

Quantum enhanced magnetometer and squeezed state of light tunable filter

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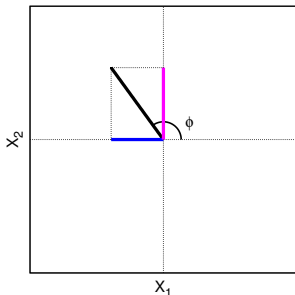
October 5, 2012

Transition from classical to quantum field

Classical analog

- Field amplitude a
- Field real part
 $X_1 = (a^* + a)/2$
- Field imaginary part
 $X_2 = i(a^* - a)/2$

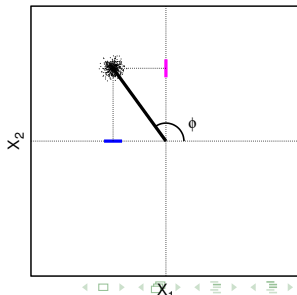
$$E(\phi) = |a|e^{-i\phi} = X_1 + iX_2$$



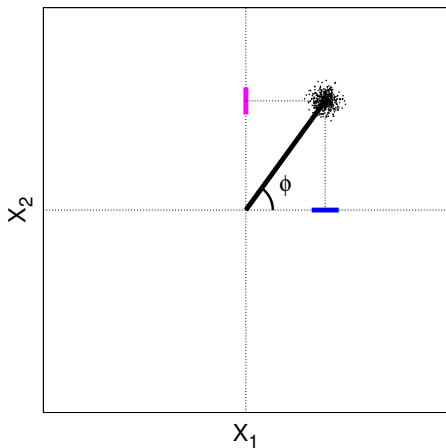
Quantum approach

- Field operator \hat{a}
- Amplitude quadrature
 $\hat{X}_1 = (\hat{a}^\dagger + \hat{a})/2$
- Phase quadrature
 $\hat{X}_2 = i(\hat{a}^\dagger - \hat{a})/2$

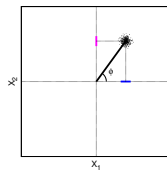
$$\hat{E}(\phi) = \hat{X}_1 + i\hat{X}_2$$



Squeezed quantum states zoo

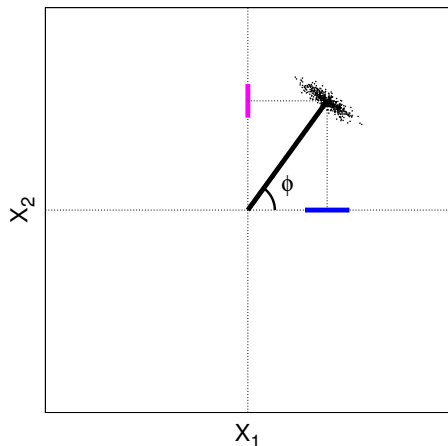


Unsqueezed
coherent

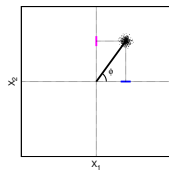


Notice $\Delta X_1 \Delta X_2 \geq \frac{1}{4}$

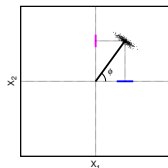
Squeezed quantum states zoo



Unsqueezed
coherent

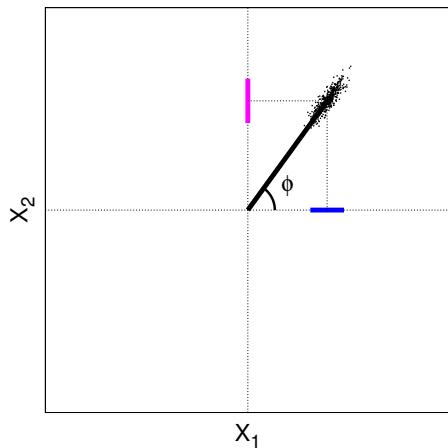


Amplitude
squeezed



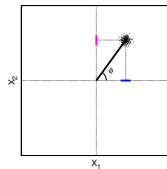
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Squeezed quantum states zoo

