Quantum optics and squeezed states of light

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From ray optics to semiclassical optics

Classical/Geometrical optics

- light is a ray
- which propagates straight
- cannot explain diffraction and interference

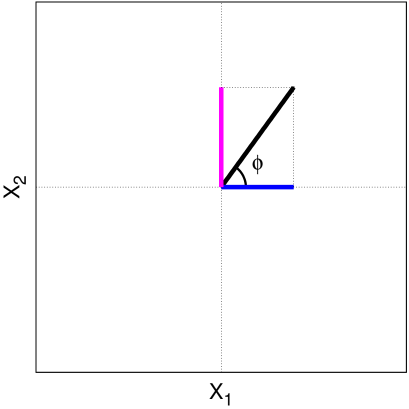
Semiclassical optics

- light is a wave
- color (wavelength/frequency) is important
- amplitude (a) and phase are important, $E(t) = ae^{i(kz-\omega t)}$
- cannot explain residual measurements noise

Classical field

$$E(\phi) = |a|e^{-i\phi} = |a|\cos(\phi) + i|a|\sin(\phi) = X_1 + iX_2, \quad \phi = \omega t - kz$$

Detectors sense the real part of the field (X_1) but there is a way to see X_2 as well



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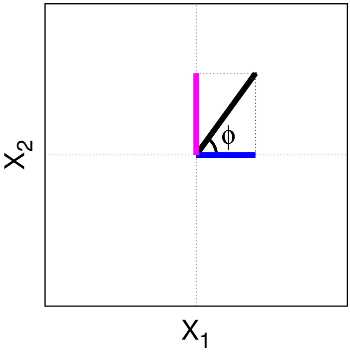

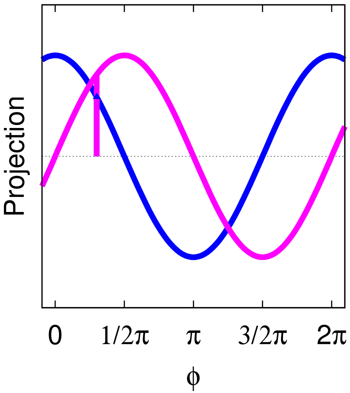

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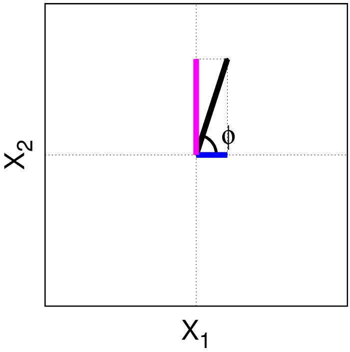

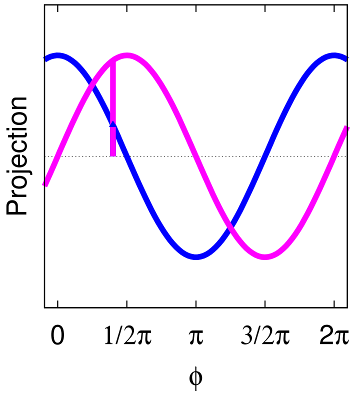

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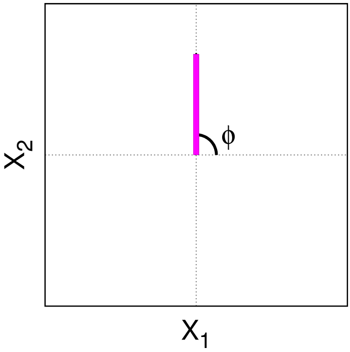

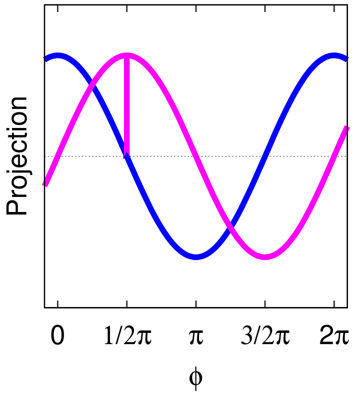

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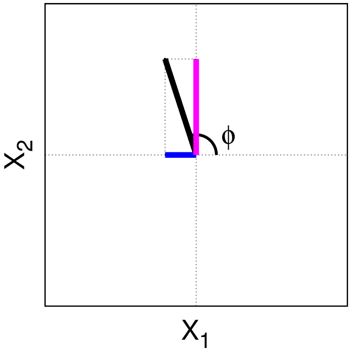

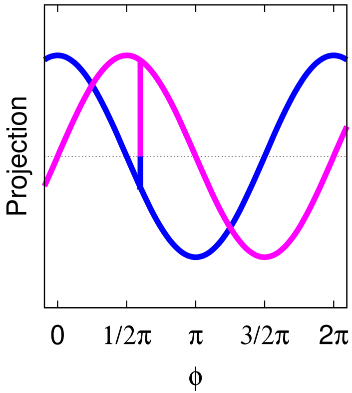

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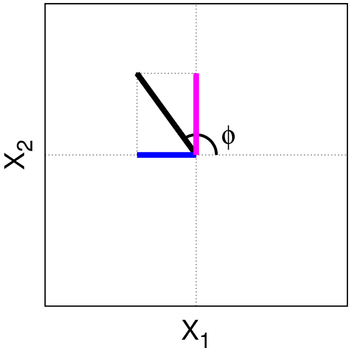

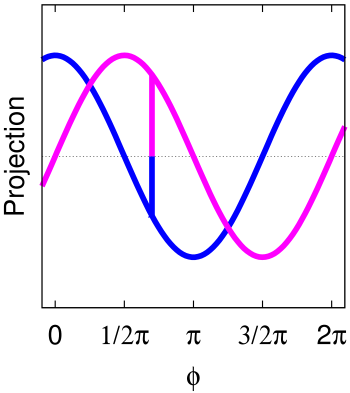

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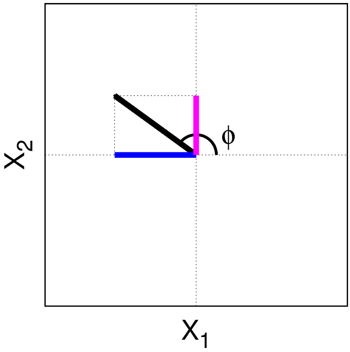

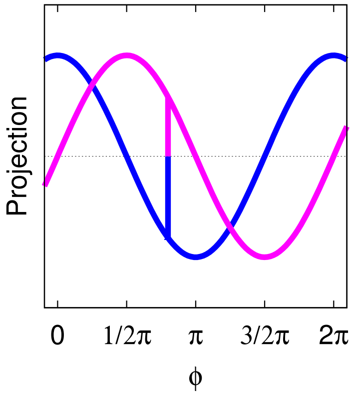

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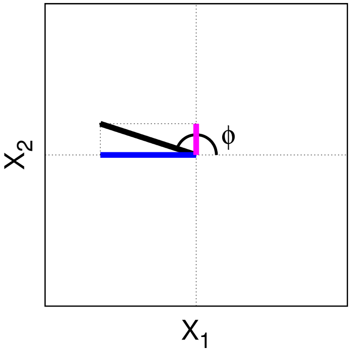

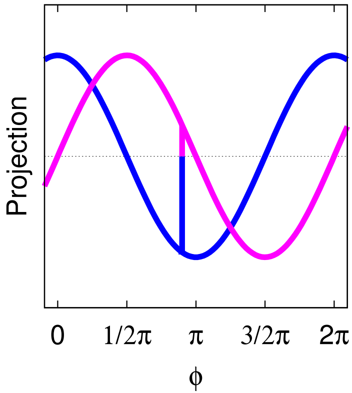

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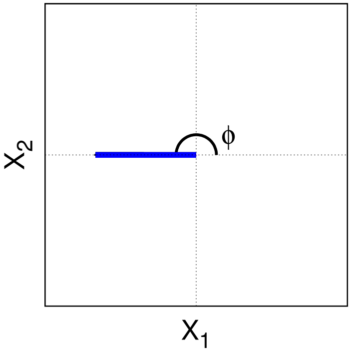

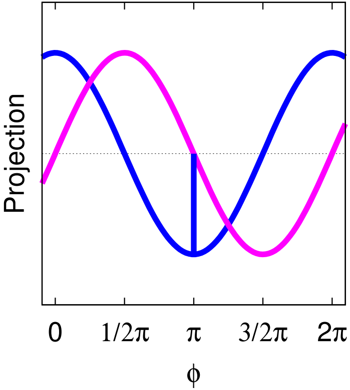

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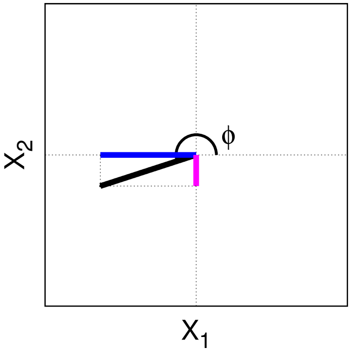

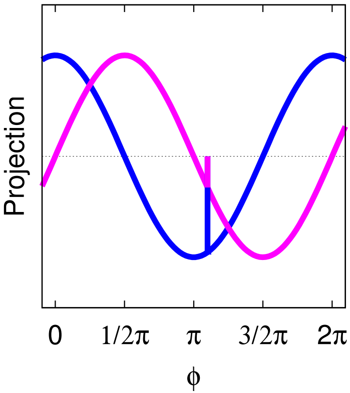

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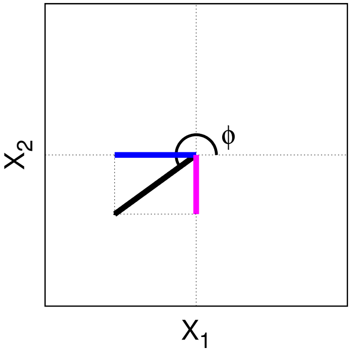

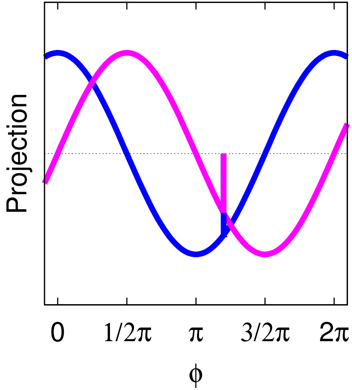

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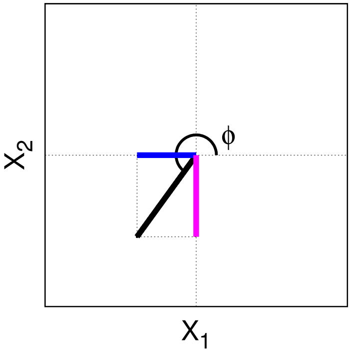

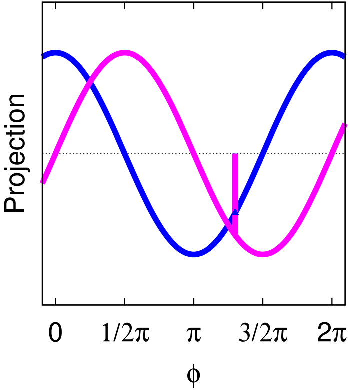

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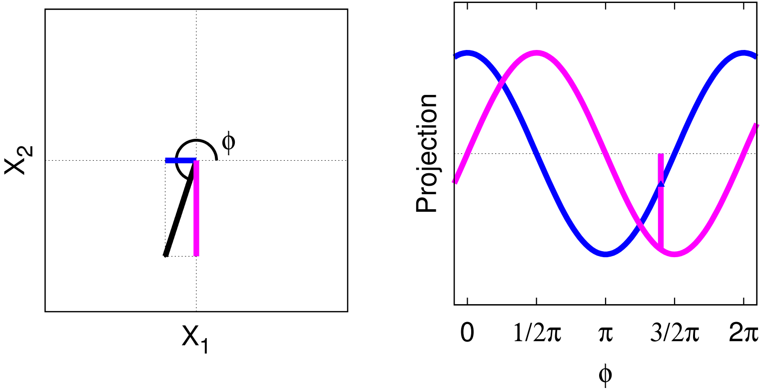

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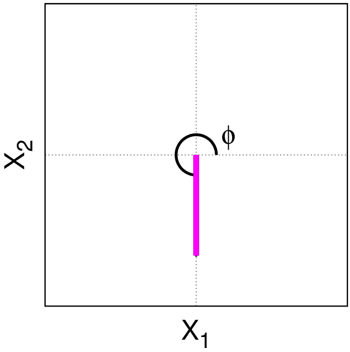

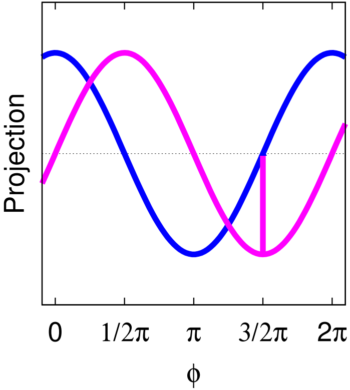

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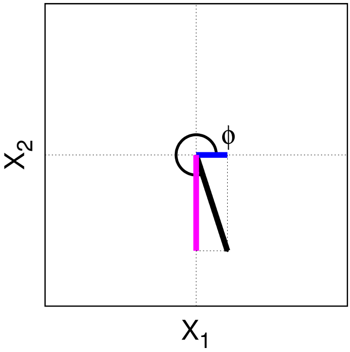

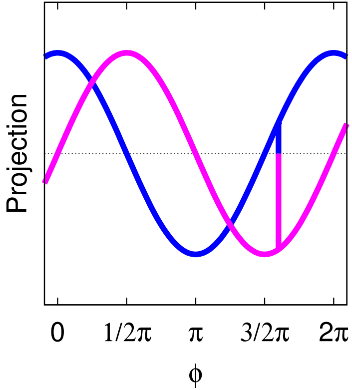

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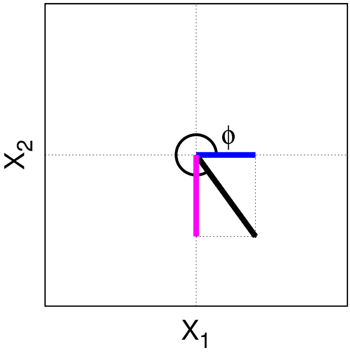

