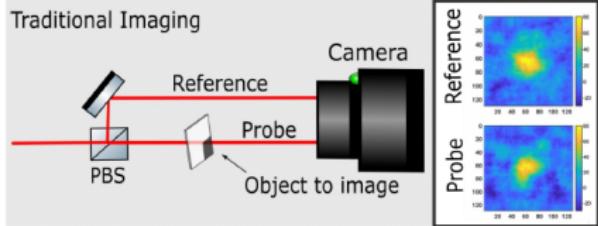
Reference

Probe

1 photon per frame
Transmission Map



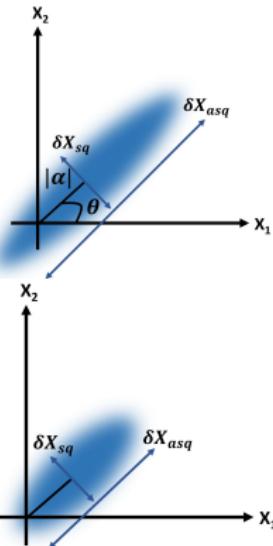
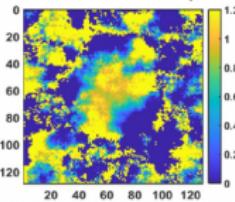
Traditional Imaging



Reference

Probe

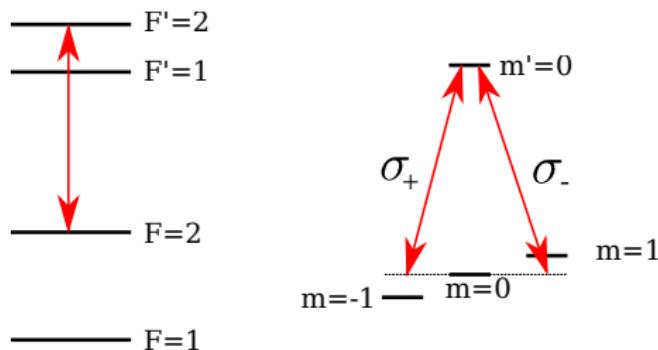
250 photons per frame
Transmission Map



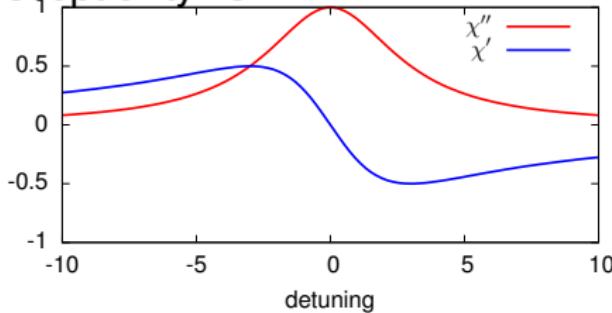
“Low-Light Shadow Imaging using Quantum-Noise Detection with a Camera” <https://arxiv.org/abs/2106.00785>

Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

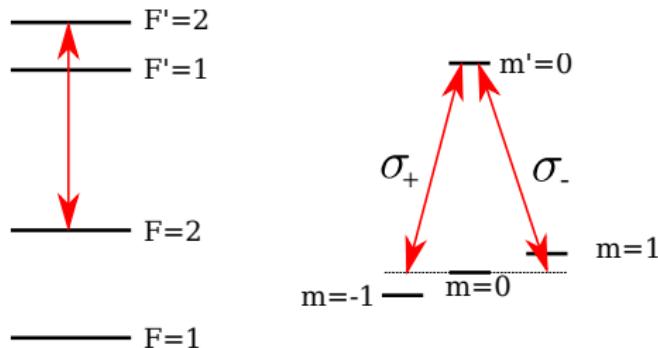


Susceptibility vs B

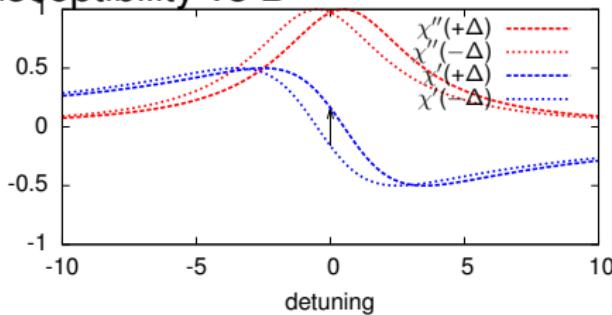


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

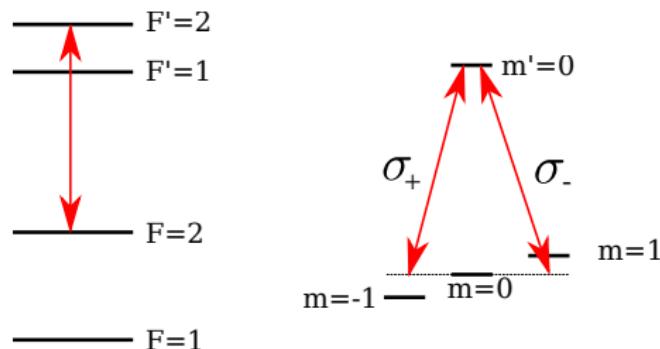


Susceptibility vs B

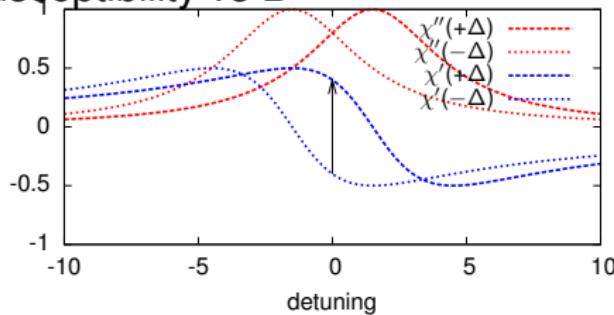


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

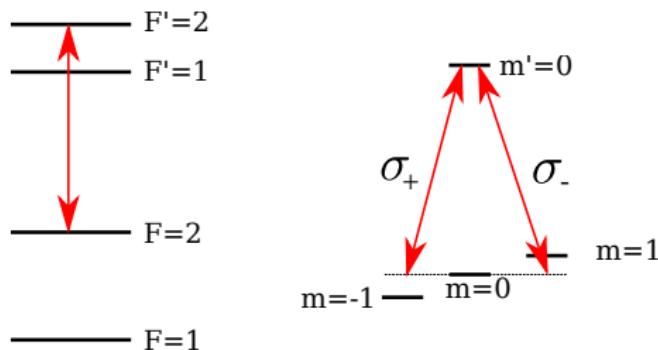


Susceptibility vs B

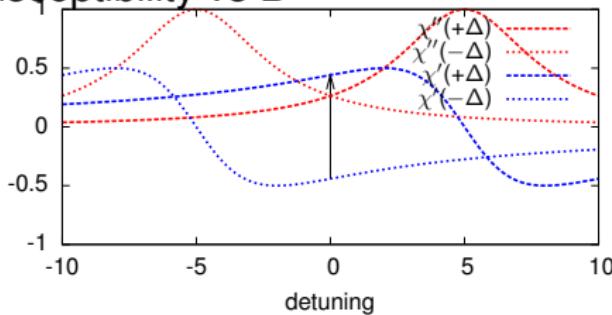


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

