

Squeezing vacuum with Rb atoms: quantum enhanced magnetometry and transient effects.

Eugeniy E. Mikhailov

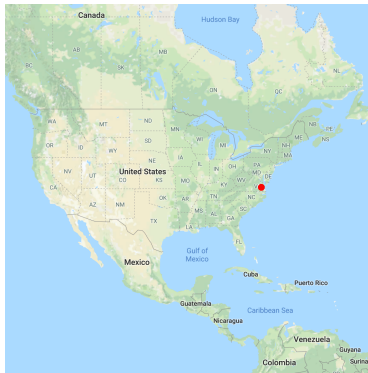


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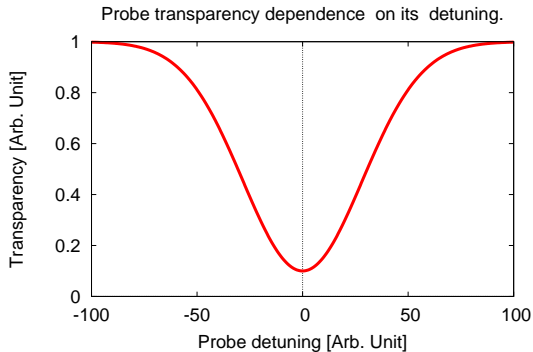
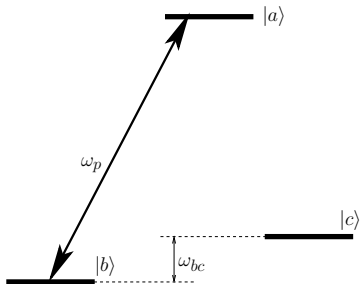
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Helmholtz-Institut Mainz, November 2, 2018

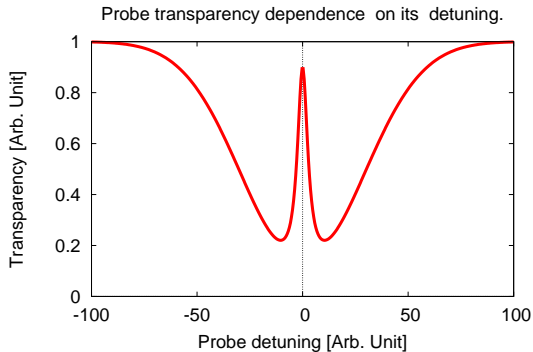
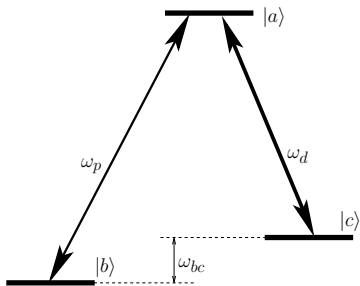
About College of William and Mary



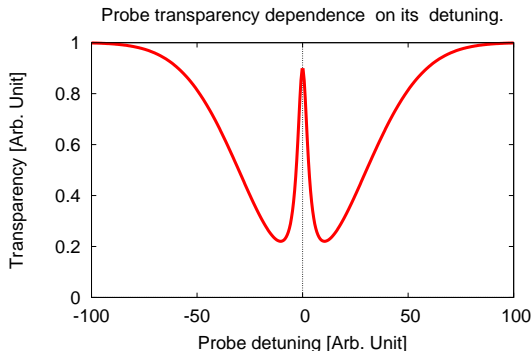
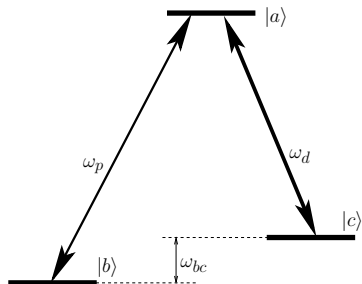
Electromagnetically Induced Transparency (EIT)



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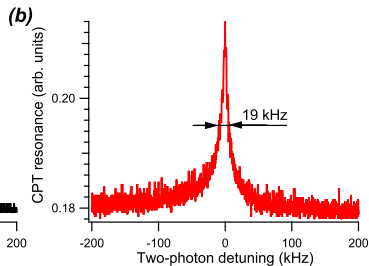
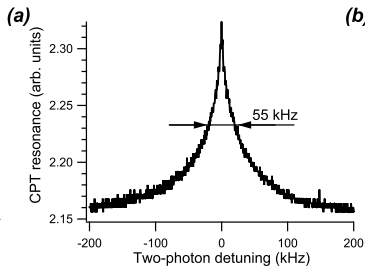
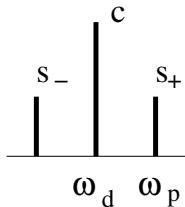
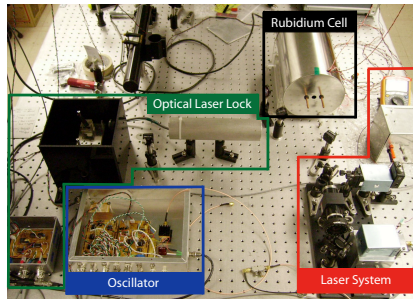
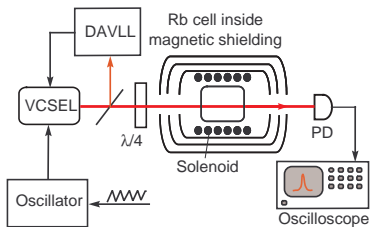
Electromagnetically Induced Transparency (EIT)



Coherent Population Trapping

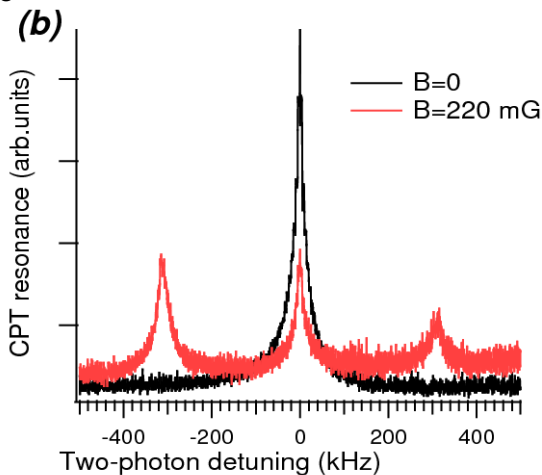
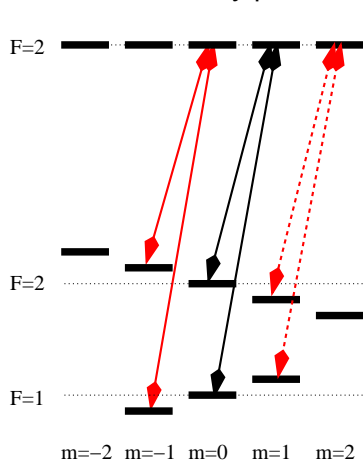
- Dark $|D\rangle = \Omega_d|b\rangle - \Omega_p|c\rangle$ and Bright $|B\rangle = \Omega_d|c\rangle + \Omega_p|b\rangle$ states
- resonance width ($\sim 10\text{kHz}$) much smaller than natural line width