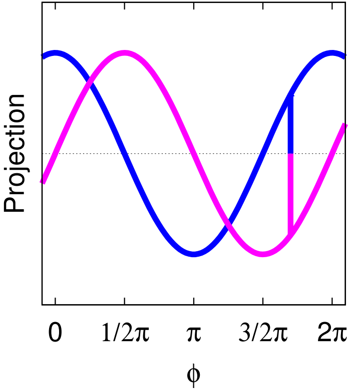

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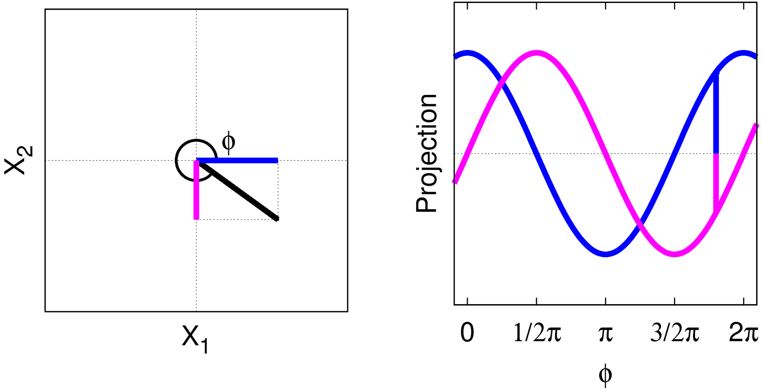

Classical quadratures vs time in a rotating frame

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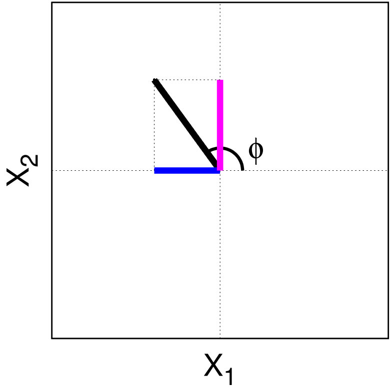

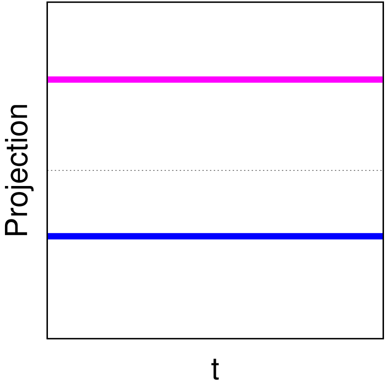

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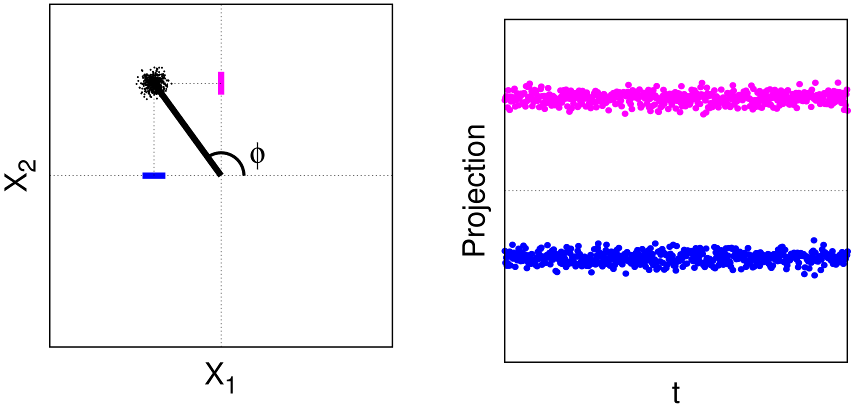
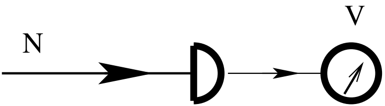


Detector quantum noise

Simple photodetector



$$V \sim N$$

$$\Delta V \sim \sqrt{N}$$

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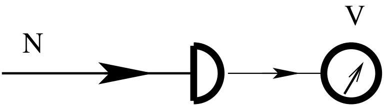

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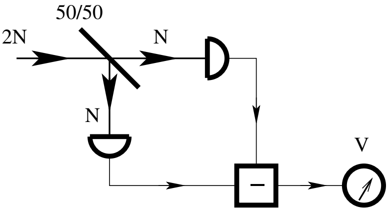
Balanced photodetector



$$V = 0$$

$$\Delta V \sim \sqrt{N}$$

Transition from classical to quantum field

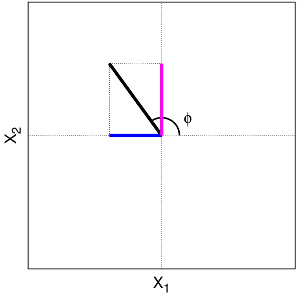
Classical analog

- Field amplitude a
- Field real part
  $X_1 = (a^* + a)/2$
- Field imaginary part
  $X_2 = i(a^* - a)/2$

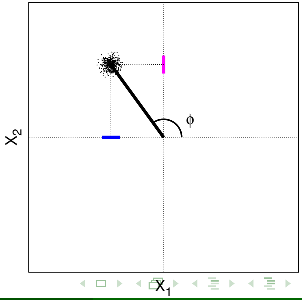
Quantum approach

- Field operator $\hat{a}$
- Amplitude quadrature
  $\hat{X}_1 = (\hat{a}^\dagger + \hat{a})/2$
- Phase quadrature
  $\hat{X}_2 = i(\hat{a}^\dagger - \hat{a})/2$

$$E(\phi) \quad = \quad |a|e^{-i\phi} = X_1 + iX_2$$



$$\hat{E}(\phi) \quad = \quad \hat{X}_1 + i\hat{X}_2$$



Heisenberg uncertainty principle and its optics equivalent

Heisenberg uncertainty principle

$\Delta p \Delta x \geq \hbar/2$

The more precisely the POSITION is determined, the less precisely the MOMENTUM is known, and vice versa

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Optics equivalent

$\Delta \phi \Delta N \geq 1$

The more precisely the PHASE is determined, the less precisely the AMPLITUDE is known, and vice versa

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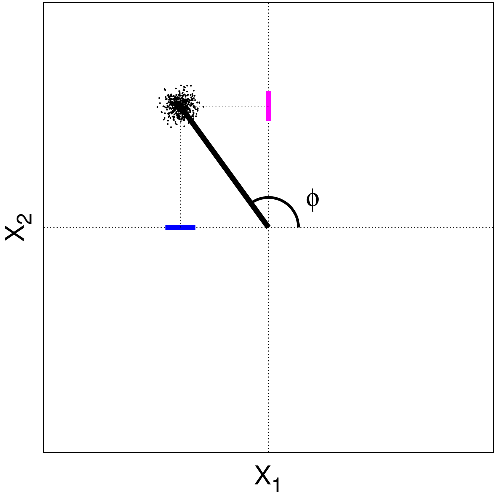
Optics equivalent strict definition

$\Delta X_1 \Delta X_2 \geq 1/4$

Quantum optics summary



Light consist of photons

- $\hat{N} = a^\dagger a$

Commutator relationship

- $[a, a^\dagger] = 1$
- $[X_1, X_2] = i/2$

Detectors measure

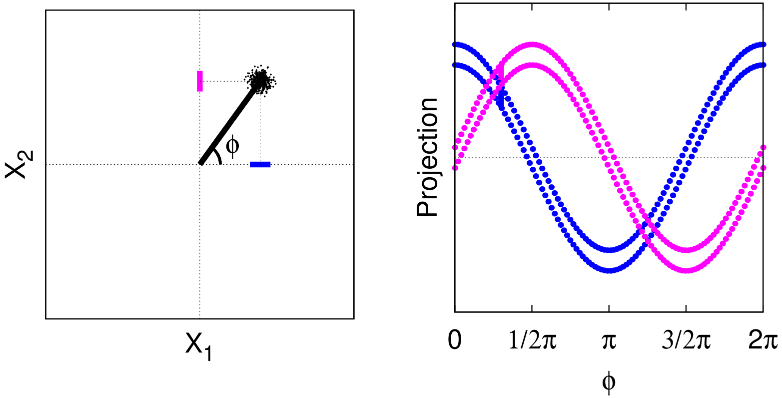
- number of photons $\hat{N}$
- Quadratures $\hat{X}_1$ and $\hat{X}_2$

Uncertainty relationship

- $\Delta X_1 \Delta X_2 \geq 1/4$

Coherent state is minimum uncertainty state

$\Delta X_1 \Delta X_2 = 1/4$



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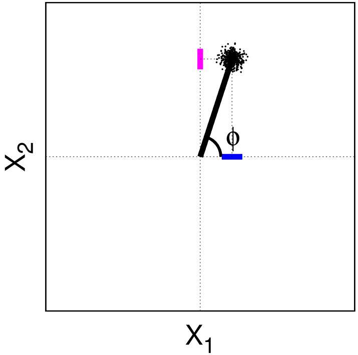

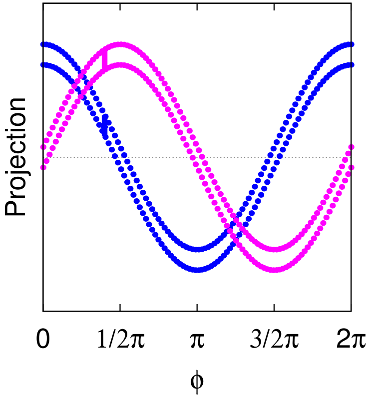

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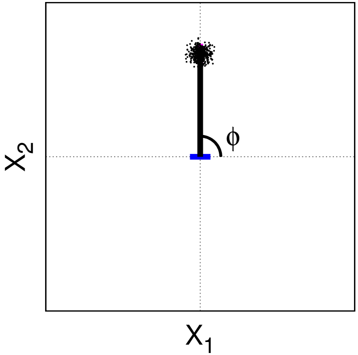

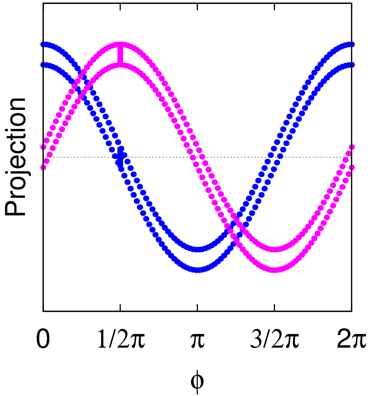

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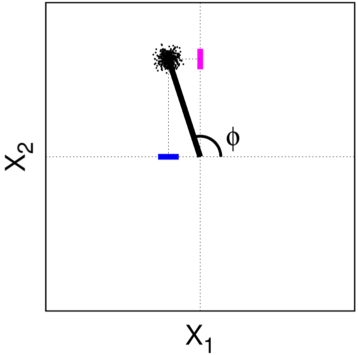

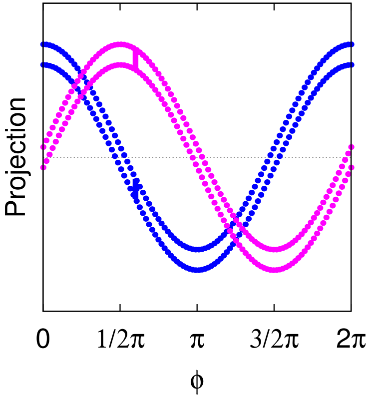

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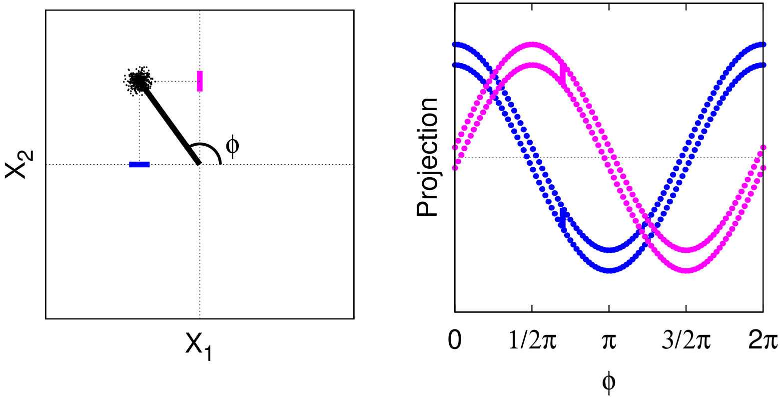

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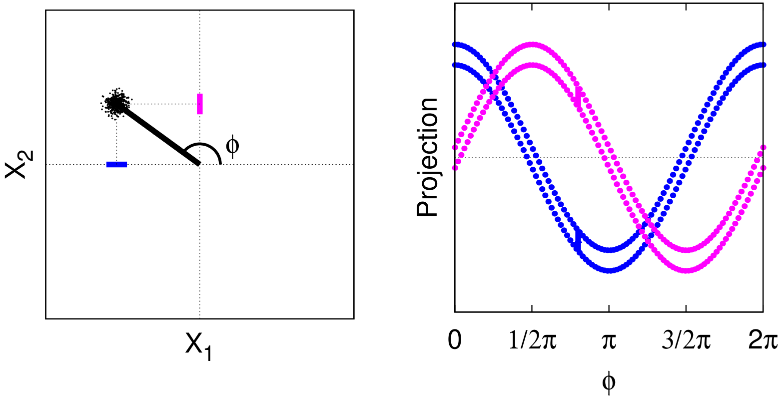

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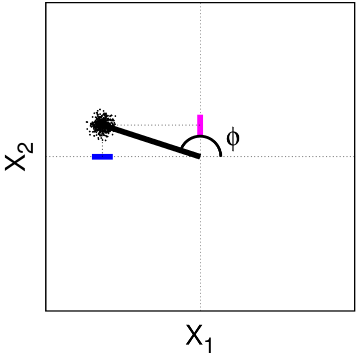