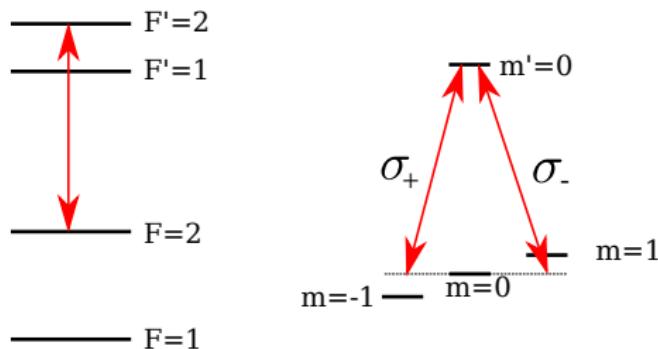


Susceptibility vs B

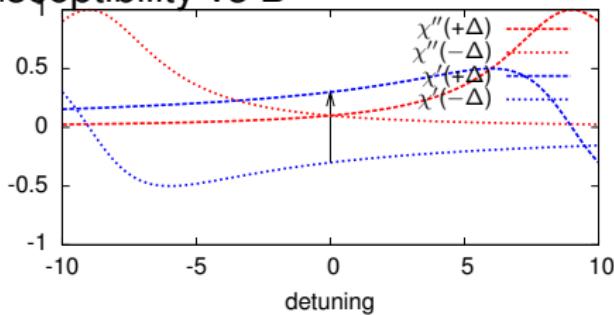


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

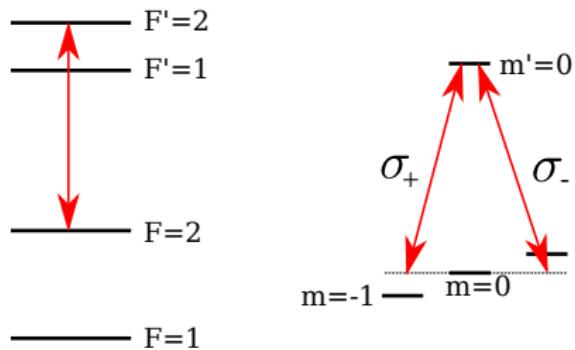


Susceptibility vs B

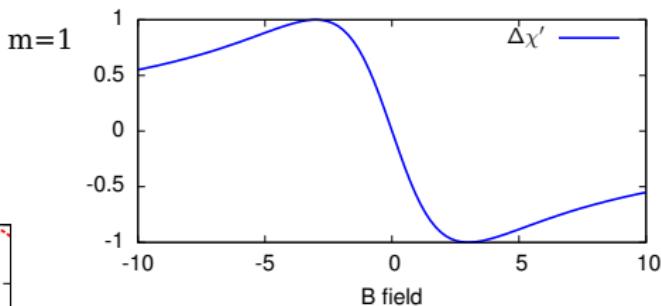


Optical magnetometer based on Faraday effect

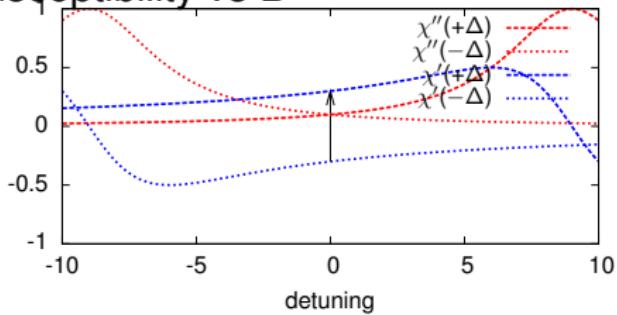
^{87}Rb D₁ line



Polarization rotation vs B



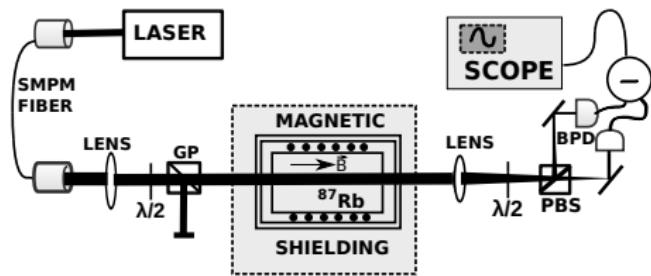
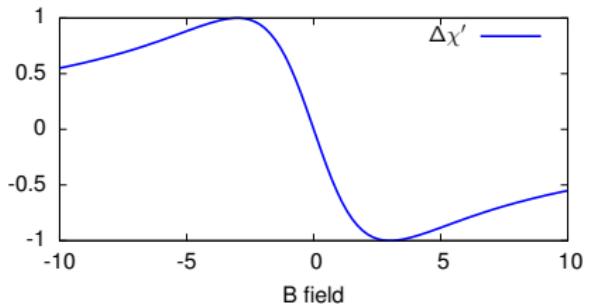
Susceptibility vs B



