

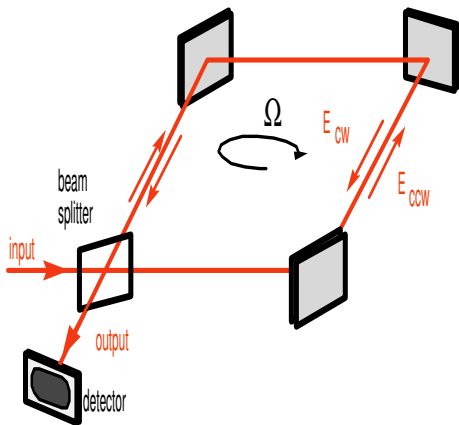
# Tuning laser intracavity dispersion: switching from vibration insensitive laser to the cavity length enhanced sensitivity.

Eugeniy E. Mikhailov, Savannah Cuozzo, Irina Novikova



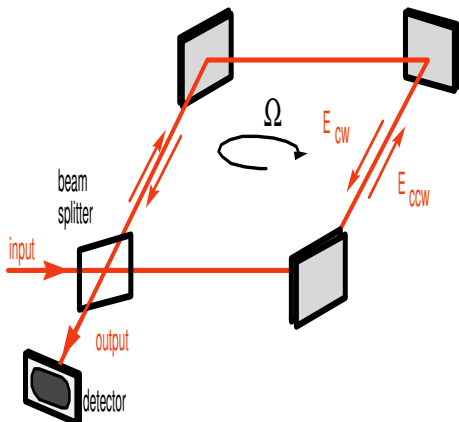
SPIE, 2 February 2019

# Dispersive cavity response



$$\Delta f = f_0 \frac{\Delta L}{L}$$

# Dispersive cavity response



$$\Delta f = f_0 \frac{\Delta L}{L} \frac{1}{n_g} = \Delta f_{empty} \frac{1}{n_g}$$

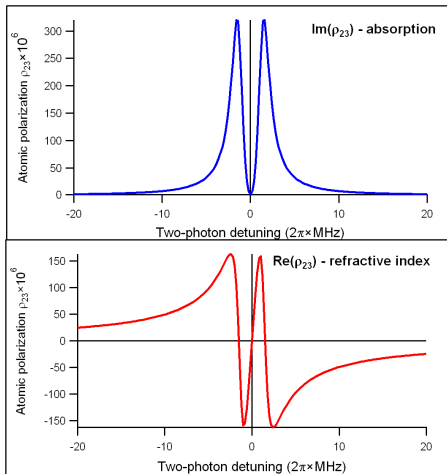
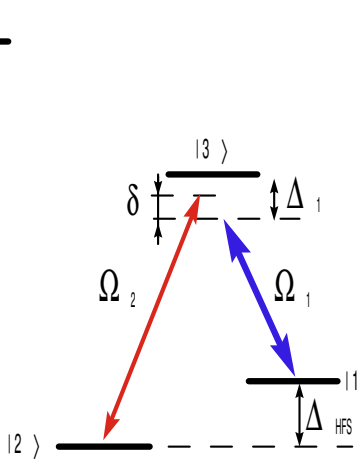
Group index

$$n_g(f) = n + f_0 \frac{\partial n}{\partial f}$$

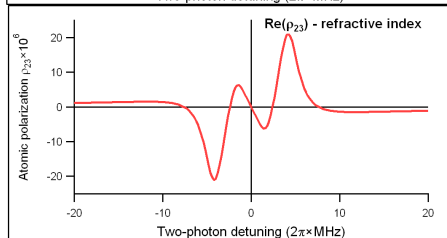
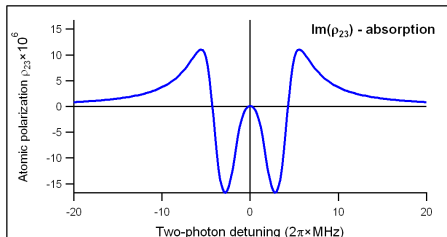
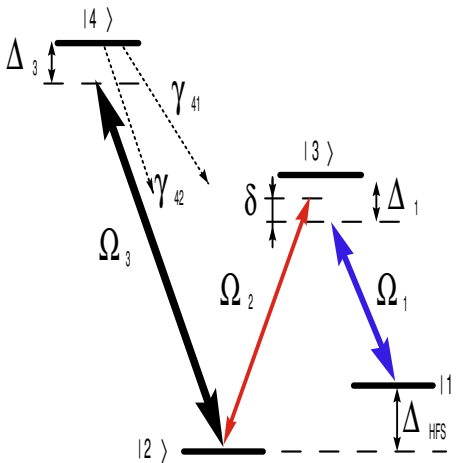
$$v_g = c/n_g$$