EIT observation

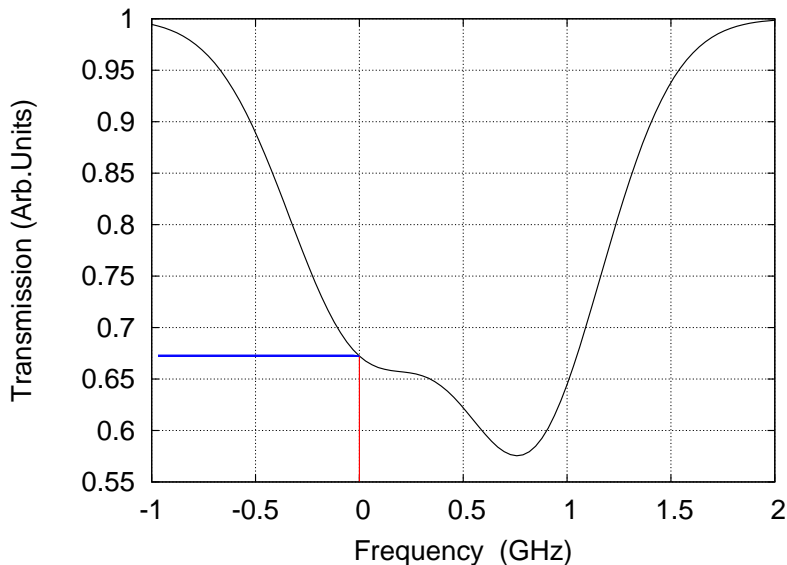


Simple EIT magnetometer

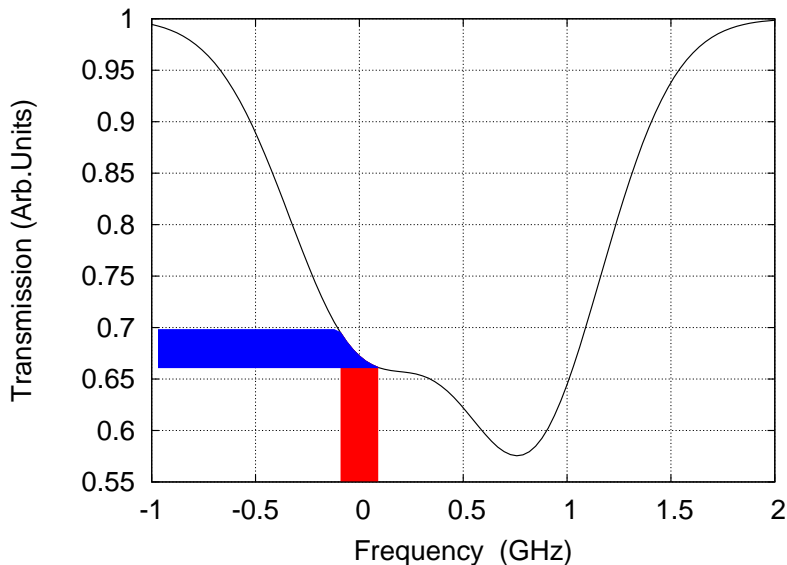
EIT with circularly polarized light



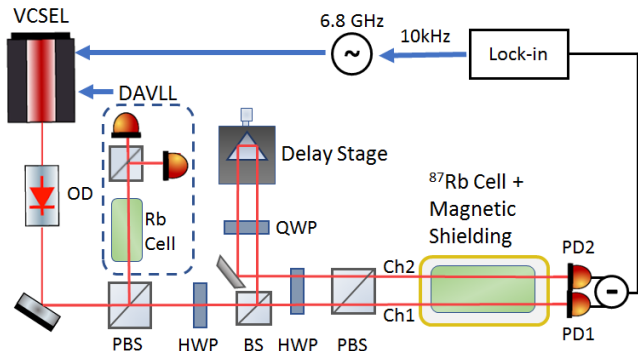
Phase noise to amplitude noise conversion



Phase noise to amplitude noise conversion

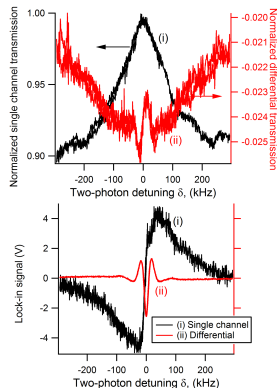
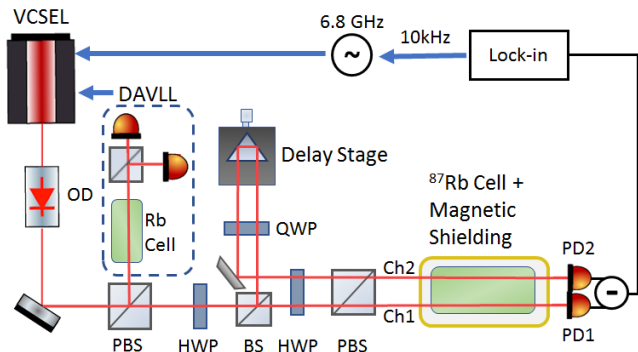


Transit Ramsey Electromagnetically Induced Transparency (TREIT)



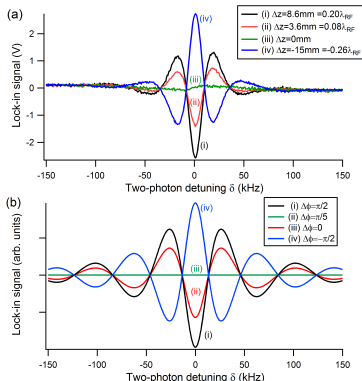
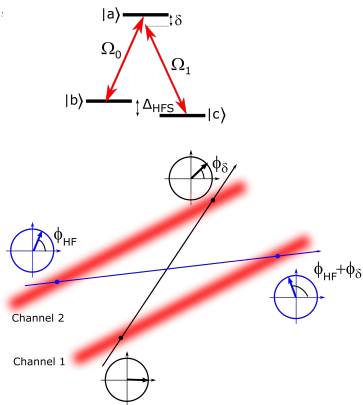
Ravn M. Jenkins, Eugeny E. Mikhailov, and Irina Novikova,
“Transit Ramsey EIT resonances in a Rb vacuum cell”,
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TREIT explanation