Optical magnetometer and non linear Faraday effect

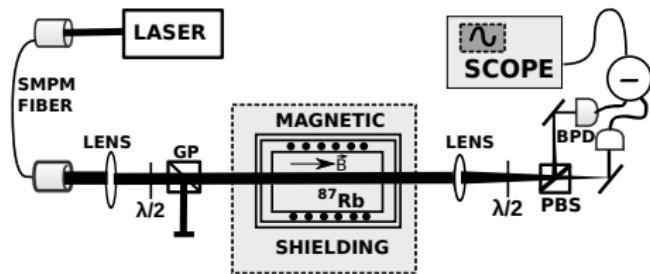
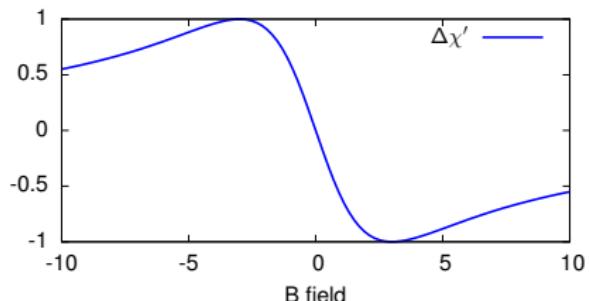
Naive model of rotation