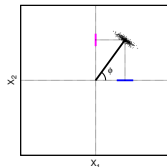


Notice $\Delta X_1 \Delta X_2 \geq \frac{1}{4}$

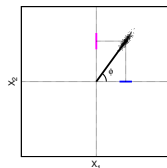
Unsqueezed
coherent



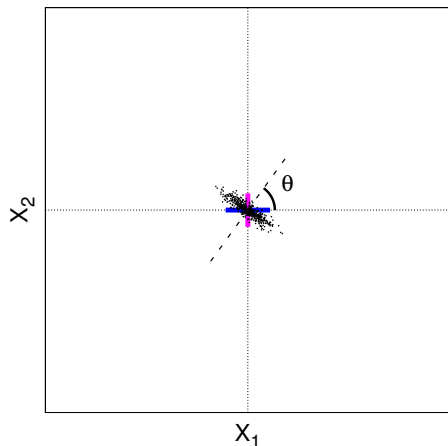
Amplitude
squeezed



Phase
squeezed

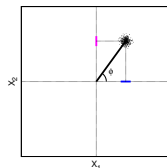


Squeezed quantum states zoo

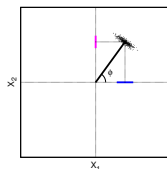


Notice $\Delta X_1 \Delta X_2 \geq \frac{1}{4}$

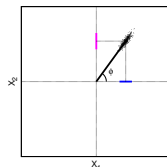
Unsqueezed
coherent



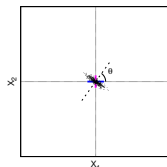
Amplitude
squeezed



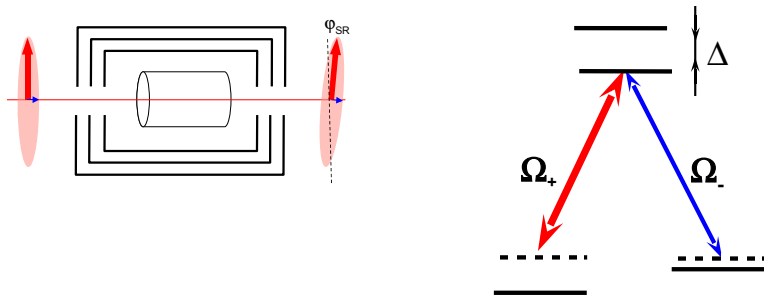
Phase
squeezed



Vacuum
squeezed



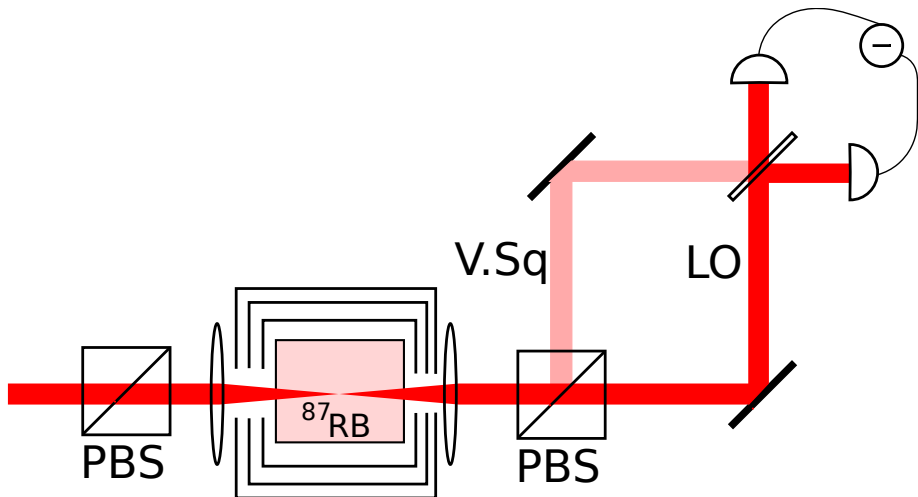
Self-rotation of elliptical polarization in atomic medium