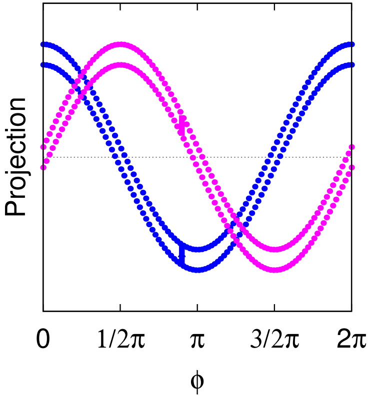

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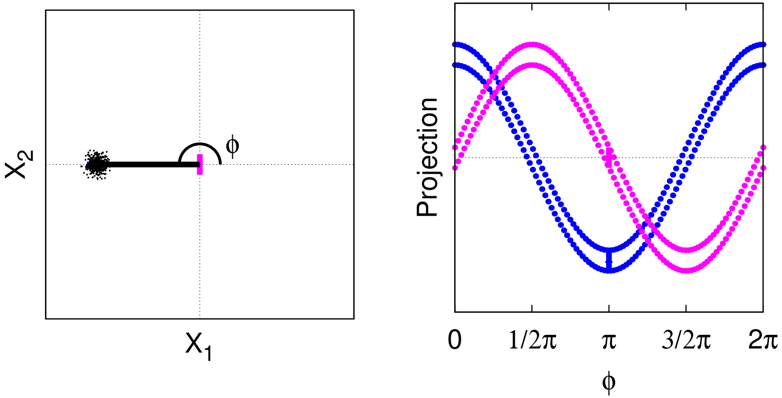

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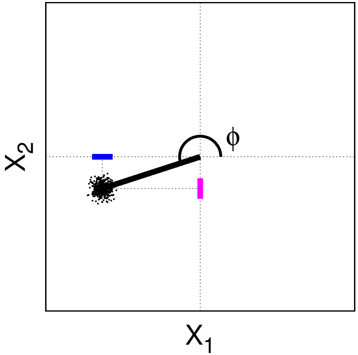

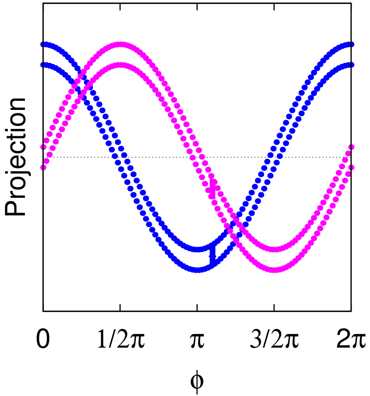

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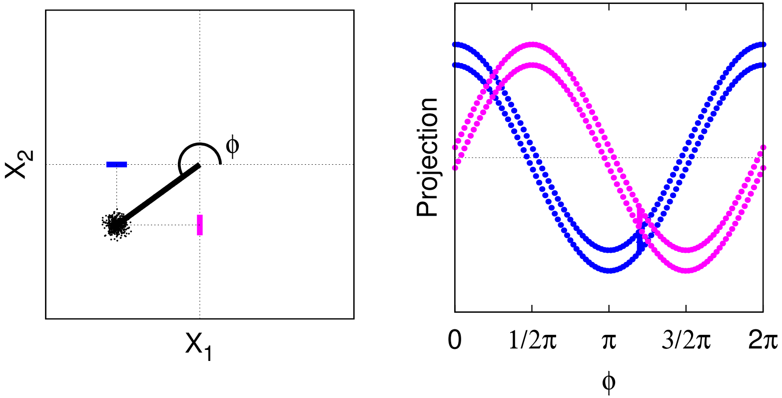

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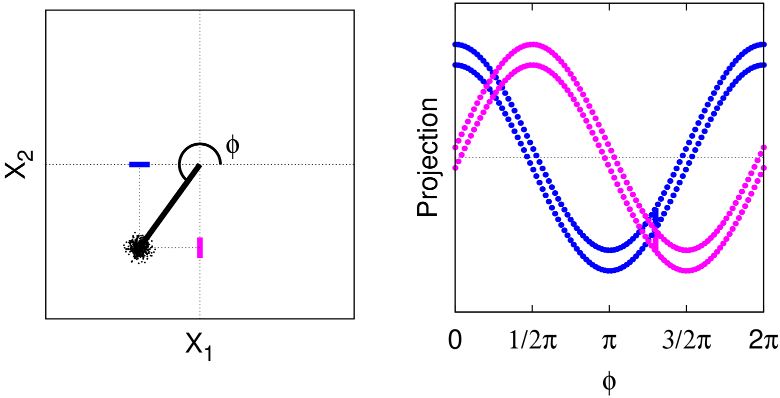

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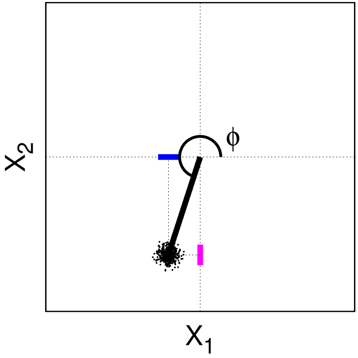
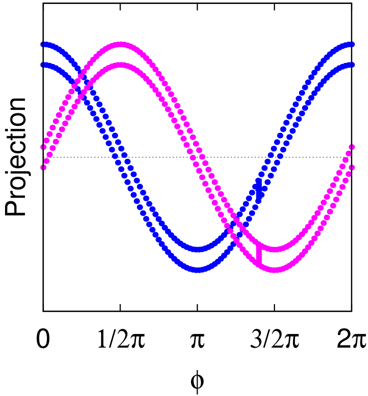
$\Delta X_1 \Delta X_2 = 1/4$





Coherent state is minimum uncertainty state

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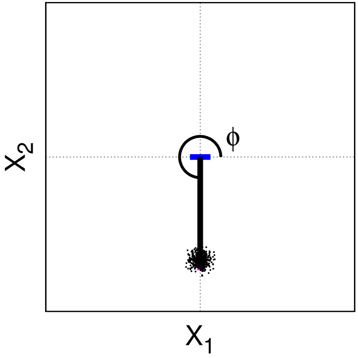

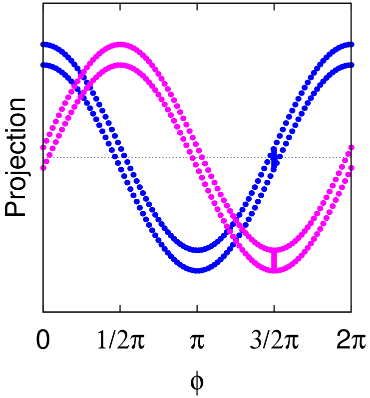

Coherent state is minimum uncertainty state

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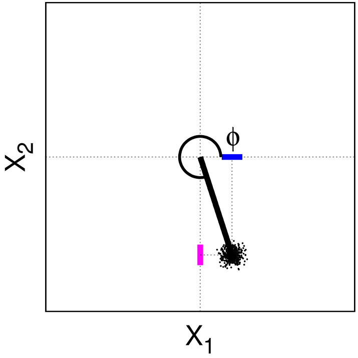

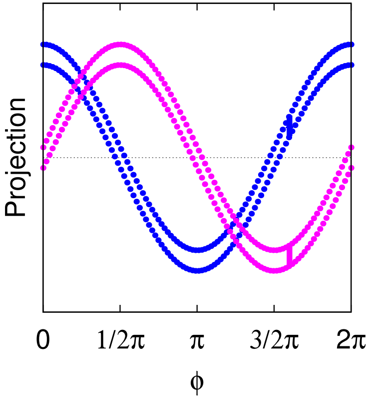

Coherent state is minimum uncertainty state

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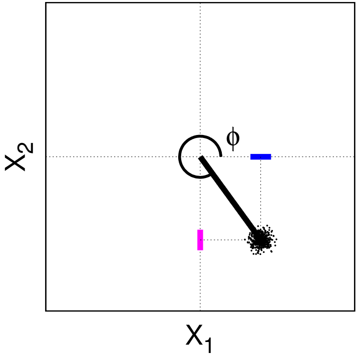

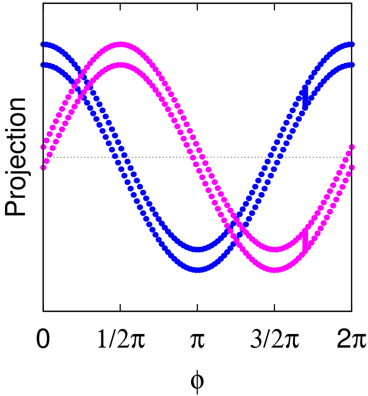

Coherent state is minimum uncertainty state

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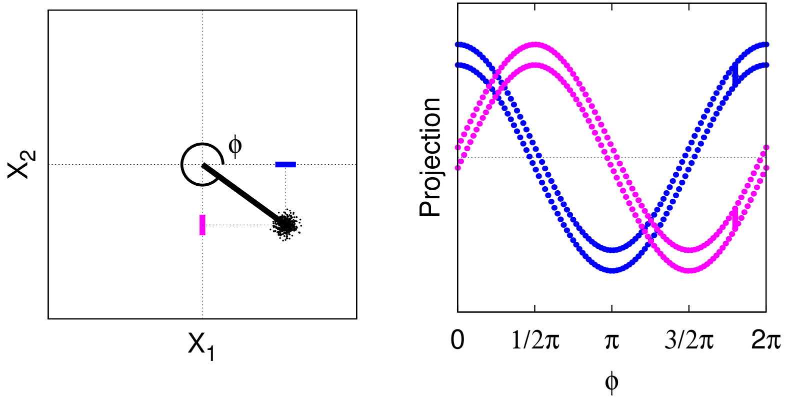
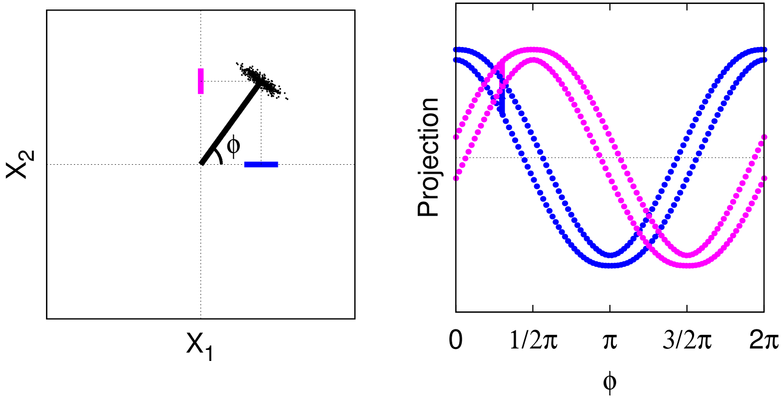


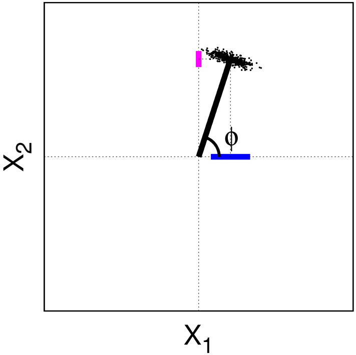
Amplitude squeezed states

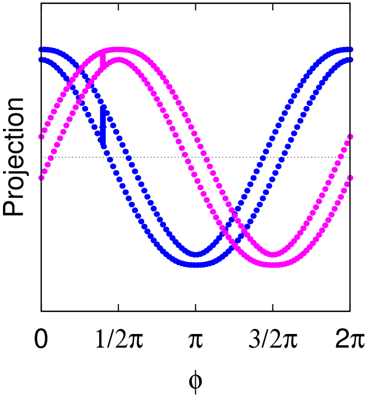
$\Delta X_1 \Delta X_2 = 1/4$



Amplitude squeezed states

$\Delta X_1 \Delta X_2 = 1/4$





Amplitude squeezed states

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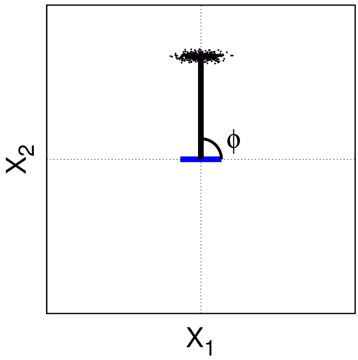

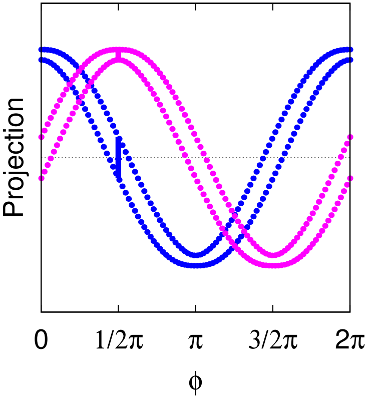

Amplitude squeezed states

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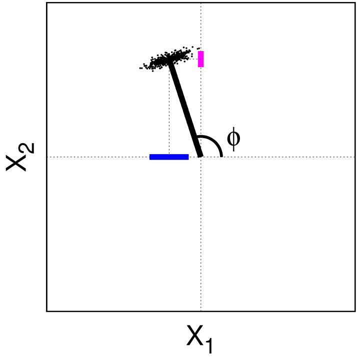

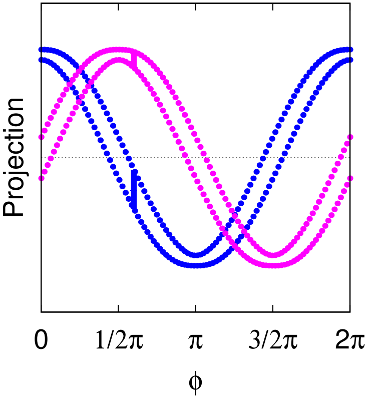

Amplitude squeezed states

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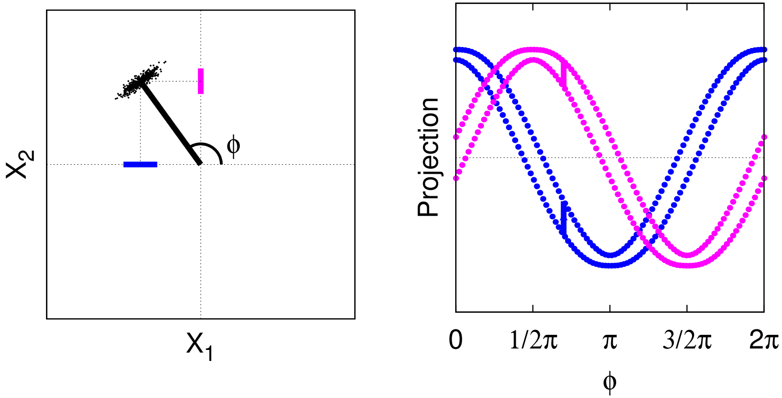