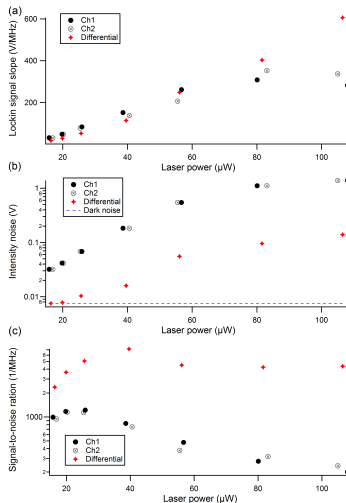
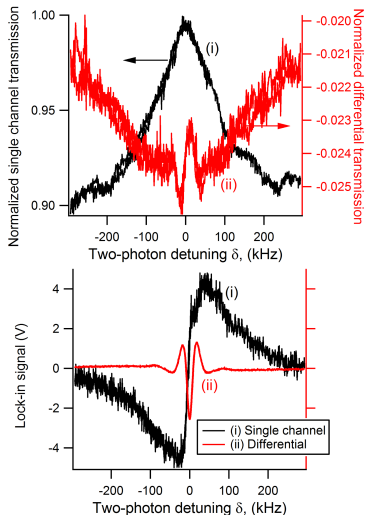


$$|D_1\rangle(t = \tau) = (\Omega_1|b\rangle - e^{i\delta \cdot \tau} \Omega_0|c\rangle)/\Omega,$$

$$|D_2\rangle(t = \tau) = (\Omega_1|b\rangle - e^{i\phi_{\text{HF}} + i\delta \cdot \tau} \Omega_0|c\rangle)/\Omega,$$

$$\Delta I(\delta) \propto \frac{|\Omega|^2}{\delta^2 + \Gamma^2} e^{-2\Gamma t_{\text{tr}}} \sin \phi_{\text{HF}} \times \sin [\delta(2t_{\text{tr}} + \tau) + \tan^{-1}(\delta/\Gamma)],$$

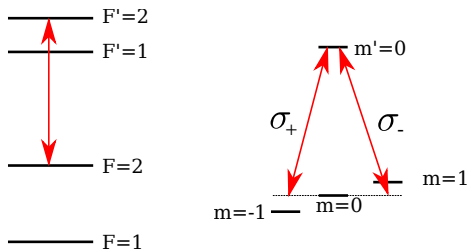
Signal to noise analysis EIT vs TREIT



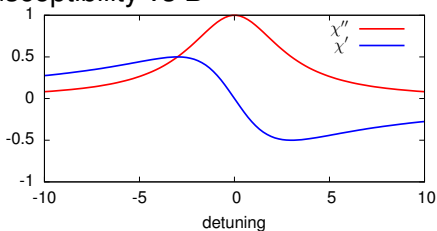
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Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

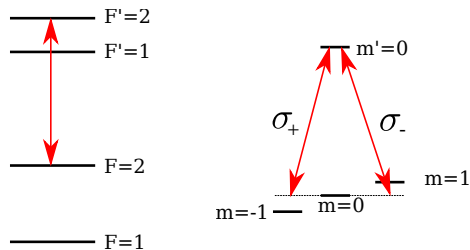


Susceptibility vs B

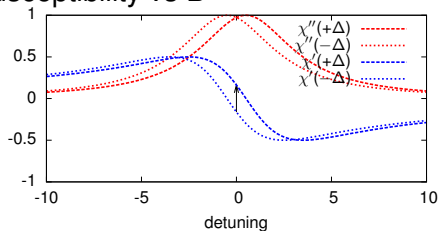


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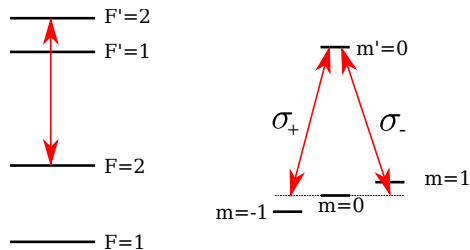


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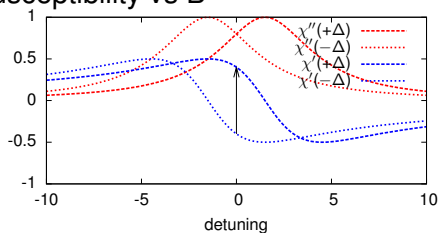


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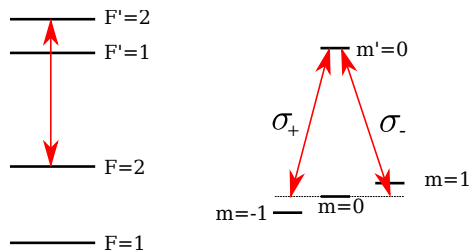


Susceptibility vs B

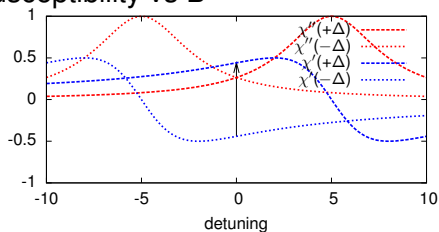


Optical magnetometer based on Faraday effect

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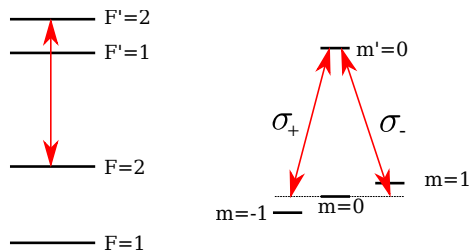


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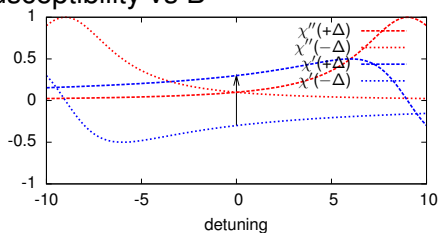