Cavity response enhanced if  $n_g < 1$  i.e. under the **fast light** condition  
Shahriar et al., PRA **75**, 053807 (2007)

# EIT - slow light

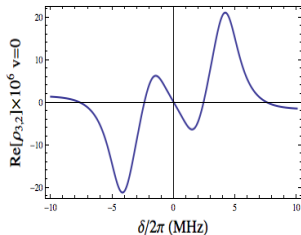


# N-bar with four-wave mixing - fast and with gain

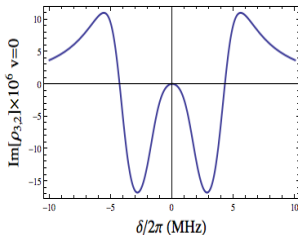


# N-bar with Doppler averaging

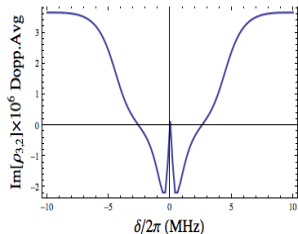
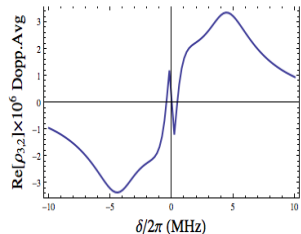
## Refractive index



## Absorption



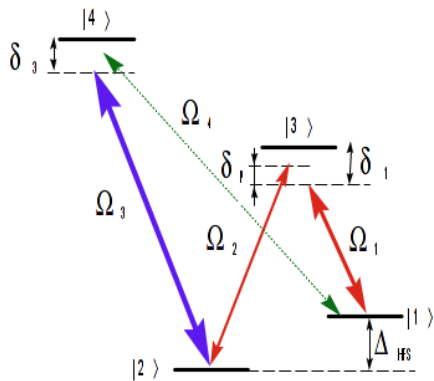
Stationary atoms



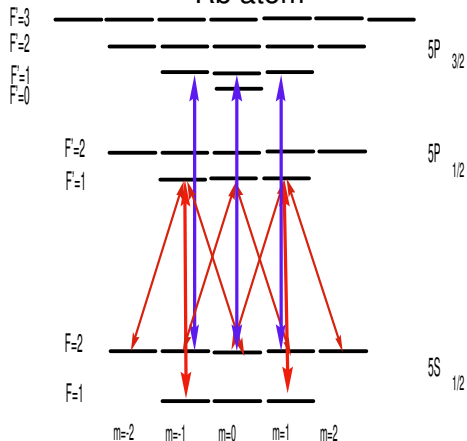
Room temperature  
Doppler averaged

# N-bar levels and fields diagram

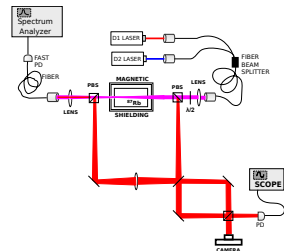
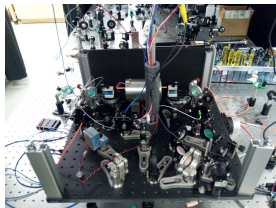
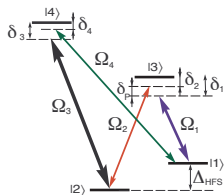
Artificial atom