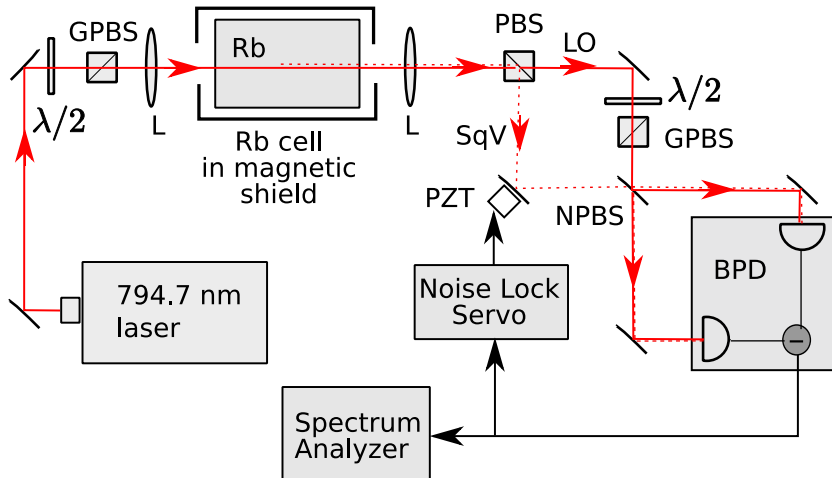
A.B. Matsko et al., PRA 66, 043815 (2002): theoretically prediction of 4-6 dB noise suppression