Amplitude squeezed states

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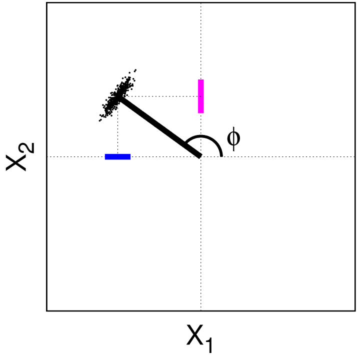

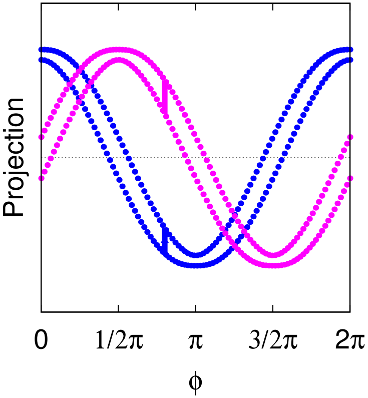

Amplitude squeezed states

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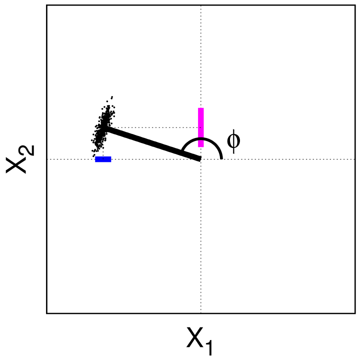

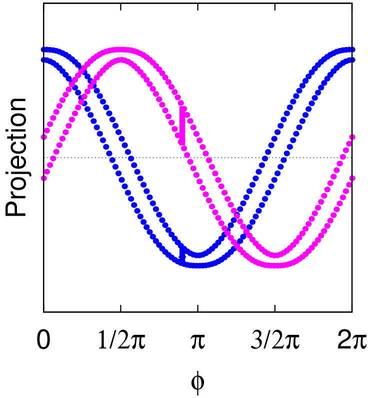

Amplitude squeezed states

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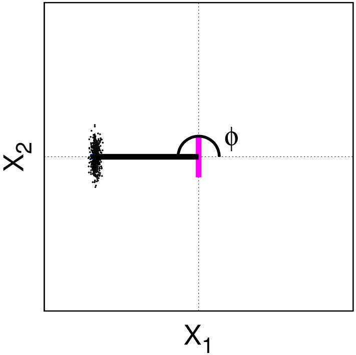

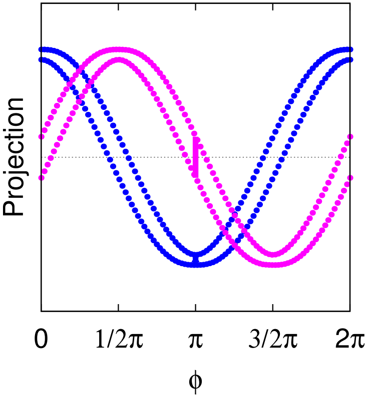

Amplitude squeezed states

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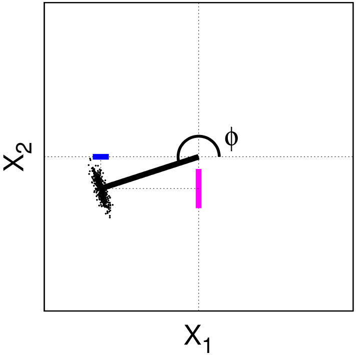

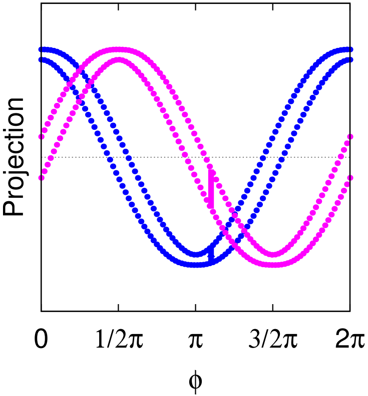

Amplitude squeezed states

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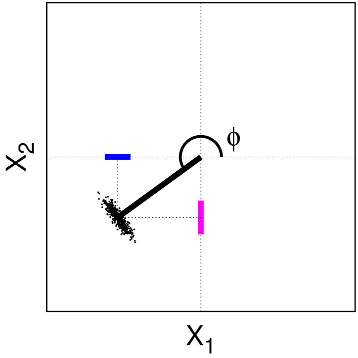

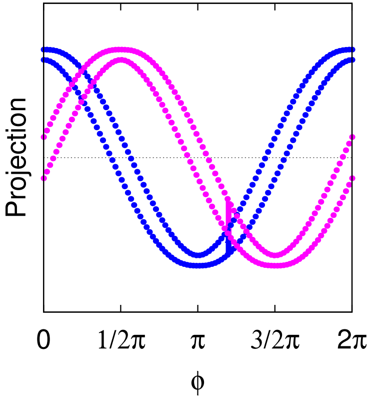

Amplitude squeezed states

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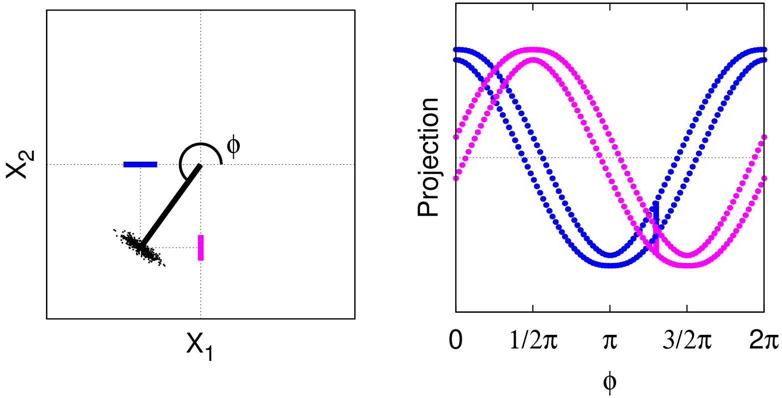

Amplitude squeezed states

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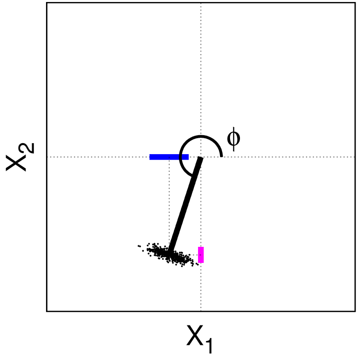

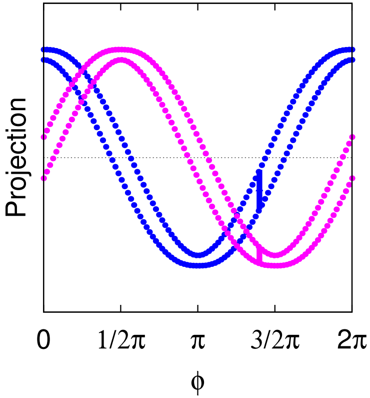

Amplitude squeezed states

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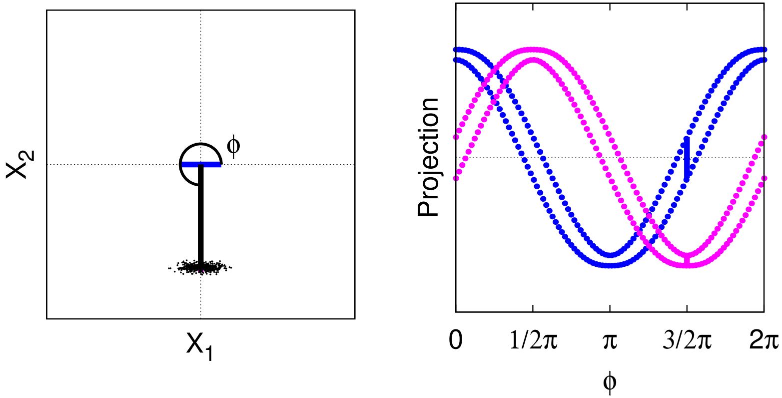

Amplitude squeezed states

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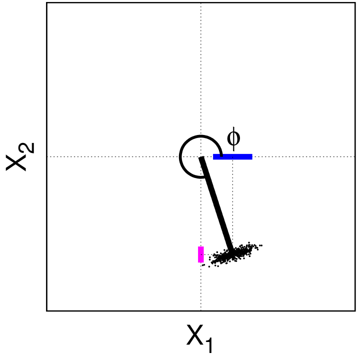

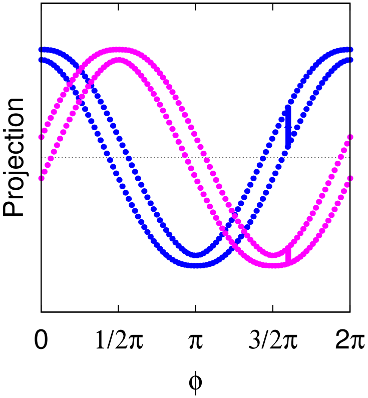

Amplitude squeezed states

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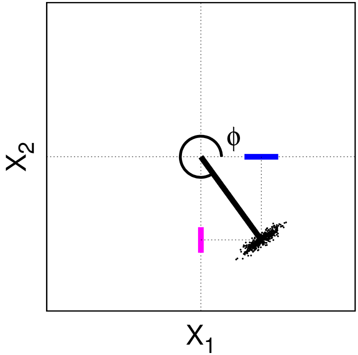

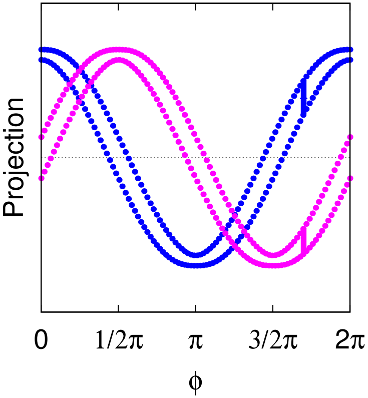

Amplitude squeezed states

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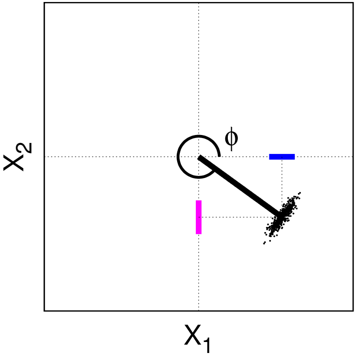

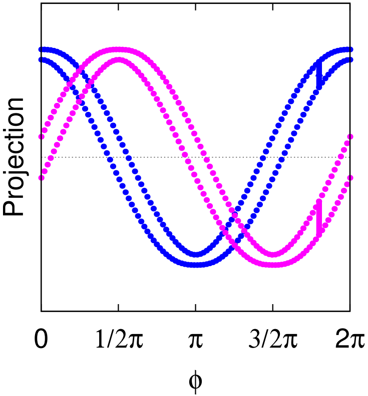

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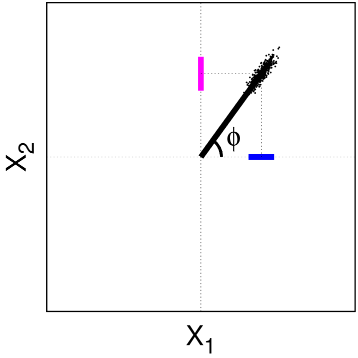
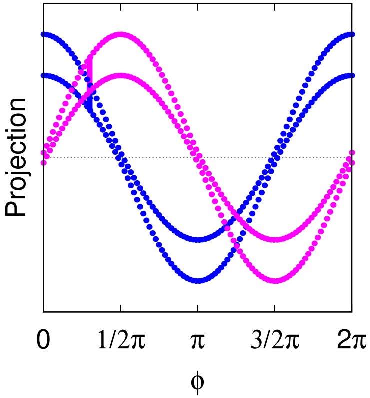
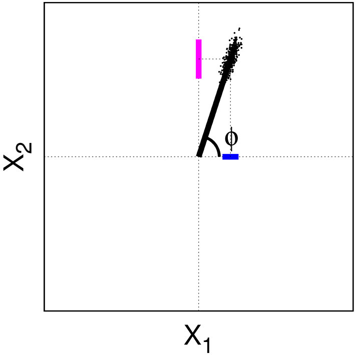
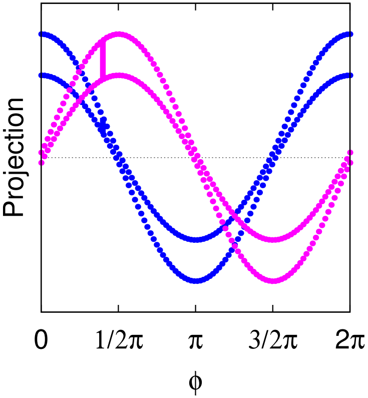




Phase squeezed states

$\Delta X_1 \Delta X_2 = 1/4$





Phase squeezed states

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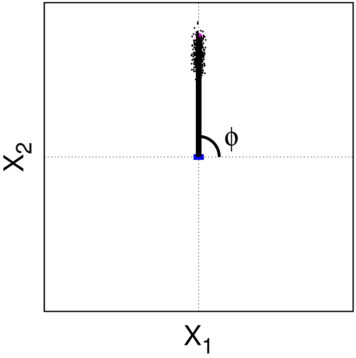

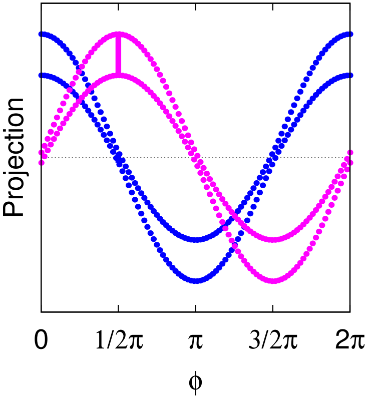

Phase squeezed states

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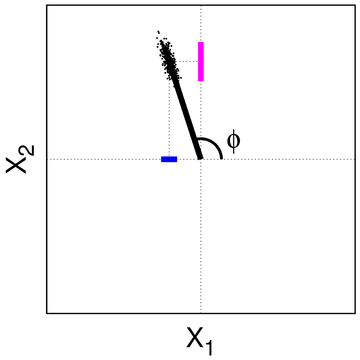