$^{87}\text{Rb}$  atom



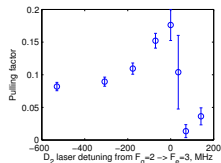
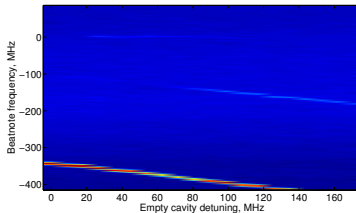
# Setup and measured pulling factor



$$P.F. = \frac{\Delta f_{\text{dispersive}}}{\Delta f_{\text{empty}}} = \frac{1}{n_g}$$

$$\Delta f_{\text{empty}} = f_0 \frac{\Delta L}{L}$$

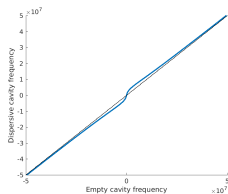
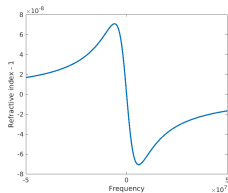
Gyro beatnote spectrum vs. empty cavity offset



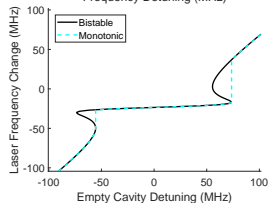
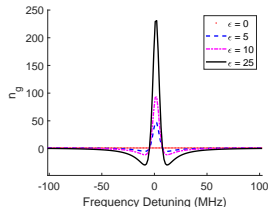
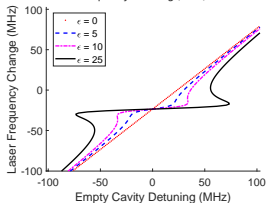
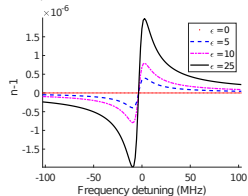