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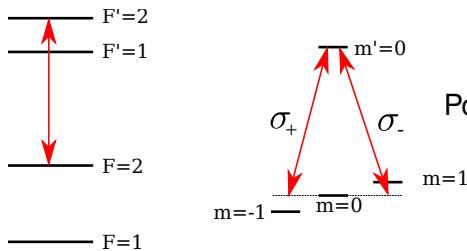


Susceptibility vs B

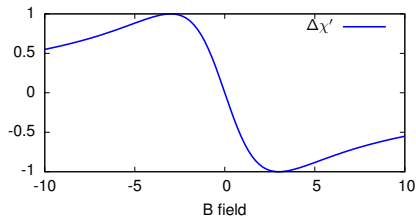


Optical magnetometer based on Faraday effect

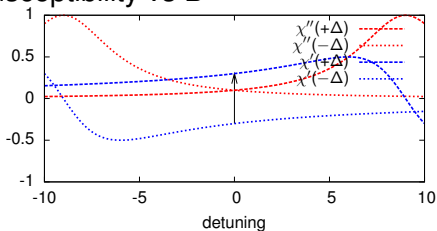
^{87}Rb D₁ line



Polarization rotation vs B

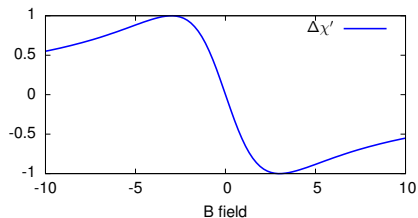


Susceptibility vs B

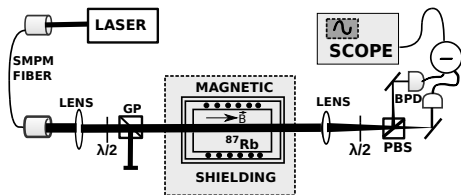


Optical magnetometer and non linear Faraday effect

Naive model of rotation

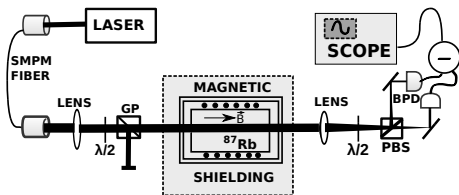
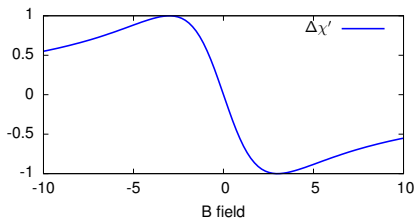


Experiment

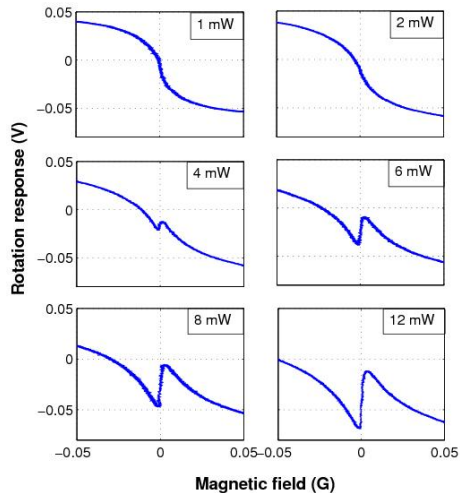


Optical magnetometer and non linear Faraday effect

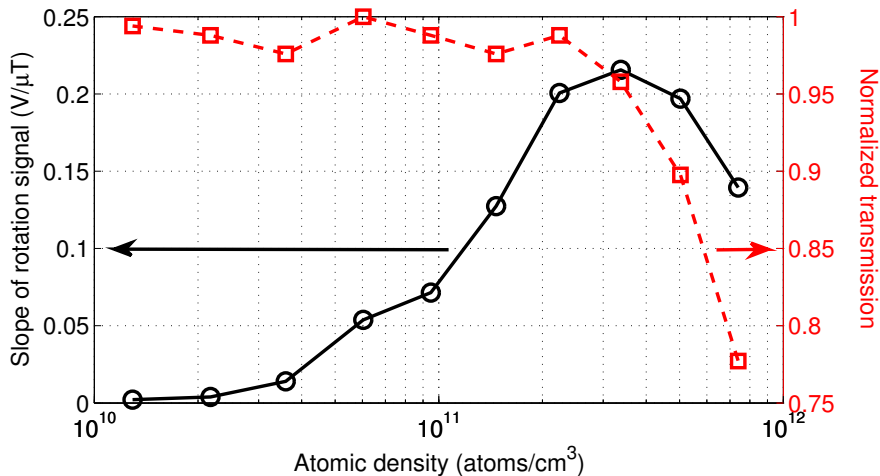
Naive model of rotation



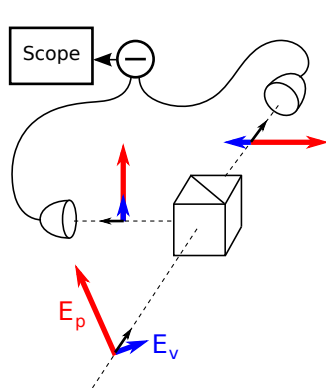
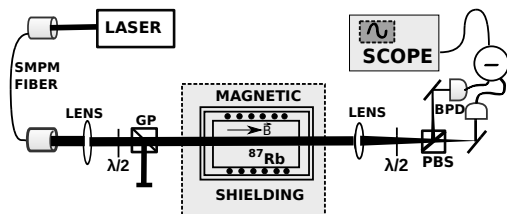
Experiment



Magnetometer response vs atomic density



Shot noise limit of the magnetometer



$$S = |E_p + E_v|^2 - |E_p - E_v|^2$$

$$S = 4E_p E_v$$

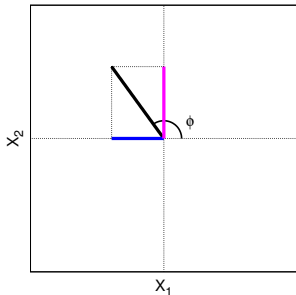
$$\langle \Delta S \rangle \sim E_p \langle \Delta E_v \rangle$$

Transition from classical to quantum field

Classical analog

- Field amplitude a
- Field real part
 $X_1 = (a^* + a)/2$
- Field imaginary part
 $X_2 = i(a^* - a)/2$