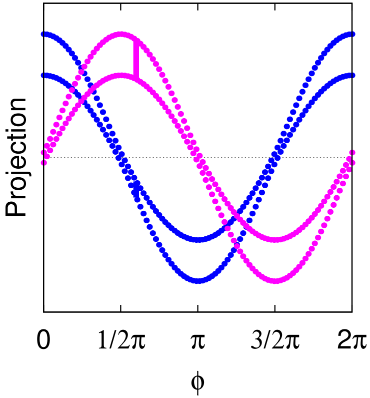

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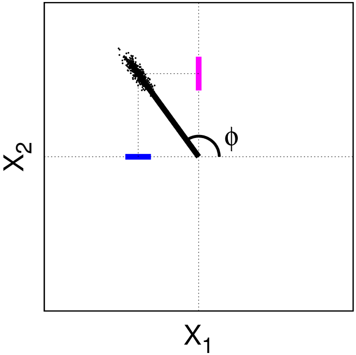

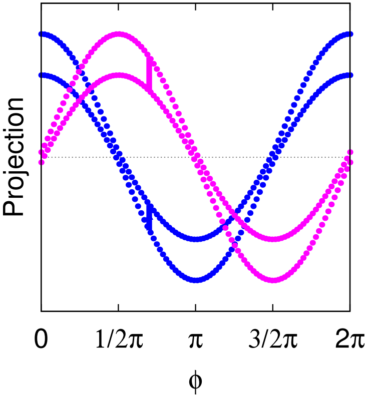

Phase squeezed states

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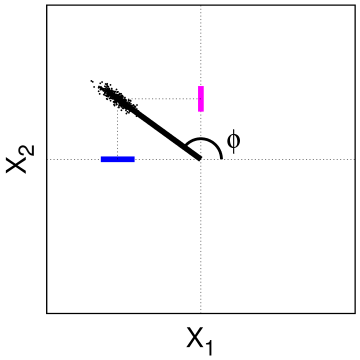

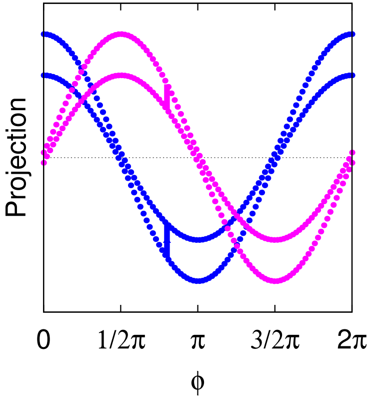

Phase squeezed states

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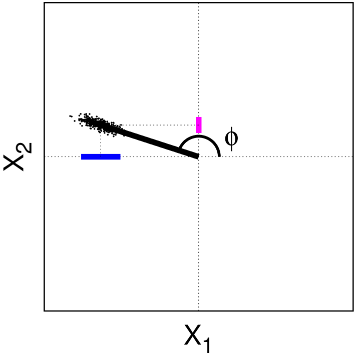

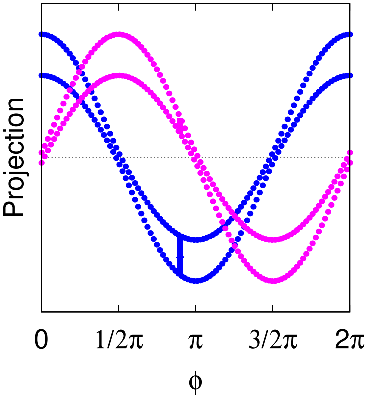

Phase squeezed states

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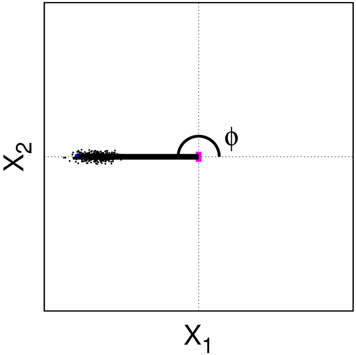

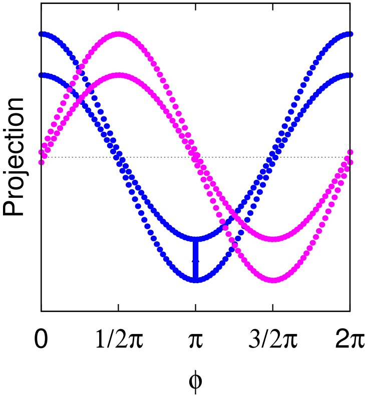

Phase squeezed states

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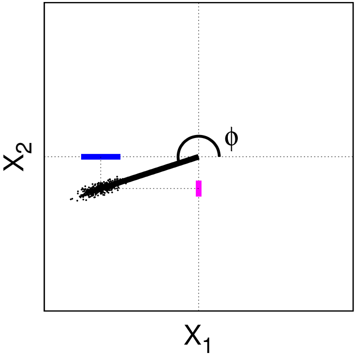

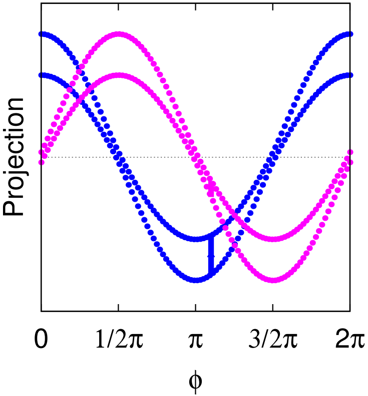

Phase squeezed states

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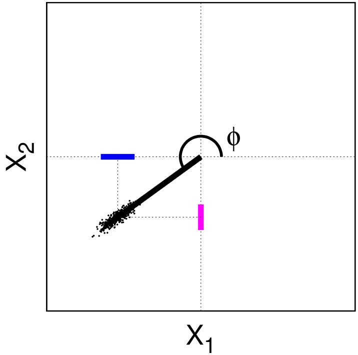

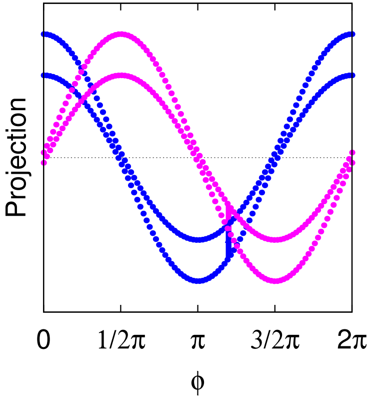

Phase squeezed states

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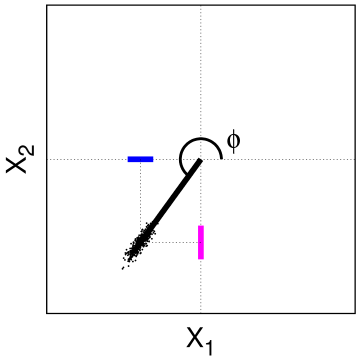

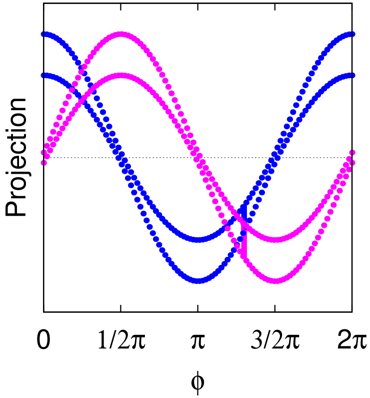

Phase squeezed states

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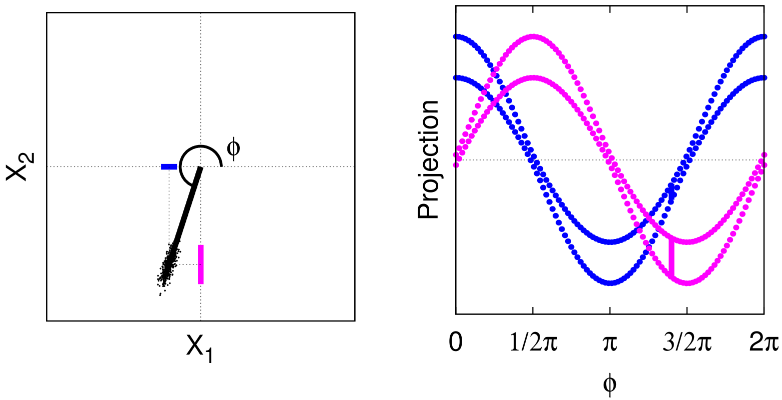

Phase squeezed states

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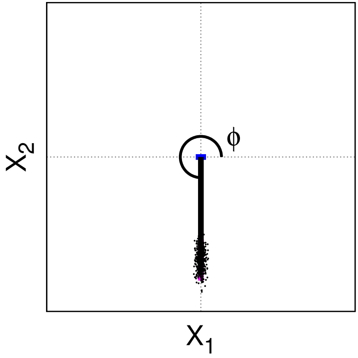

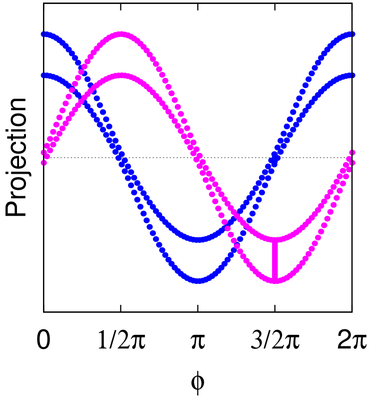

Phase squeezed states

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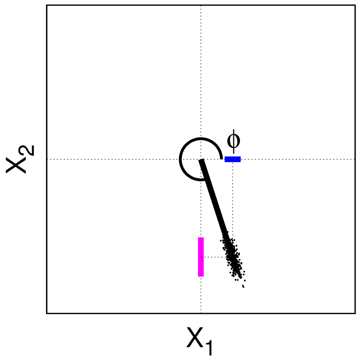

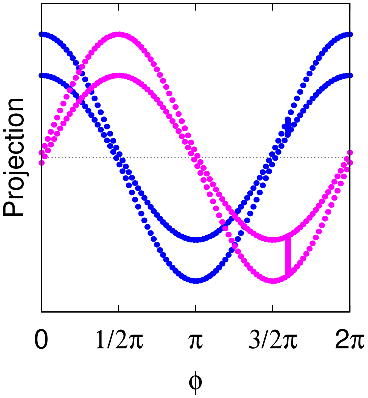

Phase squeezed states

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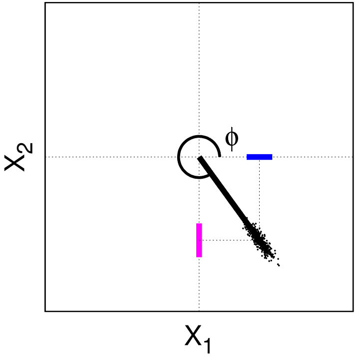

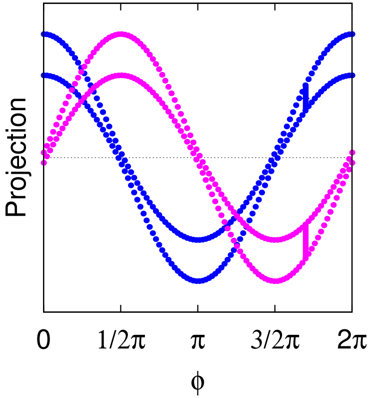

Phase squeezed states

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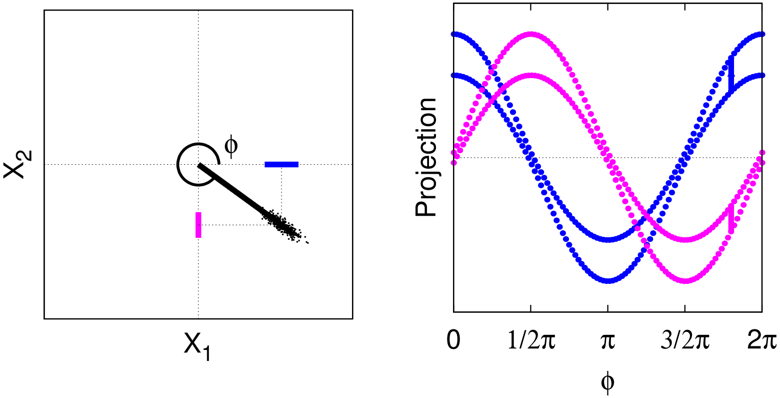
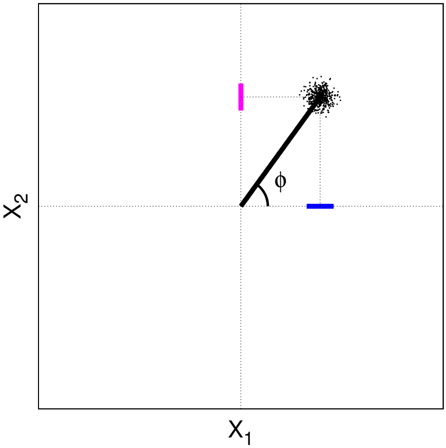
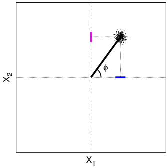