# Cavity response in fast, slow, and super slow regimes

Fast  
 $dn/d\omega < 1$



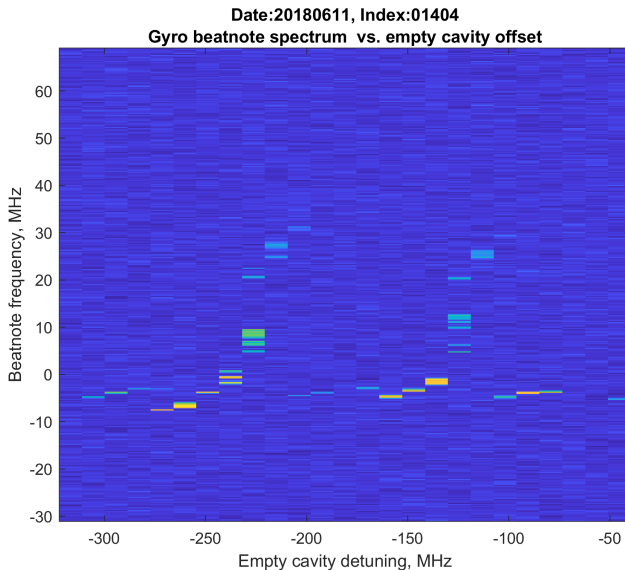
Slow  
 $dn/d\omega > 1$



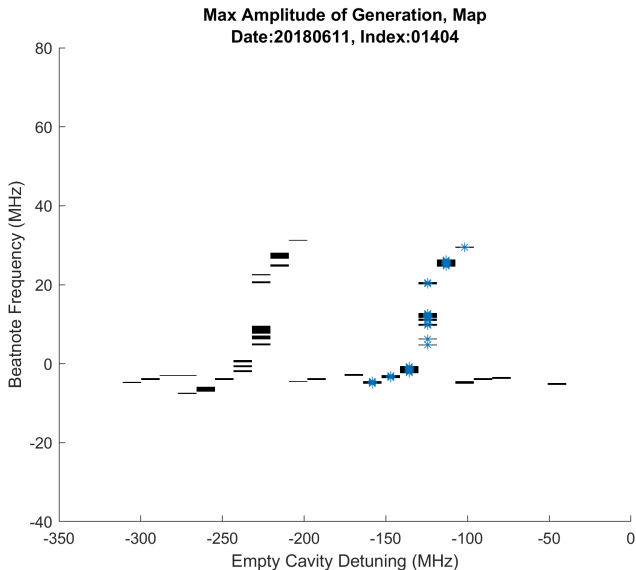
Lasing equation

$$n(\omega)L = m\lambda = mc \frac{2\pi}{\omega}$$

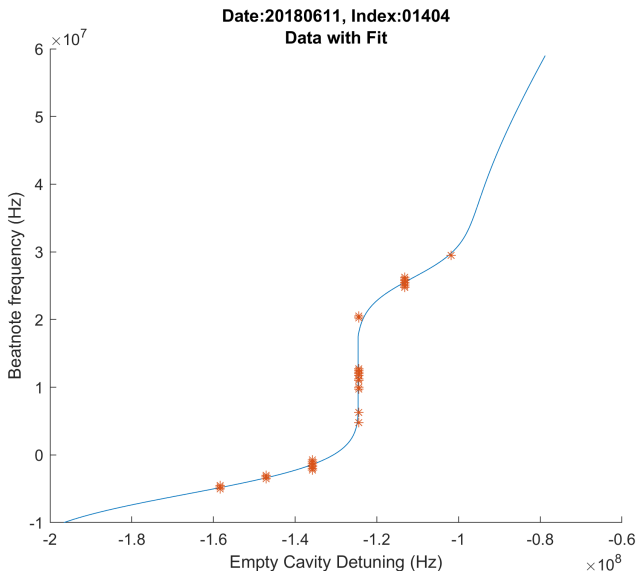
# Beatnote map with “high” pulling factor



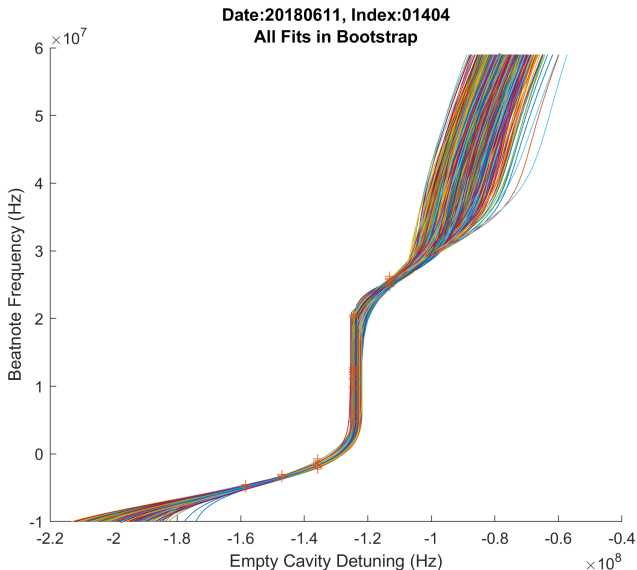
# Beatnote map with “high” pulling factor



# Beatnote map with “high” pulling factor

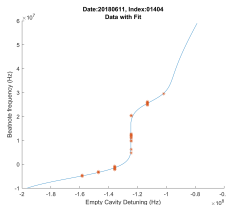


# Beatnote map with “high” pulling factor

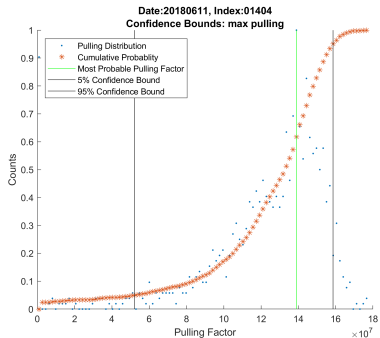
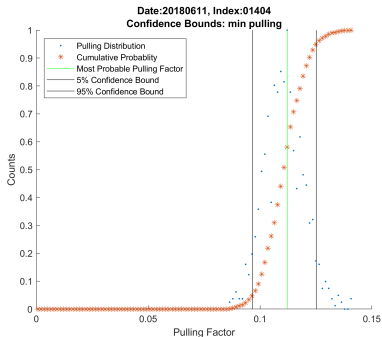