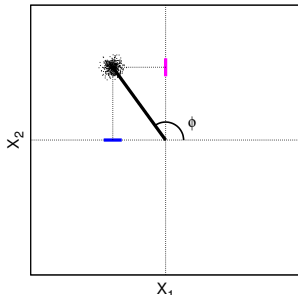
$$E(\phi) = |a|e^{-i\phi} = X_1 + iX_2$$



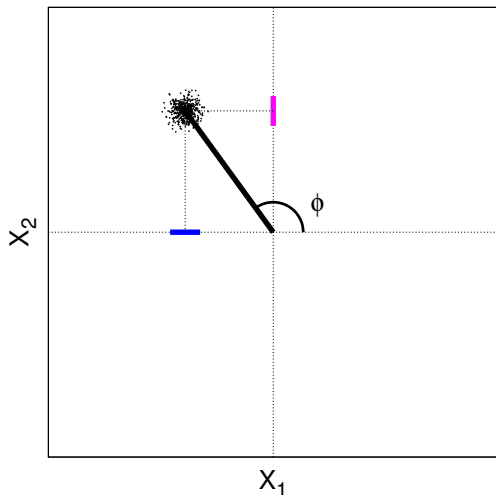
Quantum approach

- Field operator \hat{a}
- Amplitude quadrature
 $\hat{X}_1 = (\hat{a}^\dagger + \hat{a})/2$
- Phase quadrature
 $\hat{X}_2 = i(\hat{a}^\dagger - \hat{a})/2$

$$\hat{E}(\phi) = \hat{X}_1 + i\hat{X}_2$$



Quantum optics summary



Light consist of photons

- $\hat{N} = a^\dagger a$

Commutator relationship

- $[a, a^\dagger] = 1$

- $[X_1, X_2] = i/2$

Detectors measure

- number of photons \hat{N}
- Quadratures \hat{X}_1 and \hat{X}_2

Uncertainty relationship

- $\Delta X_1 \Delta X_2 \geq 1/4$

Heisenberg uncertainty principle and its optics equivalent

Heisenberg uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa



Heisenberg uncertainty principle and its optics equivalent



Heisenberg uncertainty principle

$$\Delta p \Delta x \geq \hbar/2$$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa

Optics equivalent

$$\Delta \phi \Delta N \geq 1$$

The more precisely the PHASE is determined, the less precisely the AMPLITUDE is known, and vice versa

Heisenberg uncertainty principle and its optics equivalent



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$$\Delta p \Delta x \geq \hbar/2$$

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Optics equivalent

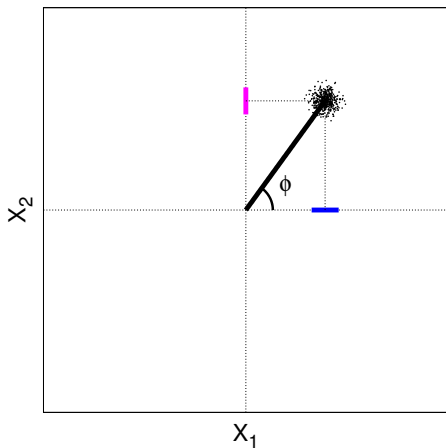
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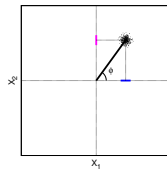
Optics equivalent strict definition

$$\Delta X_1 \Delta X_2 \geq 1/4$$

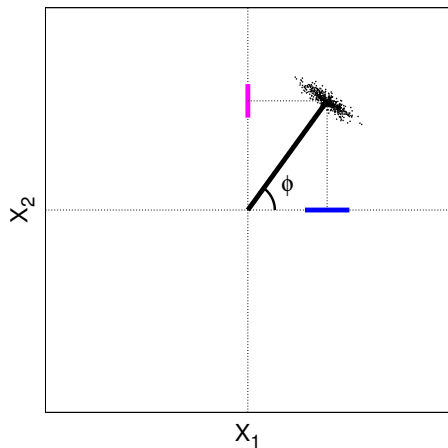
Squeezed quantum states zoo



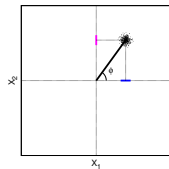
Unsqueezed
coherent



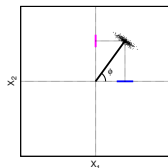
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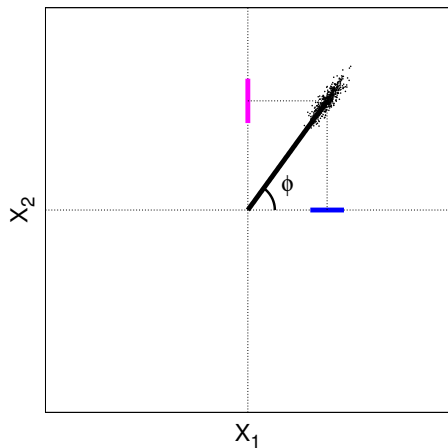
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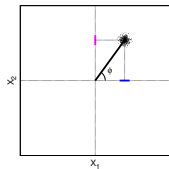
Amplitude
squeezed



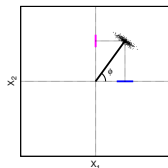
Squeezed quantum states zoo



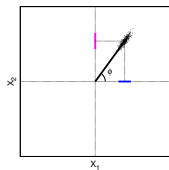
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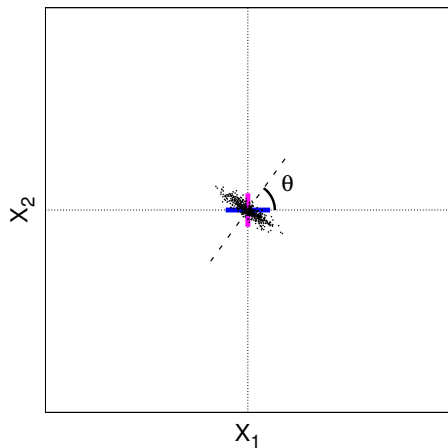
Amplitude
squeezed