Experiment

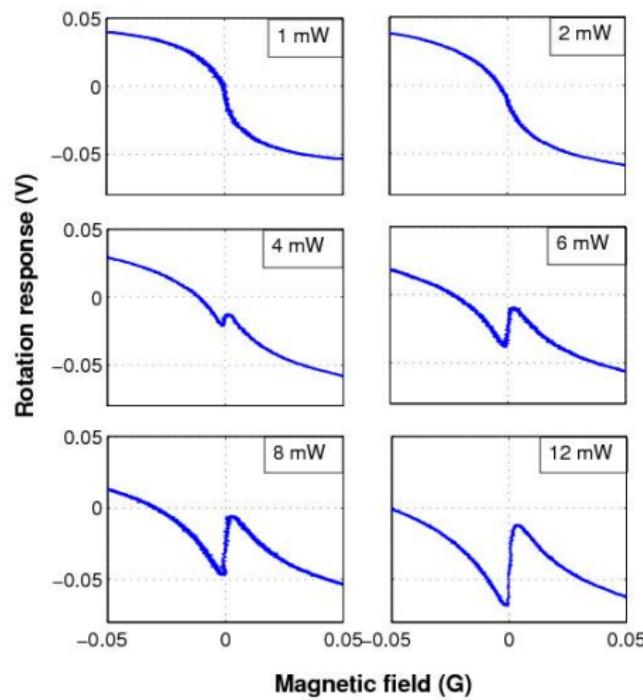


Optical magnetometer and non linear Faraday effect

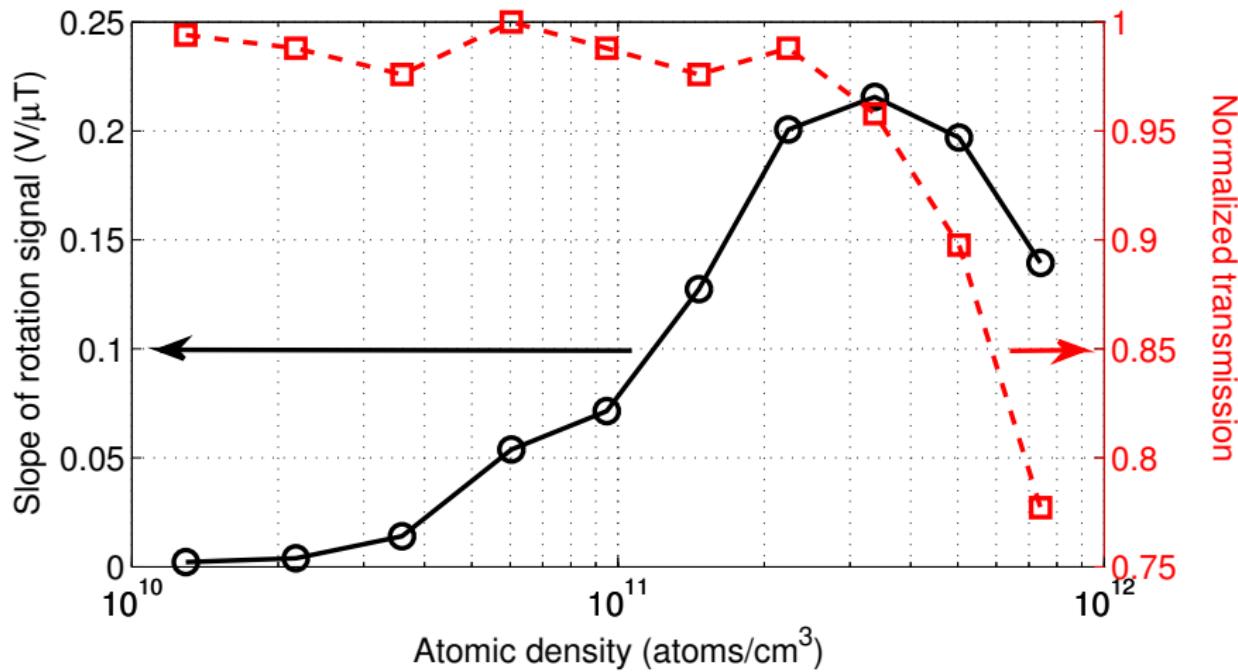
Naive model of rotation



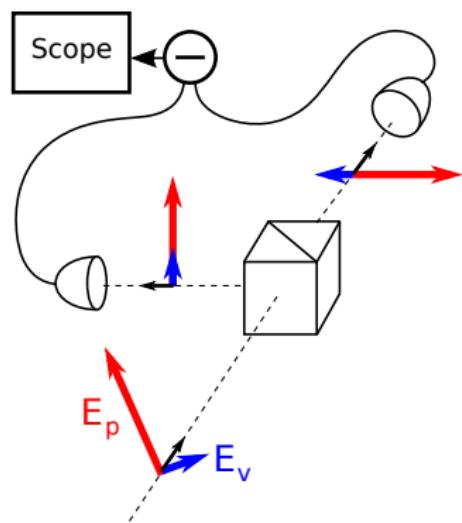
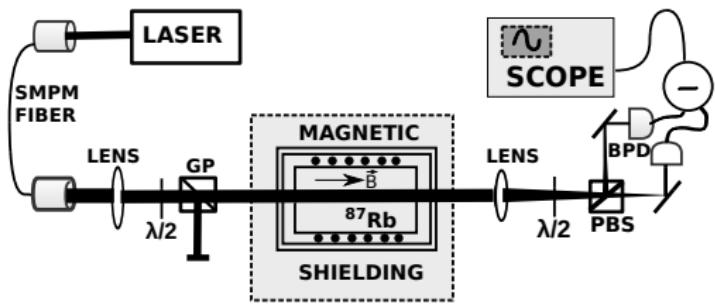
Experiment