# Confidence in “high” and “low” pulling factors

Low PF= 0.112  
with 90% bounds  
(0.096, 0.125)

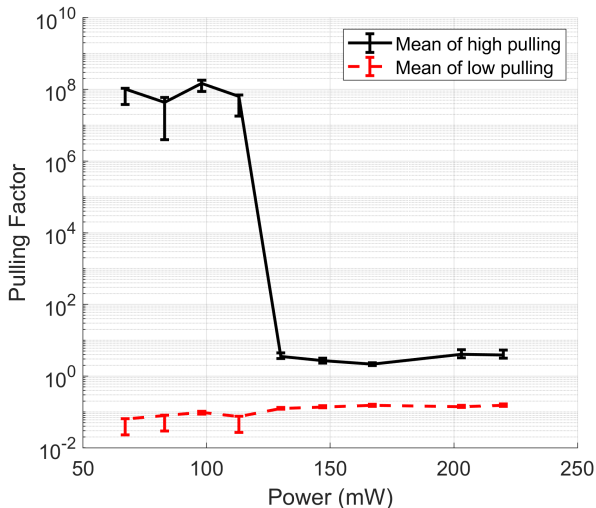


High PF=  $120 \times 10^6$   
with 90% bounds  
( $52 \times 10^6$ ,  $158 \times 10^6$ )

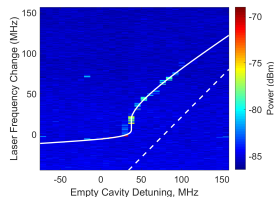


# Pulling factor vs power

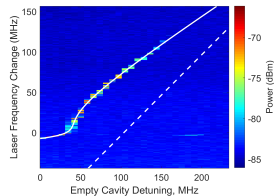
Temperature 100°C,  $\text{Power}_{D1}/\text{Power}_{D2} = 3/2$