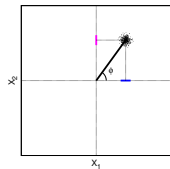
Phase
squeezed



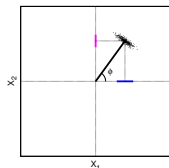
Squeezed quantum states zoo



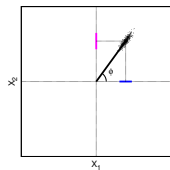
Unsqueezed
coherent



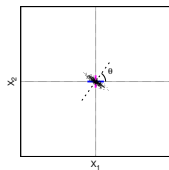
Amplitude
squeezed



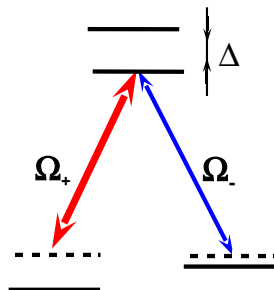
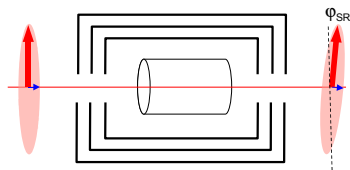
Phase
squeezed



Vacuum
squeezed



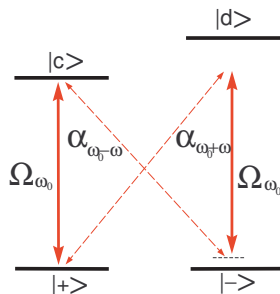
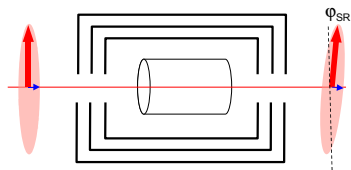
Self-rotation of elliptical polarization in atomic medium



A. B. Matsko, I. Novikova, G. R. Welch, D. Budker, D. F. Kimball, and S. M. Rochester, PRA 66, 043815 (2002):
theoretical prediction of 4–6 dB noise suppression

$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^\dagger - a_{in})$$

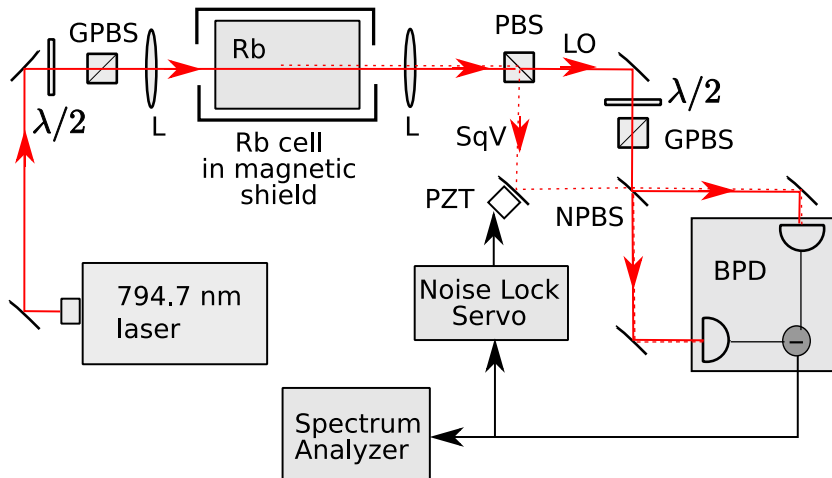
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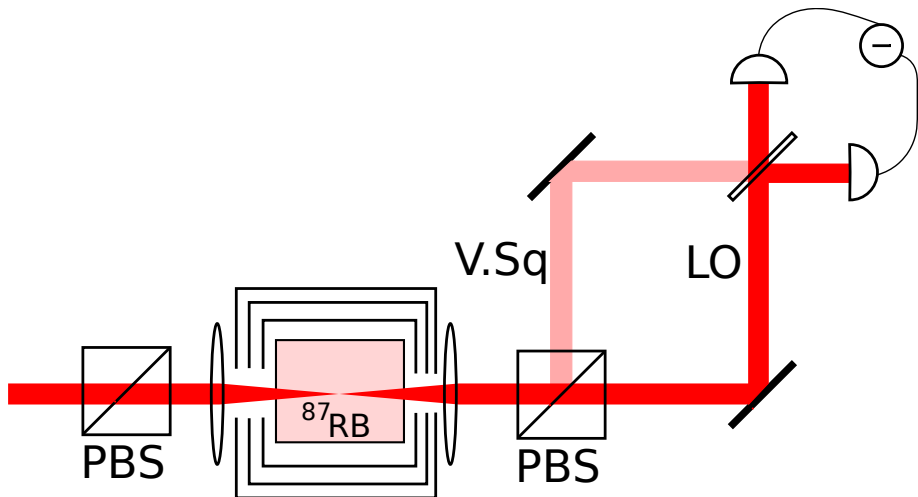
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$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^\dagger - a_{in})$$