Magnetometer response vs atomic density



Shot noise limit of the magnetometer

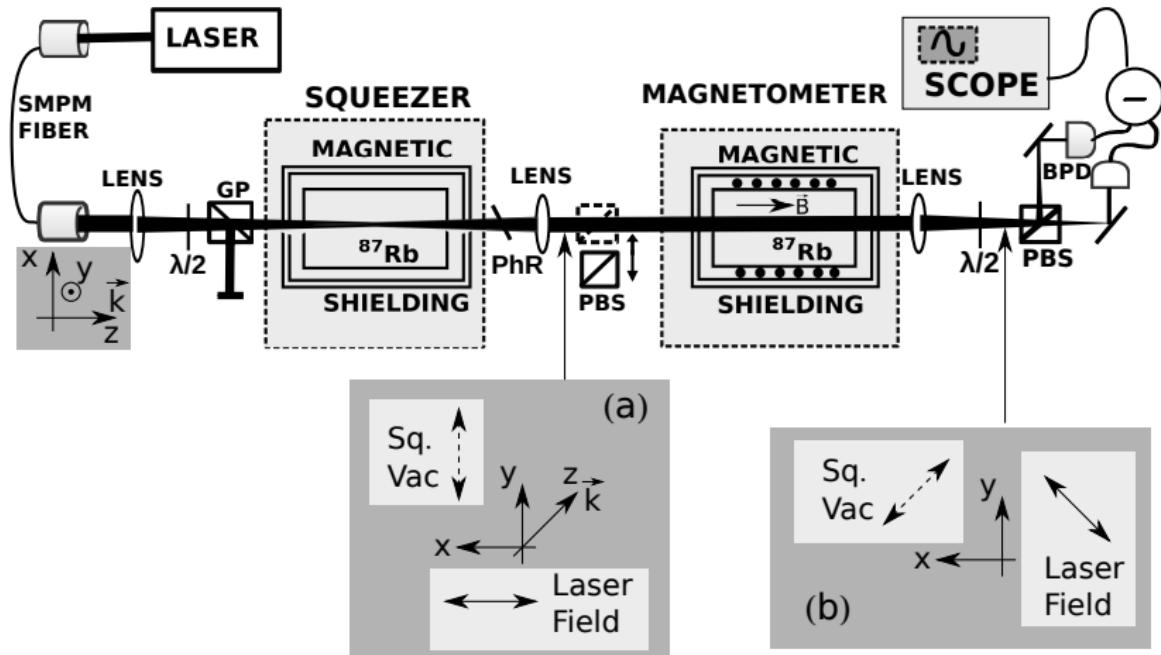


$$S = |E_p + E_v|^2 - |E_p - E_v|^2$$

$$S = 4E_p E_v$$

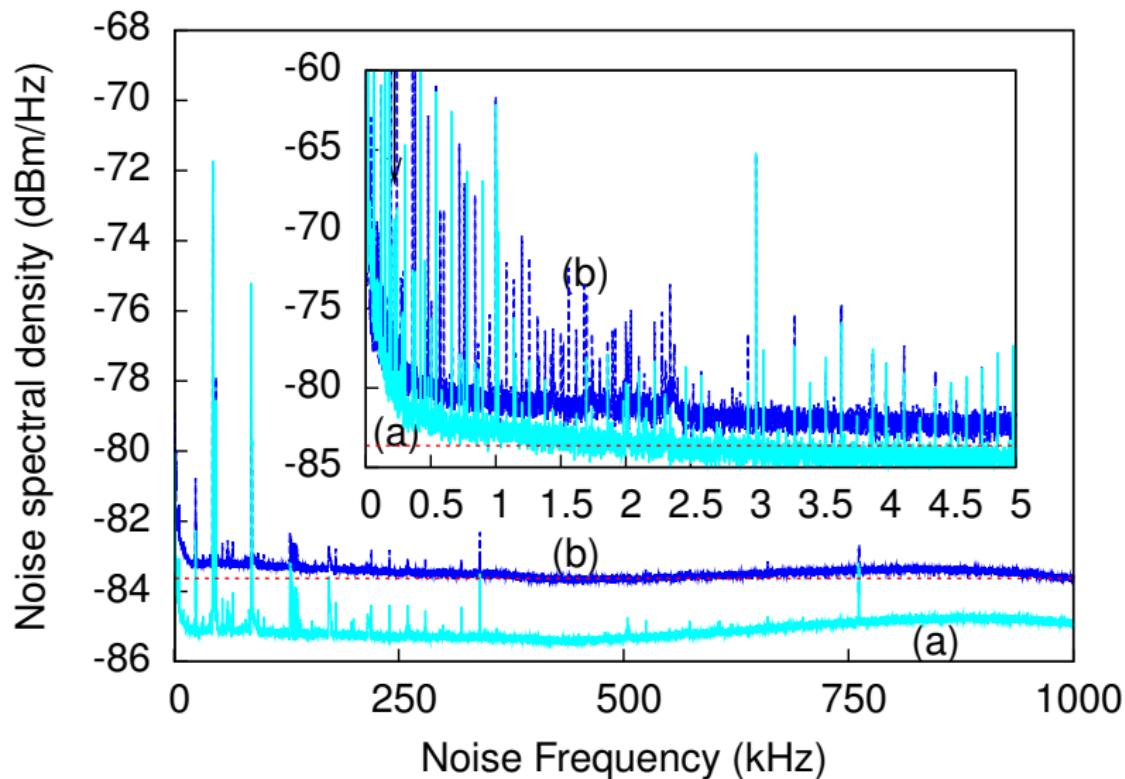
$$\langle \Delta S \rangle \sim E_p \langle \Delta E_v \rangle$$