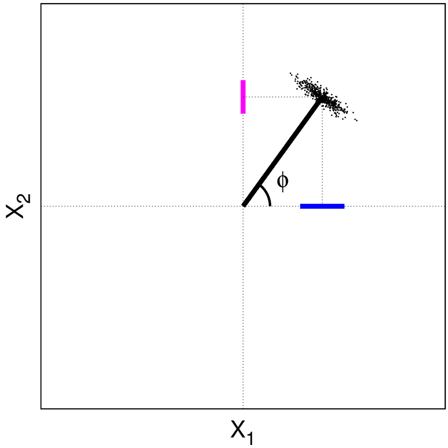
Squeezed quantum states zoo



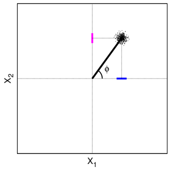
Unsqueezed coherent



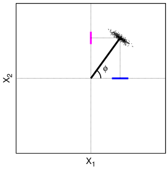
Squeezed quantum states zoo



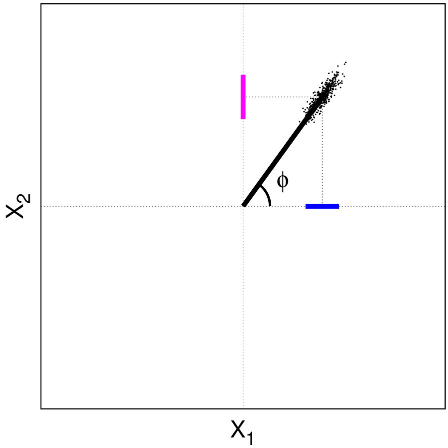
Unsqueezed coherent



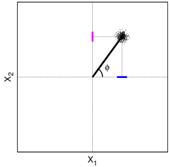
Amplitude squeezed



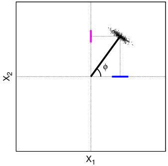
Squeezed quantum states zoo



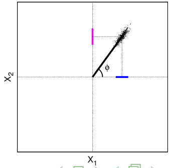
Unsqueezed coherent



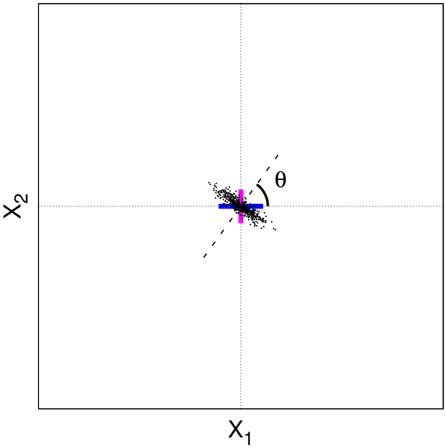
Amplitude squeezed



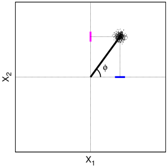
Phase squeezed



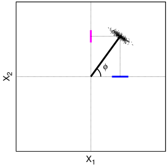
Squeezed quantum states zoo



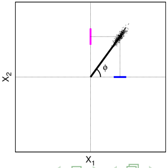
Unsqueezed coherent



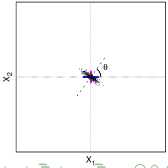
Amplitude squeezed



Phase squeezed



Vacuum squeezed

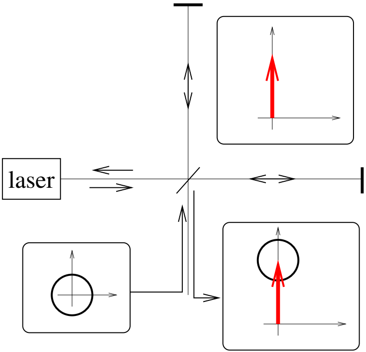


Possible squeezing applications

- improvements any shot noise limited optical sensors
- noiseless signal amplification
- secure communications (you would notice eavesdropper)
- photon pair generation, entanglement, true single photon sources
- interferometers sensitivity boost (for example gravitational wave antennas)
- light free measurements
- quantum memory probe and information carrier

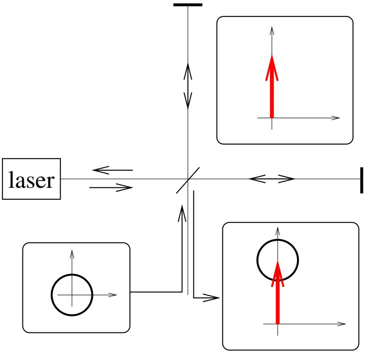
Squeezing and interferometer

Vacuum input

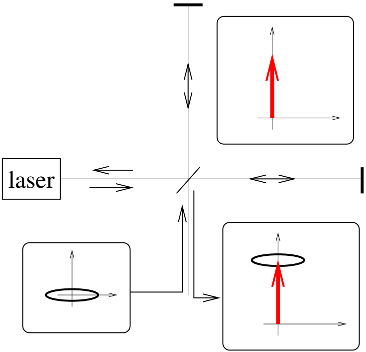


Squeezing and interferometer

Vacuum input

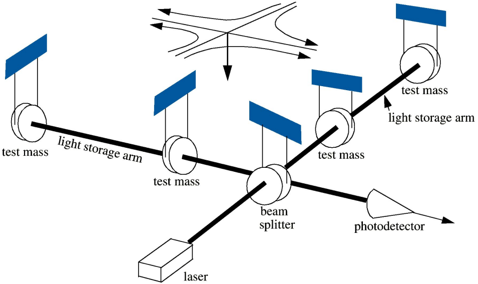


Squeezed input

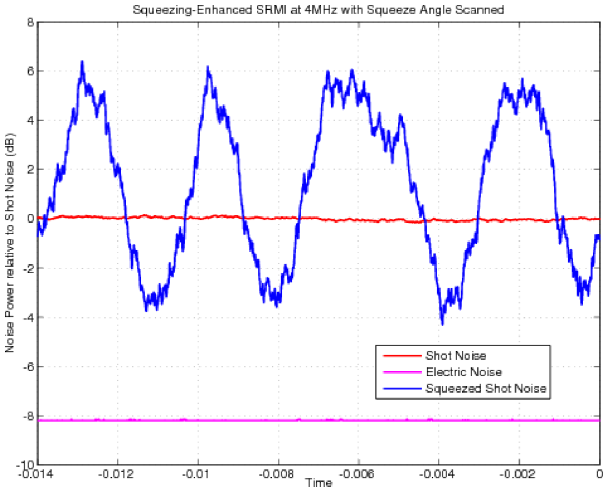


Laser Interferometer Gravitational-wave Observatory

- $L = 4$ km
- $h \sim 10^{-21}$
- $\Delta L \sim 10^{-18}$ m
- $\Delta\phi \sim 10^{-10}$ rad

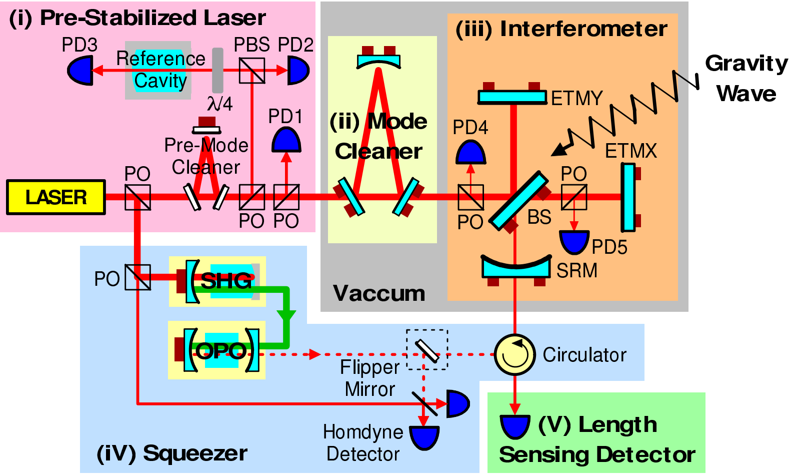


Squeezing level vs time (unlocked)

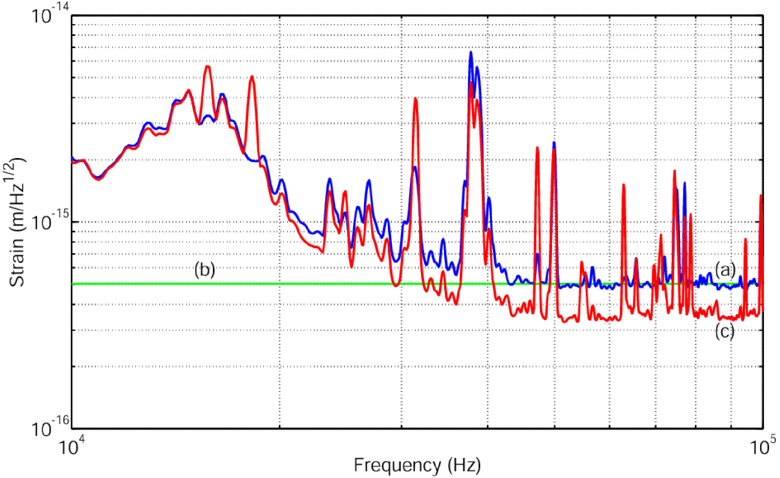


“A quantum-enhanced prototype gravitational-wave detector”, Nature Physics, **4**, 472-476, (2008).

GW 40m detector and squeezer



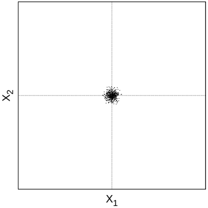
GW 40m detector with 4dB of squeezed vacuum



Signal to noise improvement by factor of 1.43

Squeezed field generation recipe

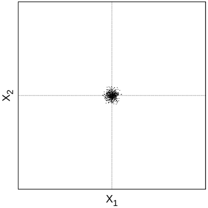
Take a vacuum state $|0>$



$$H = \frac{1}{2}$$

Squeezed field generation recipe

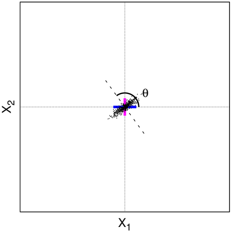
Take a vacuum state $|0>$



$$H = \frac{1}{2}$$

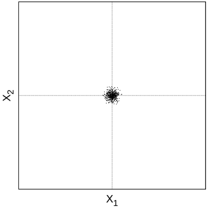
Apply squeezing operator $|\xi>=\hat{S}(\xi)|0>$

$$\hat{S}(\xi) = e^{\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi a^{\dagger 2}}$$



Squeezed field generation recipe

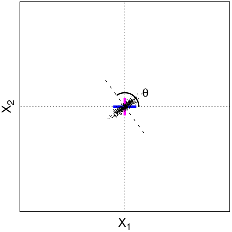
Take a vacuum state $|0>$



$$H = \frac{1}{2}$$

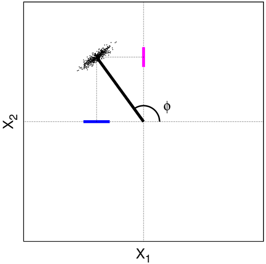
Apply squeezing operator $|\xi> = \hat{S}(\xi)|0>$

$$\hat{S}(\xi) = e^{\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi a^{\dagger 2}}$$



Apply displacement operator $|\alpha, \xi> = \hat{D}(\alpha)|s>$

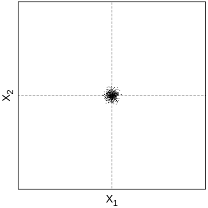
$$\hat{D}(\alpha) = e^{\alpha a^\dagger - \alpha^* a}$$



$$< \alpha, \xi | X_1 | \alpha, \xi > = Re(\alpha),$$
$$< \alpha, \xi | X_2 | \alpha, \xi > = Im(\alpha)$$

Squeezed field generation recipe

Take a vacuum state $|0>$

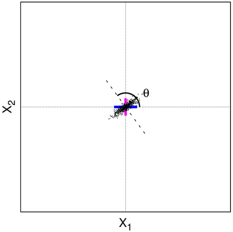


$$H = \frac{1}{2}$$

Notice $\Delta X_1 \Delta X_2 = \frac{1}{4}$

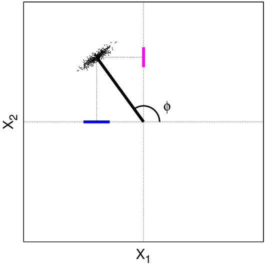
Apply squeezing operator $|\xi >= \hat{S}(\xi)|0>$

$$\hat{S}(\xi) = e^{\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi a^{\dagger 2}}$$



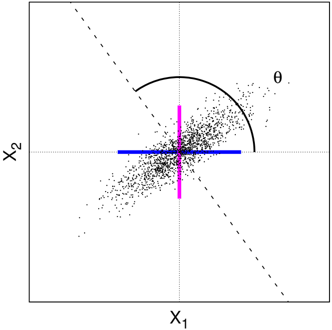
Apply displacement operator $|\alpha, \xi >= \hat{D}(\alpha)|s>$

$$\hat{D}(\alpha) = e^{\alpha a^\dagger - \alpha^* a}$$



$$< \alpha, \xi | X_1 | \alpha, \xi > = Re(\alpha),$$
$$< \alpha, \xi | X_2 | \alpha, \xi > = Im(\alpha)$$

Squeezed state $|\xi>=\hat{S}(\xi)|0>$ properties



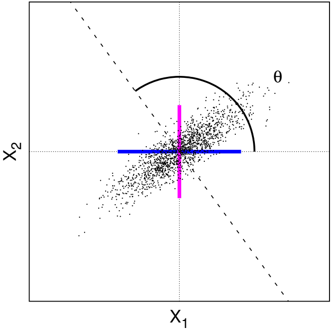
$$\hat{S}(\xi) = e^{\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi a^{\dagger 2}}, \xi = re^{i\theta}$$

If $\theta = 0$

$$< \xi|(\Delta X_1)^2|\xi > \quad = \quad \frac{1}{4}e^{-2r}$$
$$< \xi|(\Delta X_2)^2|\xi > \quad = \quad \frac{1}{4}e^{2r}$$

$$< \xi|(\Delta X_1)^2|\xi > \quad = \quad \frac{1}{4}(\cosh^2 r + \sinh^2 r - 2\sinh r \cosh r \cos\theta)$$
$$< \xi|(\Delta X_2)^2|\xi > \quad = \quad \frac{1}{4}(\cosh^2 r + \sinh^2 r + 2\sinh r \cosh r \cos\theta)$$

Photon number of squeezed state $|\xi>$



Probability to detect given number of photons $C =< n|\xi >$ for squeezed vacuum

- even

$$C_{2m} = (-1)\frac{\sqrt{(2m)!}}{2^m m!}\frac{(e^{i\theta}\tanh r)^m}{\sqrt{\cosh r}}$$