Setup



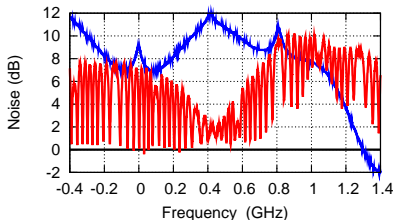
Setup



Noise contrast vs detuning in hot ^{87}Rb vacuum cell

$$F_g = 2 \rightarrow F_e = 1, 2$$

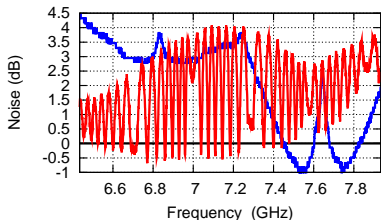
Noise vs detuning



Transmission — PSR noise

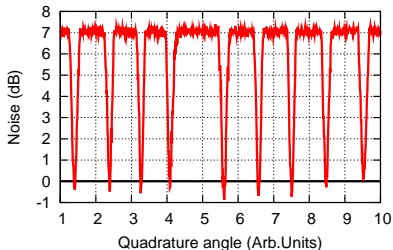
$$F_g = 1 \rightarrow F_e = 1, 2$$

Noise vs detuning

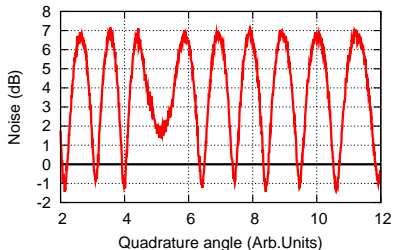


Transmission — PSR noise

Noise vs quadrature angle

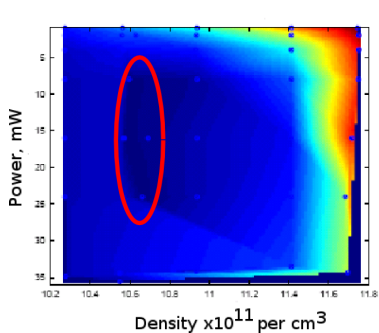


Noise vs quadrature angle

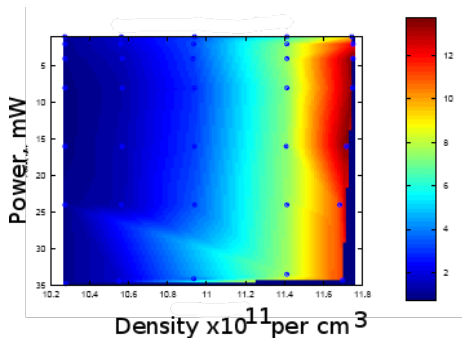


Squeezing region

Squeezing



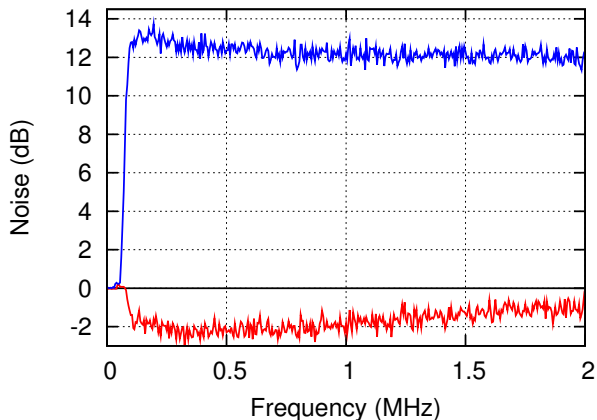
Anti-squeezing



Observation of reduction of quantum noise below the shot noise limit is corrupted by the excess noise due to atomic interaction with atoms.

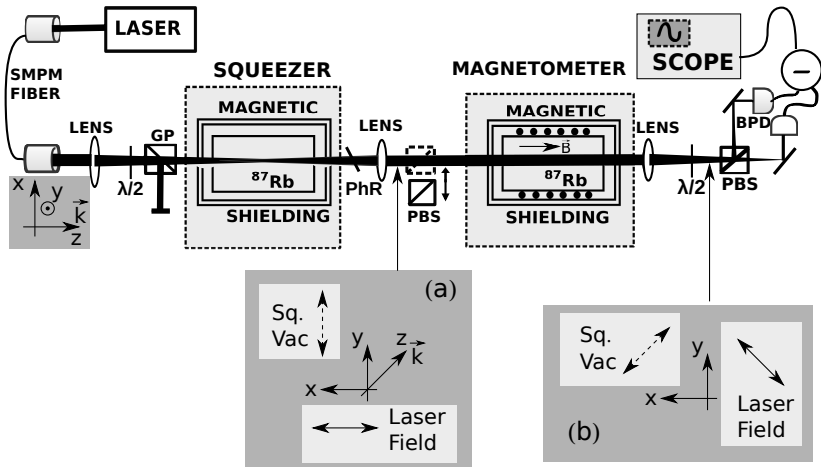
Maximally squeezed spectrum with ^{87}Rb

W&M team. ^{87}Rb $F_g = 2 \rightarrow F_e = 2$, laser power 7 mW, $T=65^\circ\text{C}$



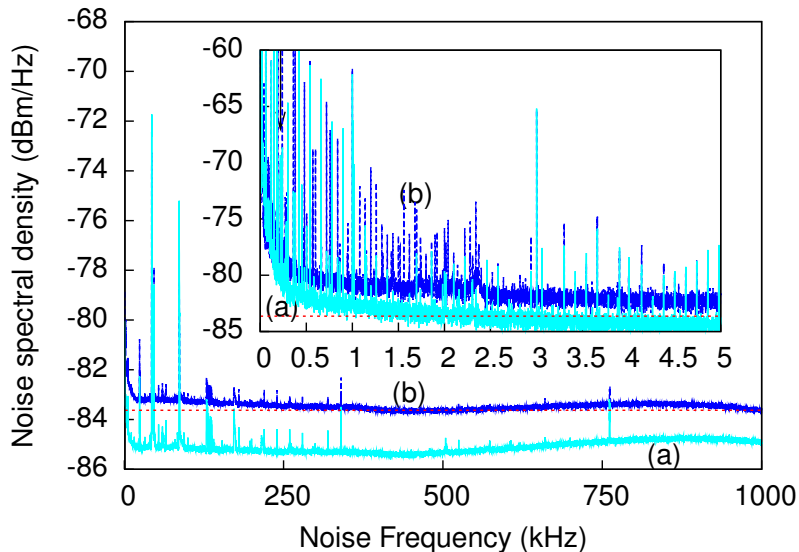
Lezama et.al report 3 dB squeezing in similar setup
Phys. Rev. A 84, 033851 (2011)

Squeezed enhanced magnetometer setup

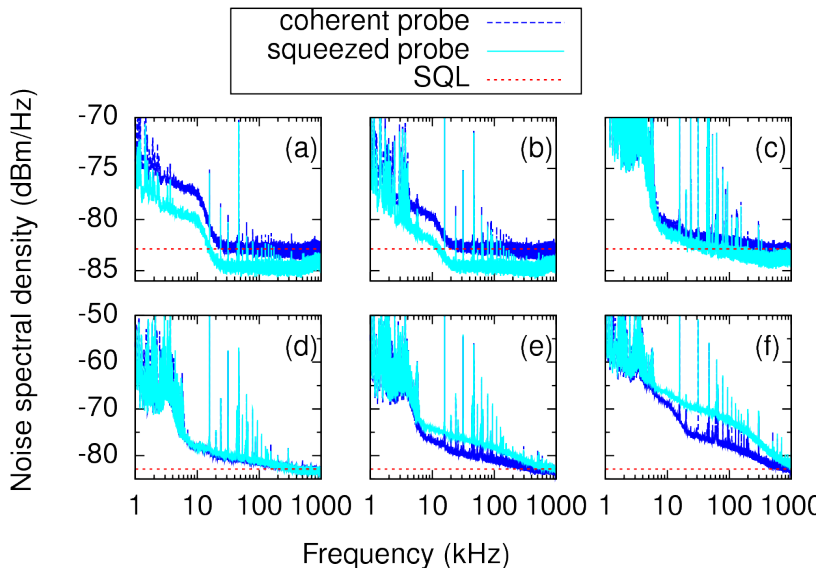


Note: Squeezed enhanced magnetometer was first demonstrated by Wolfgramm *et. al*/ Phys. Rev. Lett, **105**, 053601, 2010.