Squeezed enhanced magnetometer setup

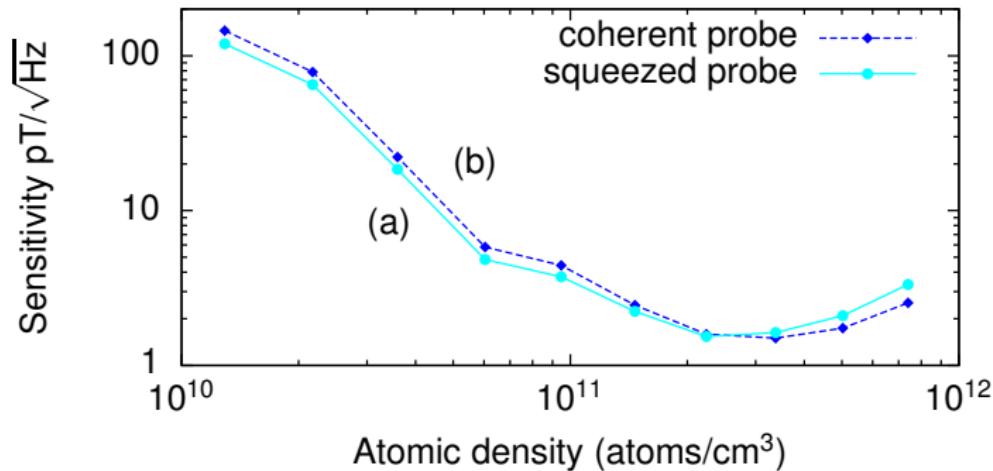
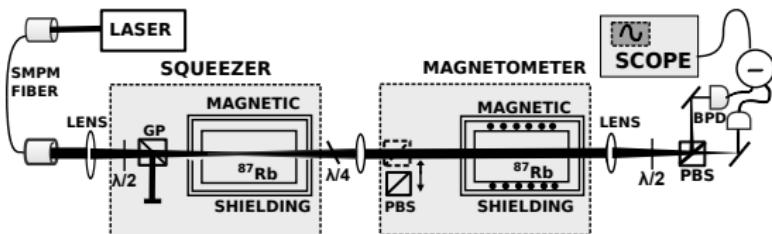


Note: Squeezed enhanced magnetometer was first demonstrated by Wolfgramm *et. al* Phys. Rev. Lett, **105**, 053601, 2010.

Magnetometer noise floor improvements

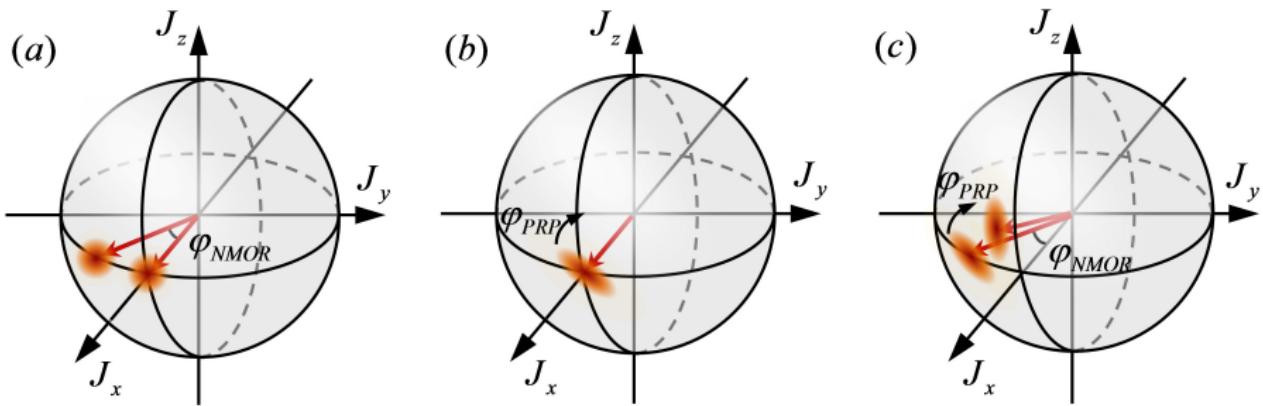
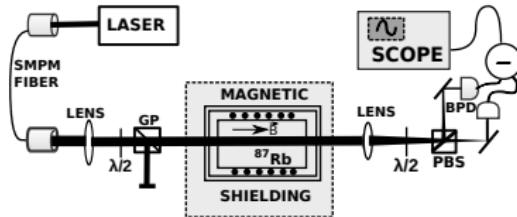


Magnetometer with squeezing enhancement



T. Horrom, et al. **PRA**, 86, 023803, (2012).

Self-squeezed magnetometry



Irina Novikova, Eugeniy E. Mikhailov, Yanhong Xiao, "Excess optical quantum noise in atomic sensors", Phys. Rev. A, **91**, 051804(R), (2015).