- odd

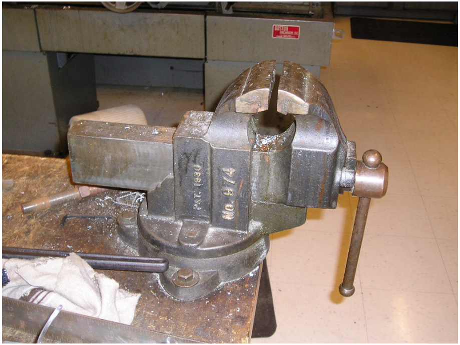
$$C_{2m+1} = 0$$

Average number of photons in general squeezed state

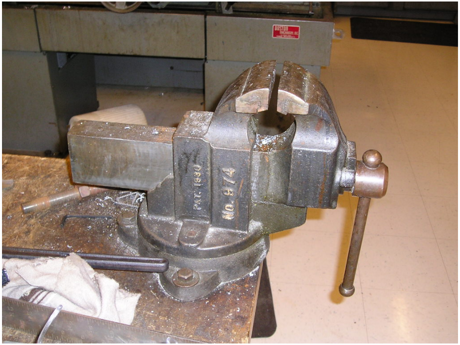
$$< \alpha, \xi|a^\dagger a|\alpha, \xi >= \alpha + \sinh^2 r$$

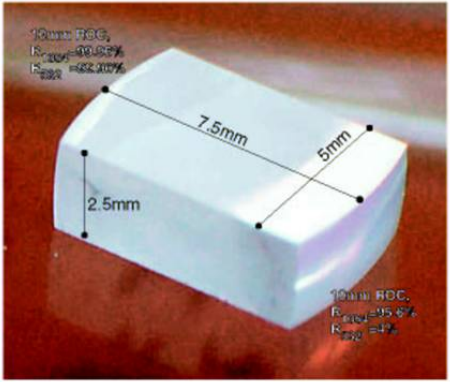
Tools for squeezing

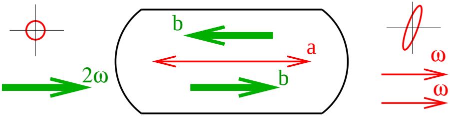
Tools for squeezing



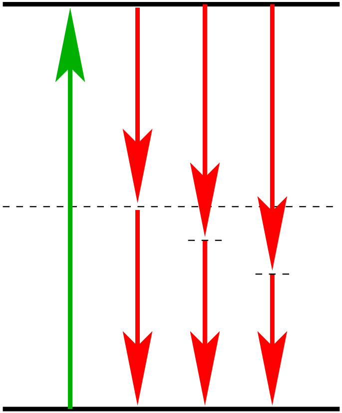
Tools for squeezing







Two photon squeezing picture

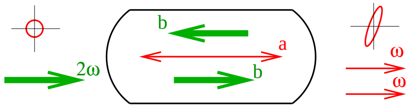


Squeezing operator

$$\hat{S}(\xi) = e^{\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi a^{\dagger 2}}$$

Parametric down-conversion in crystal

$$\hat{H} = i\hbar\chi^{(2)}(a^2 b^\dagger - a^{\dagger 2} b)$$

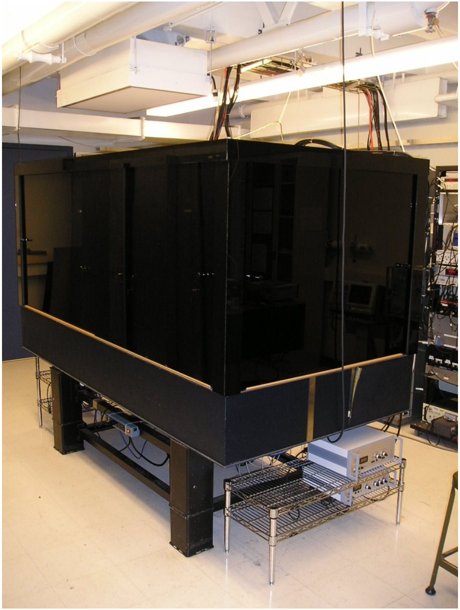


Squeezing

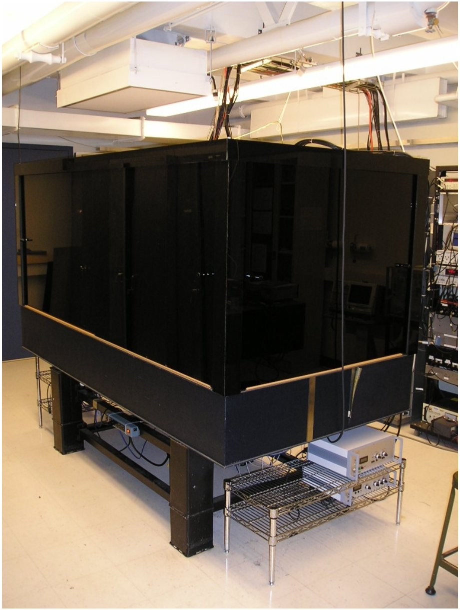
result of correlation of upper and lower sidebands

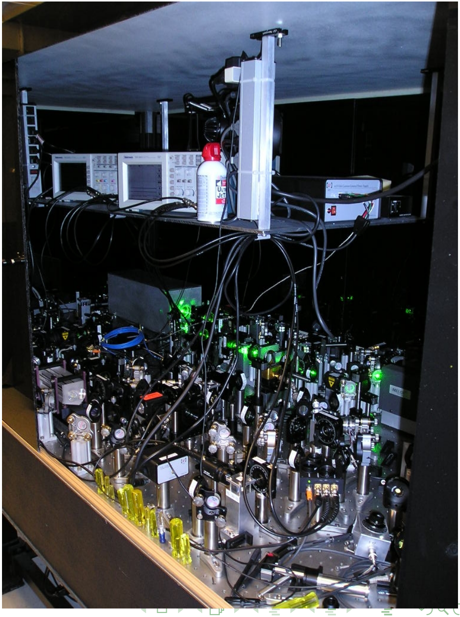
Squeezer appearance

Squeezer appearance

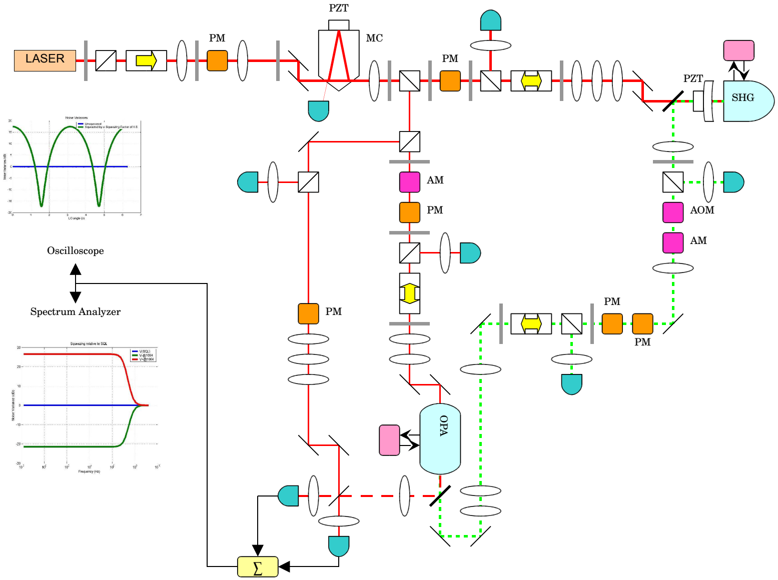


Squeezer appearance

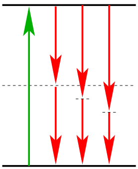




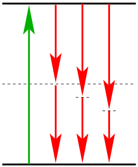
Crystal squeezing setup scheme

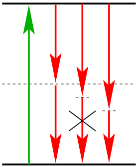


Cavity parameters with squeezing

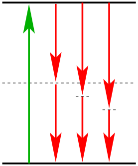


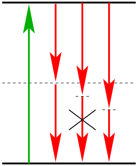
Cavity parameters with squeezing

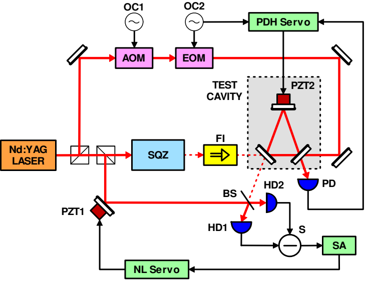




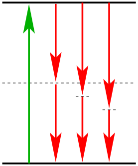
Cavity parameters with squeezing

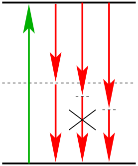


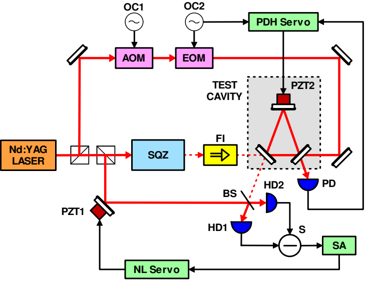


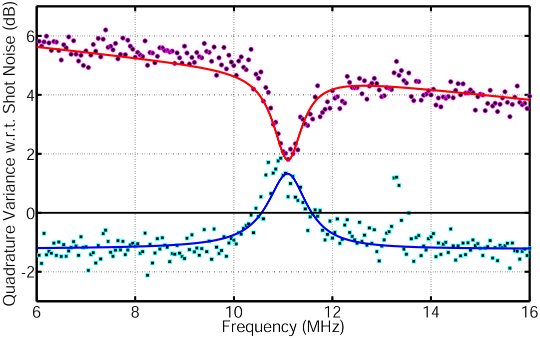


Cavity parameters with squeezing









“Noninvasive measurements of cavity parameters by use of squeezed vacuum”, Physical Review A, **74**, 033817, (2006).

Summary for crystal squeezing

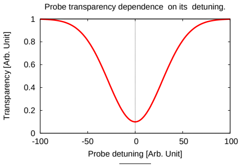
Pros

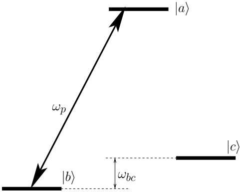
- mainstream: many different nonlinear crystals available
- so far the best squeezers
  - maximum squeezing value detected 11.5 dB at 1064 nm
  - Moritz Mehmet, Henning Vahlbruch, Nico Lastzka, Karsten Danzmann, and Roman Schnabel, "Observation of squeezed states with strong photon-number oscillations", Phys. Rev. A **81**, 013814 (2010)
- well understood

Cons

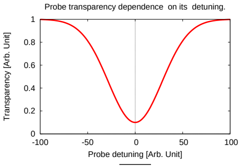
- crystals have limited transparency window
- thus squeezing is hard to generate at visible wavelength
  - at 795 nm only 4-6 dB squeezing is reported
- this limits applications of such squeezers for spectroscopy

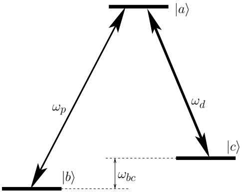
Quantum memory with atomic ensembles

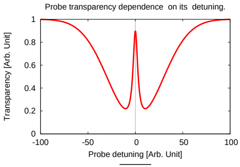




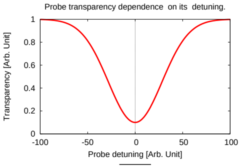
Quantum memory with atomic ensembles

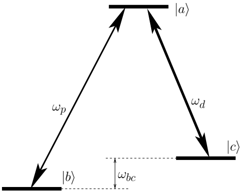


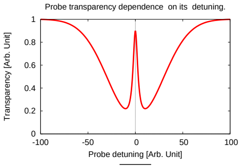




Quantum memory with atomic ensembles

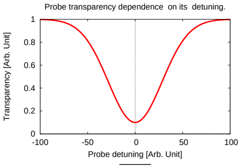


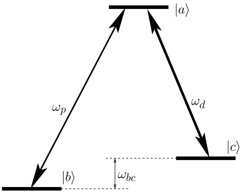


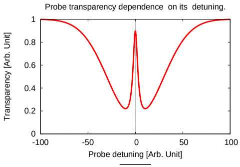


Storage and retrieval

Quantum memory with atomic ensembles



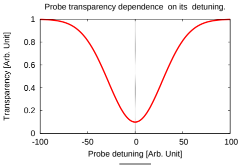


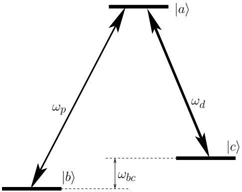


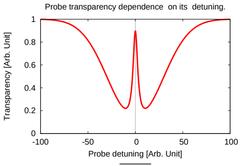
Storage and retrieval

- single photon

Quantum memory with atomic ensembles







Storage and retrieval

- single photon
- squeezed state (Furusawa and Lvovsky PRL **100** 2008)

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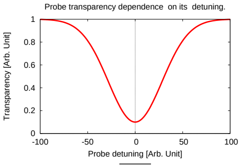

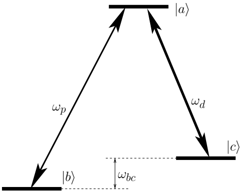

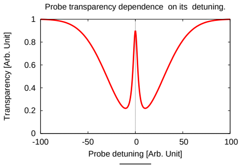

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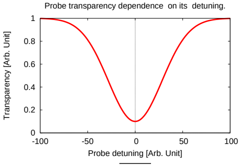
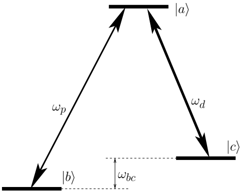
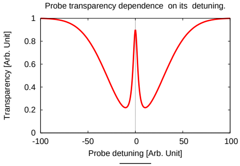
Squeezed state requirements for a quantum memory probe

Quantum memory with atomic ensembles







Storage and retrieval

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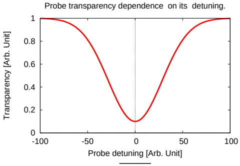
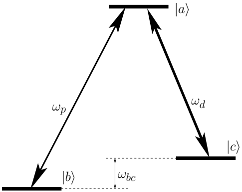
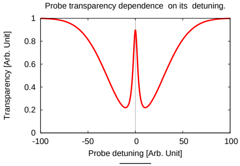
Squeezed state requirements for a quantum memory probe

- squeezing carrier at atomic wavelength (780nm, 795nm)
- squeezing within narrow resonance window at frequencies($<$100kHz)

Quantum memory with atomic ensembles







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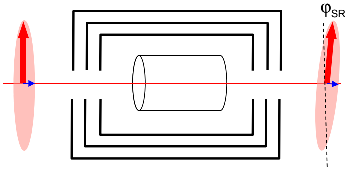
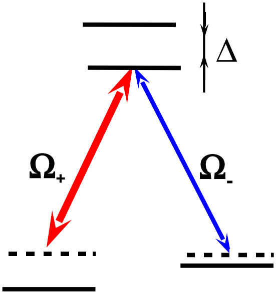
Traditional nonlinear crystal based squeezers are capable of it, but they are extremely technically challenging especially at short wave length.