## Power 98 mW

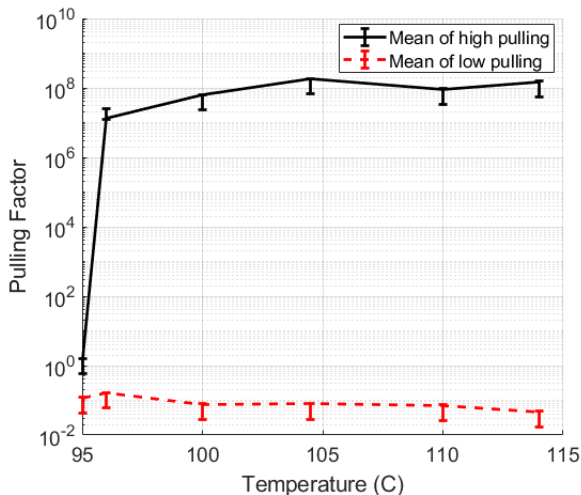


## Power 147 mW



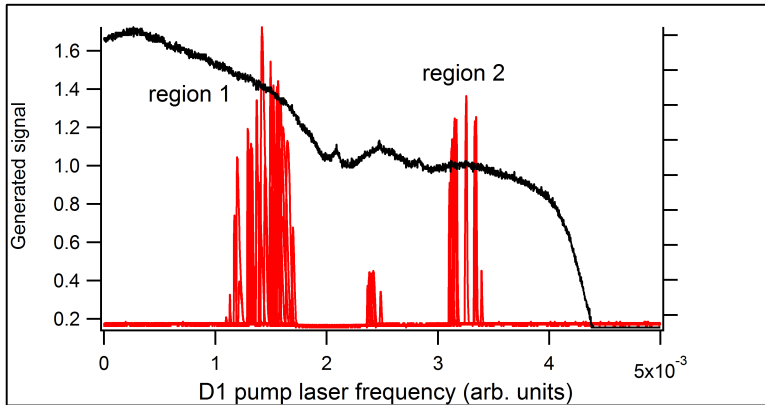
# Pulling factor vs temperature

Total power 82 mW,  $\text{Power}_{D1}/\text{Power}_{D2} = 2/1$



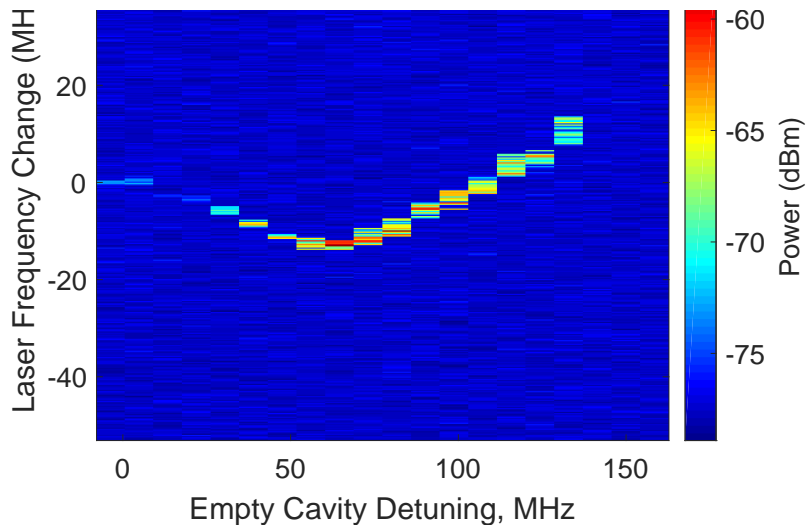


# Pulling factor vs detuning dependence

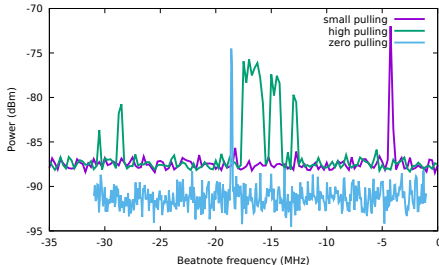
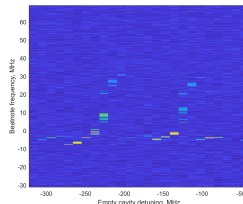
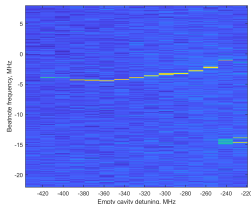


- Region 1: Pulling factor  $\leq 1$  (no discontinuities), high laser output
- Region 2: Large pulling  $\gg 1$
- Region 3 (middle): vibration free regime

# Laser insensitivity to cavity motion



# Beatnotes width comparison



Vibrations-  
broadened  
linewidth  
 $\sim 1/n_g = P.F.$

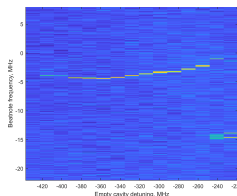
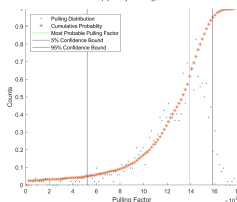
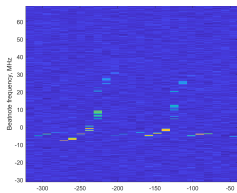
Quantum  
linewidth  
 $\sim 1/n_g^2 = (P.F.)^2$

C. Henry, IEEE Journal of Quantum Electronics 18, 259 (1982).

# People



# Summary



- We obtained enhanced cavity response as high  $10^8$  and potentially reached  $\infty$
- Under certain condition the laser output does not depend on cavity length, i.e. we have vibration insensitive laser

Savannah L. Cuozzo, Eugeny E. Mikhailov, "Dispersion enhanced laser frequency sensitivity and stability", arXiv:1812.08260, (2018).



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