$$a_{out} = a_{in} + \frac{igL}{2}(a_{in}^\dagger - a_{in}) \quad (1)$$

Simplified setup



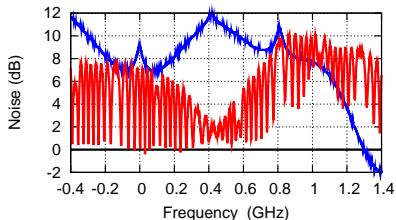
Setup



Noise contrast vs detuning in hot ^{87}Rb vacuum cell

$$F_g = 2 \rightarrow F_e = 1, 2$$

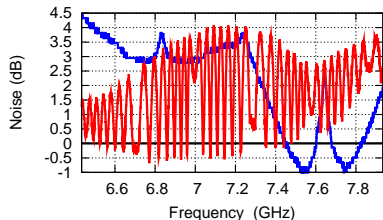
Noise vs detuning



Transmission — PSR noise

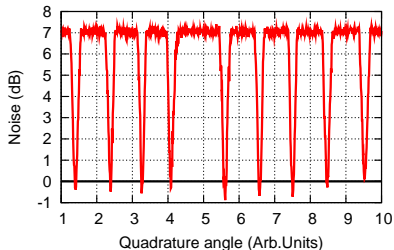
$$F_g = 1 \rightarrow F_e = 1, 2$$

Noise vs detuning

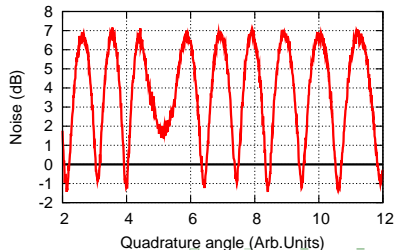


Transmission — PSR noise

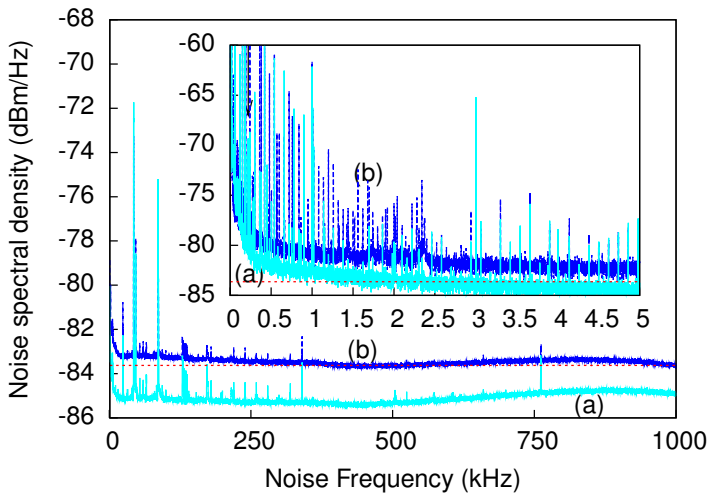
Noise vs quadrature angle



Noise vs quadrature angle

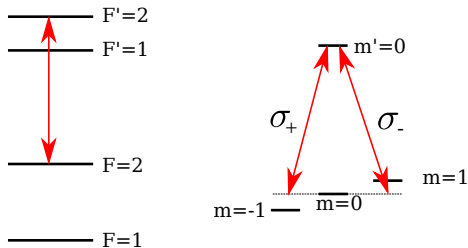


Atomic low frequency squeezing source

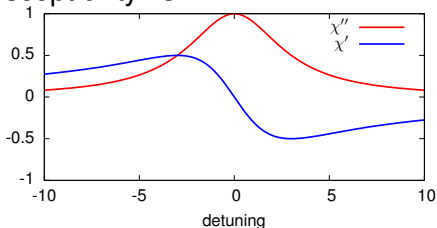


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

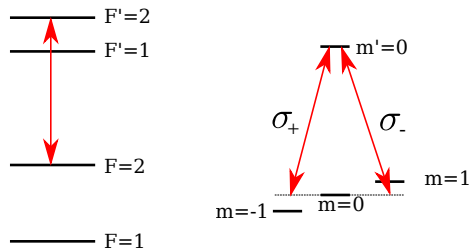


Susceptibility vs B

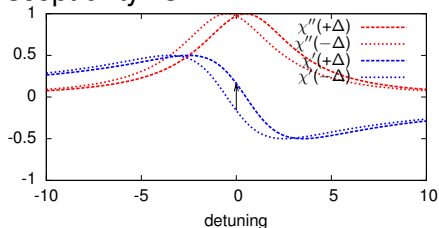


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

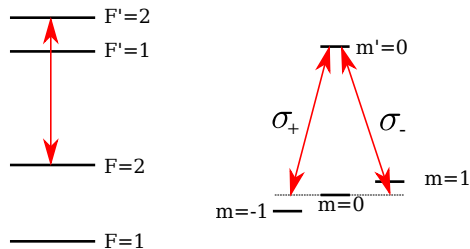


Susceptibility vs B

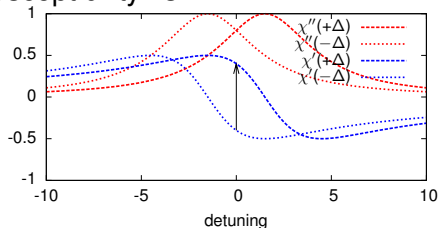


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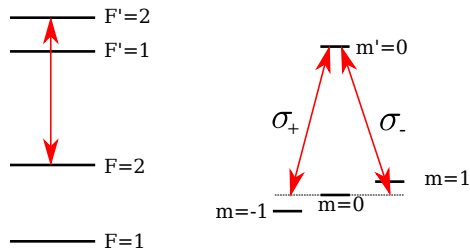


Susceptibility vs B

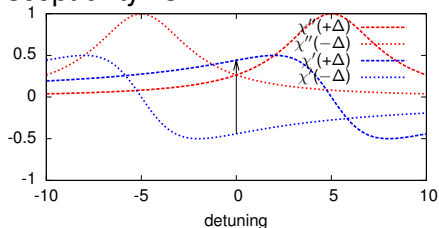


Optical magnetometer based on Faraday effect

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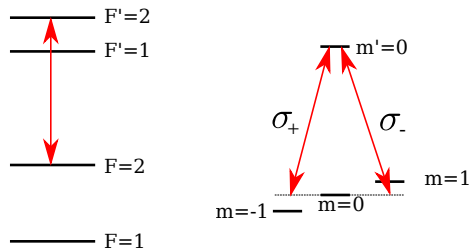