Magnetometer noise floor improvements

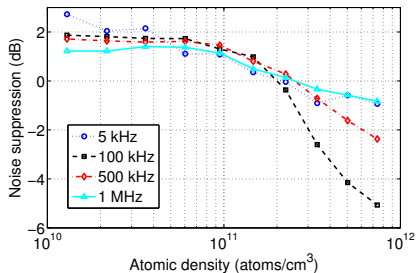


Magnetometer noise spectra

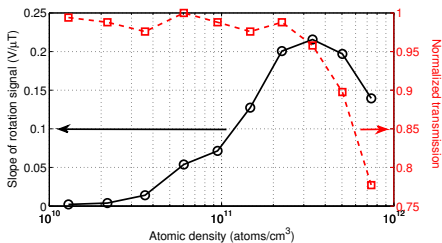


Noise suppression and response vs atomic density

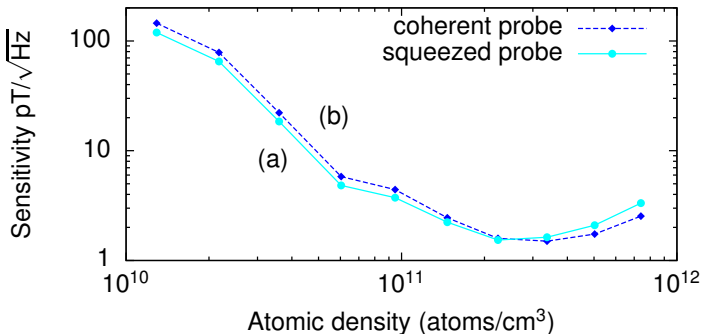
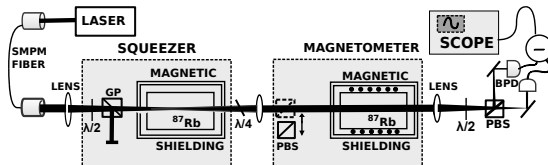
Noise suppression



Response

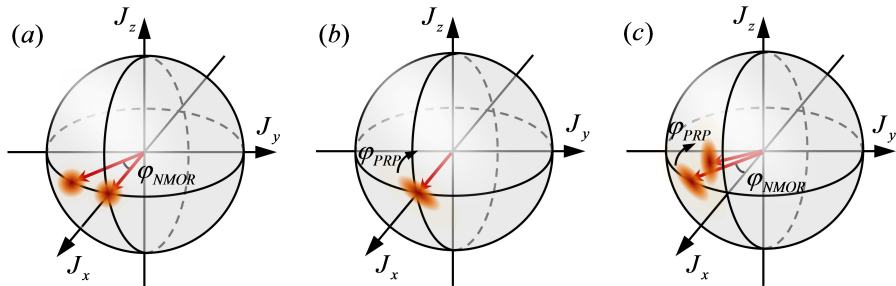
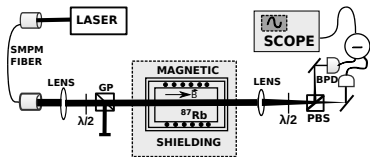


Magnetometer with squeezing enhancement



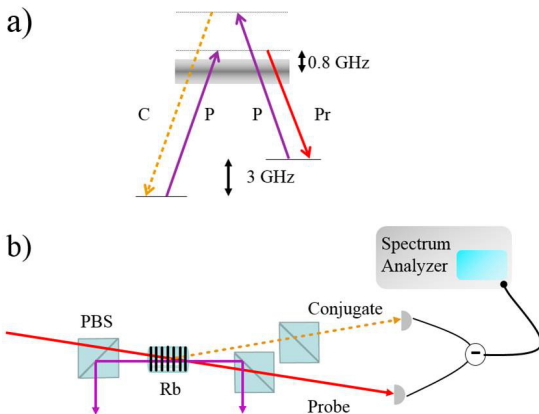
T. Horrom, et al. **PRA**, 86, 023803, (2012).

Self-squeezed magnetometry



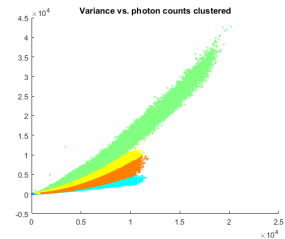
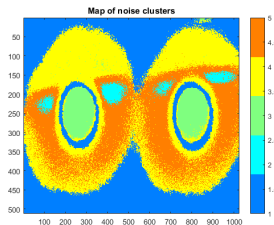
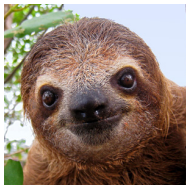
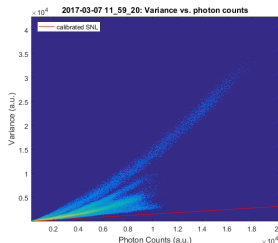
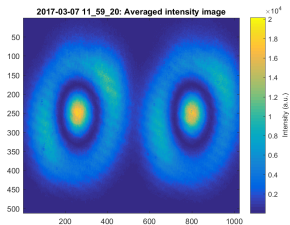
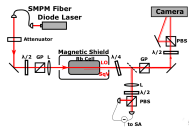
Irina Novikova, Eugeny E. Mikhailov, Yanhong Xiao, "Excess optical quantum noise in atomic sensors", Phys. Rev. A, **91**, 051804(R), (2015)

20 pT/ $\sqrt{\text{Hz}}$ self-squeezed magnetometry with 4WM



N. Otterstrom, R. C. Pooser, and B. J. Lawrie, "Nonlinear optical magnetometry with accessible in situ optical squeezing", *Optics Letters*, **39**, Issue 22, pp. 6533-6536 (2014)

Quantum imaging effort: from owl to sloth



People

WM: Irina Novikova, Ravn M. Jenkins, Savannah Cuozzo, Austin Kalasky

Former members: George Denny, Melissa A. Guidry, Mi Zhang, Travis Horrom, Gleb Romanov, Demetrious Kutzke



Summary

- fully atomic squeezed enhanced magnetometer with sensitivity as low as $1 \text{ pT}/\sqrt{\text{Hz}}$
- superluminal squeezing propagation with $v_g \approx -7'000 \text{ m/s} \approx -c/43'000$ or time advancement of $11 \text{ } \mu\text{S}$
- We were able to improve squeezing by multipass configuration
- Our squeezed state is a set of competing multimodes
- We are working on quantum modes extraction and imaging

Financial support by AFOSR, ARO, and NSF.