Self-rotation of elliptical polarization in atomic medium

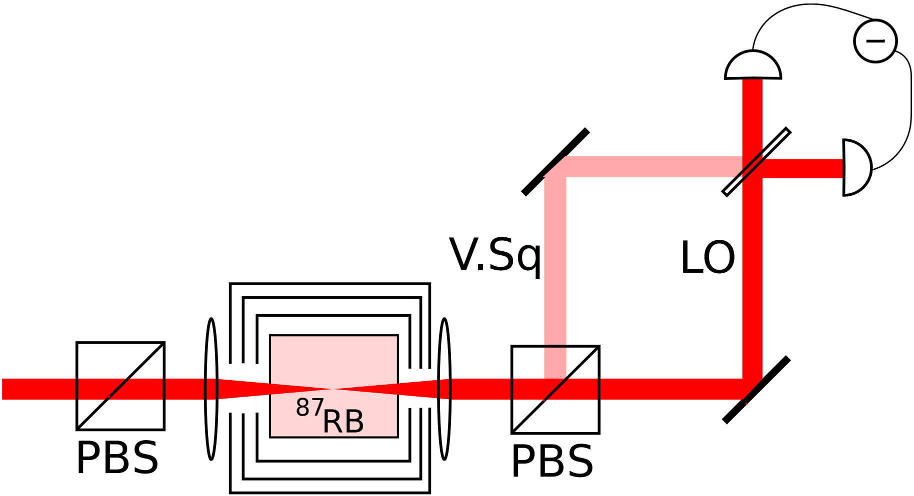




A.B. Matsko et al., PRA 66, 043815 (2002): theoretically prediction of 4-6 dB noise suppression

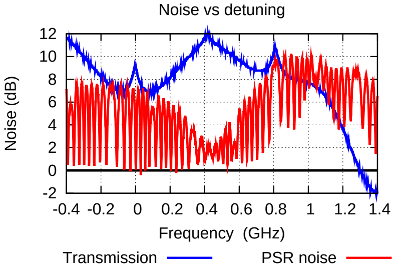
$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^{\dagger} - a_{in}) \quad (1)$$

Setup

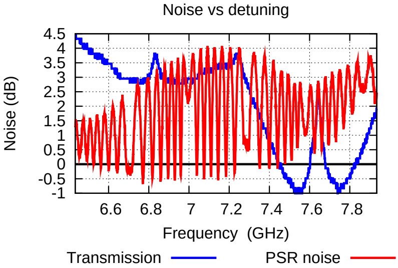


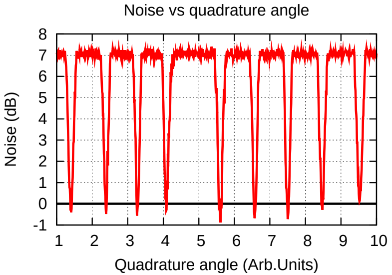
Noise contrast vs detuning in hot ^{87}Rb vacuum cell

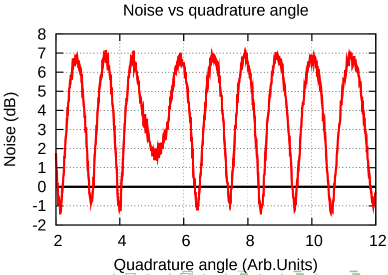
$F_g = 2 \to F_e = 1, 2$



$F_g = 1 \to F_e = 1, 2$

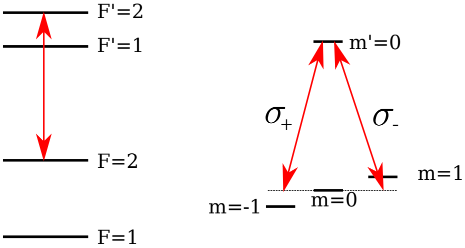


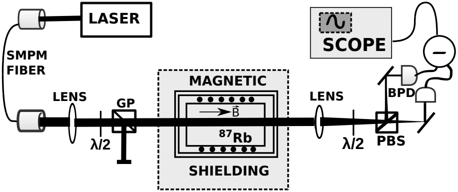




Optical magnetometer and non linear Faraday effect

^{87}Rb D_1 line





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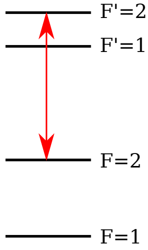
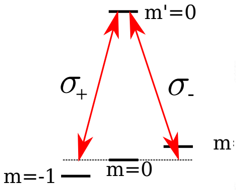
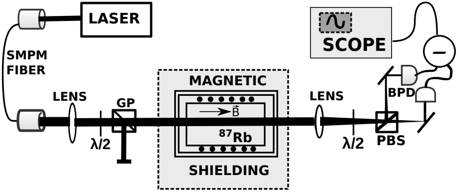
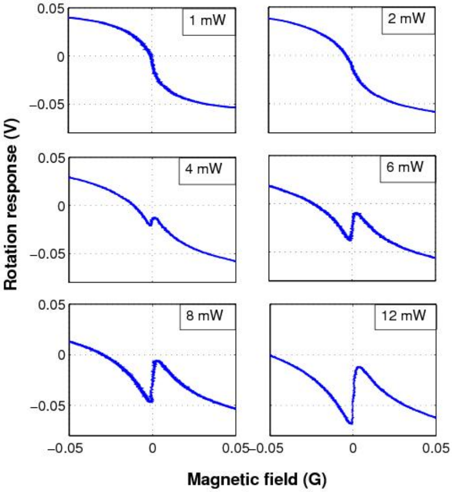
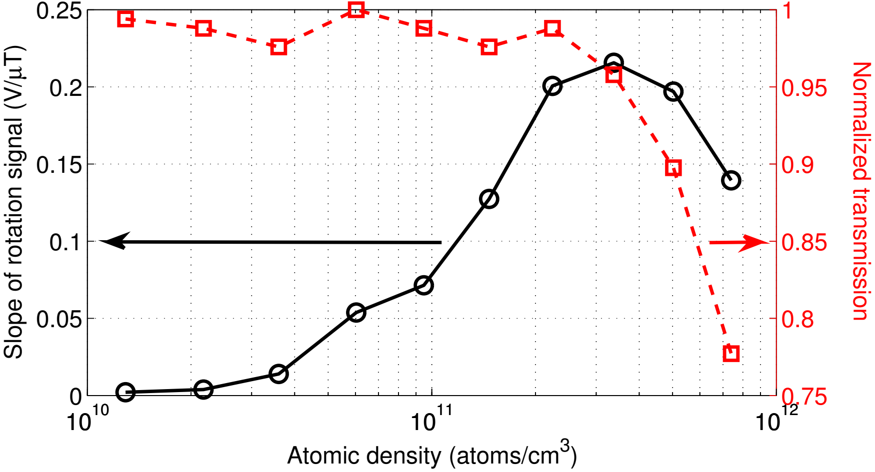




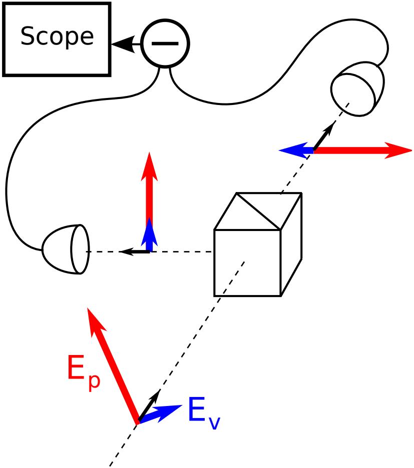




Magnetometer response vs atomic density

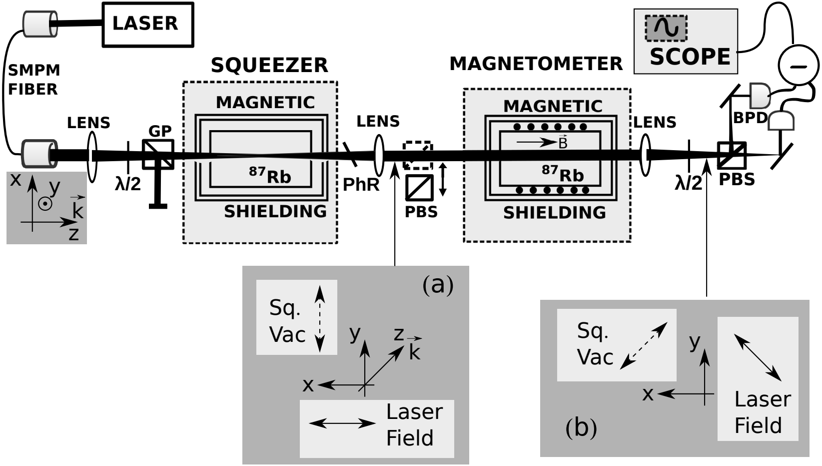


Shot noise limit of the magnetometer



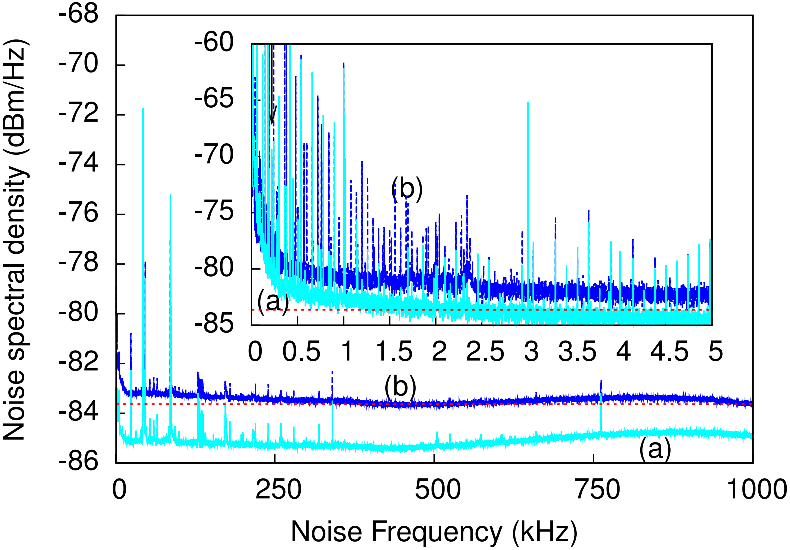
$$S = |E_p + E_v|^2 - |E_p - E_v|^2$$
$$S = 4E_pE_v$$
$$< \Delta S > \sim E_p < \Delta E_v >$$

Squeezed enhanced magnetometer setup

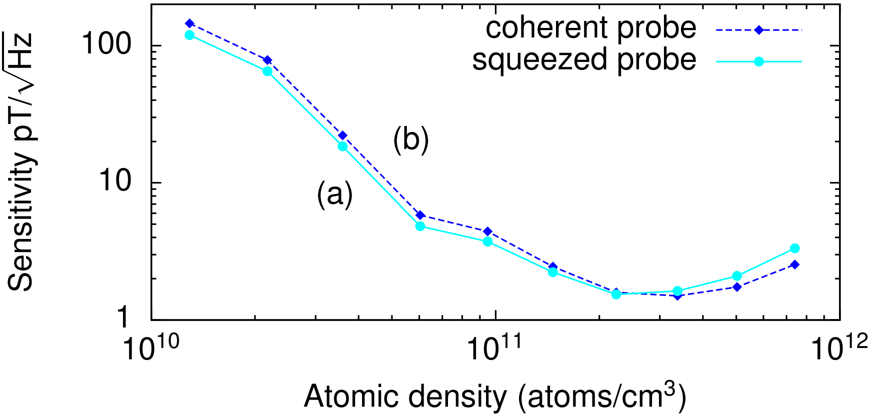


Note: Squeezed enhanced magnetometer was first demonstrated by Wolfgramm *et. al* Phys. Rev. Lett, **105**, 053601, 2010.

Magnetometer noise floor improvements

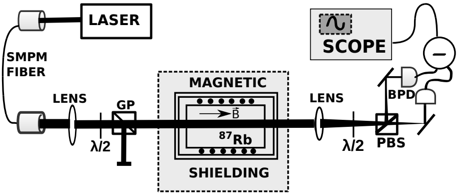


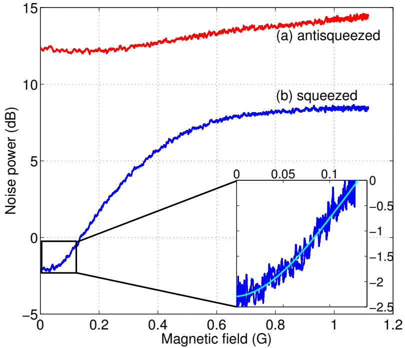
Magnetometer sensitivity vs atomic density



Squeezing vs magnetic field

Spectrum analyzer settings: Central frequency = 1 MHz, VBW = 3 MHz, RBW = 100 kHz





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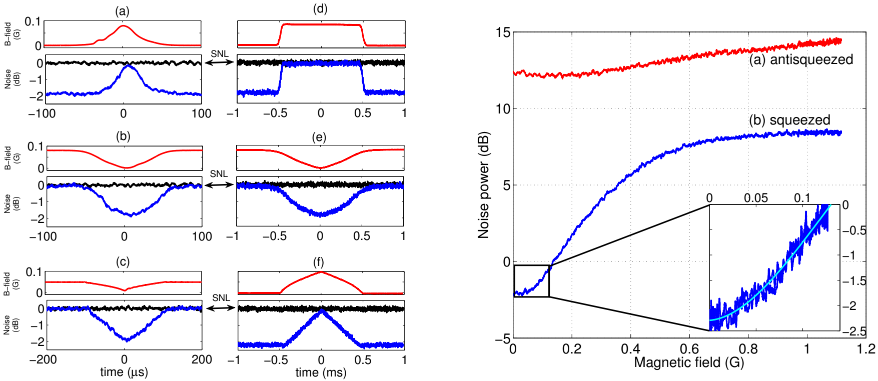

Summary

- Squeezing is exiting
- Many applications benefit from squeezing
  - interferometers
  - gravitational antennas
  - magnetometers
  - quantum memory
- There is still a lot of interesting physics to do