Susceptibility vs B

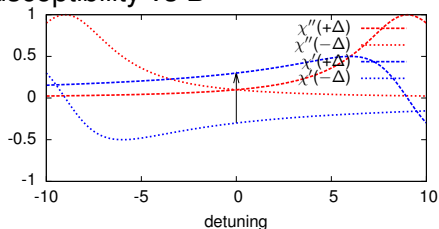


Optical magnetometer based on Faraday effect

^{87}Rb D₁ line

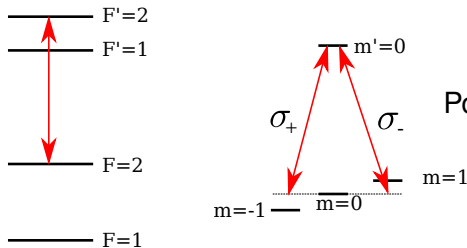


Susceptibility vs B

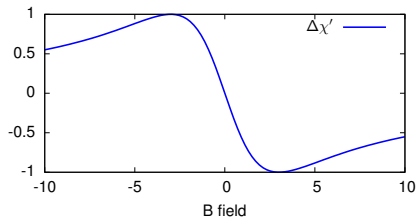


Optical magnetometer based on Faraday effect

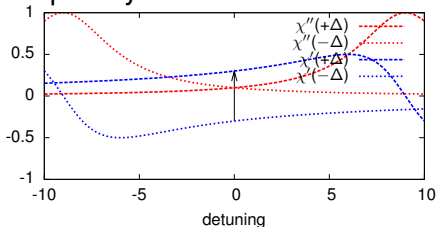
^{87}Rb D₁ line



Polarization rotation vs B

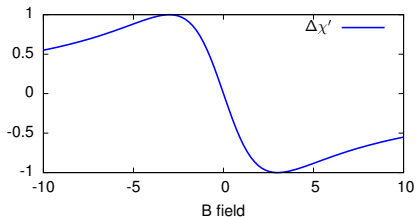


Susceptibility vs B

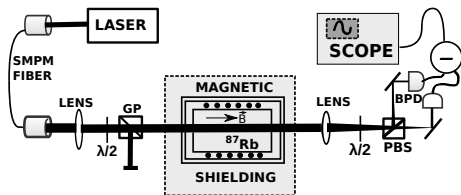


Optical magnetometer and non linear Faraday effect

Naive model of rotation

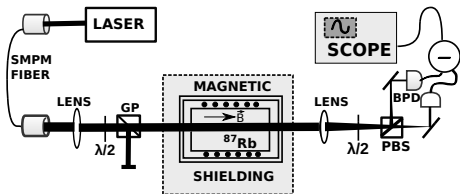
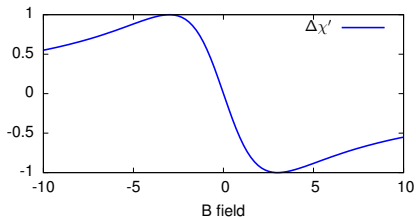


Experiment

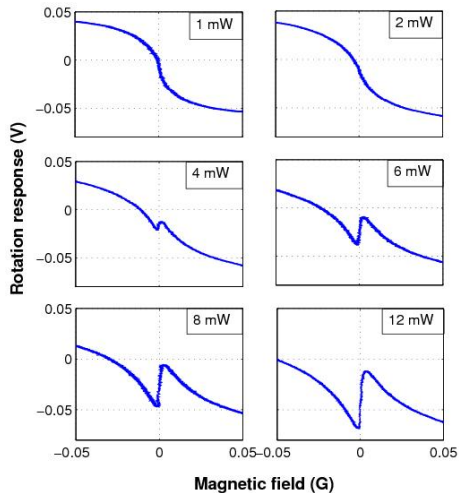


Optical magnetometer and non linear Faraday effect

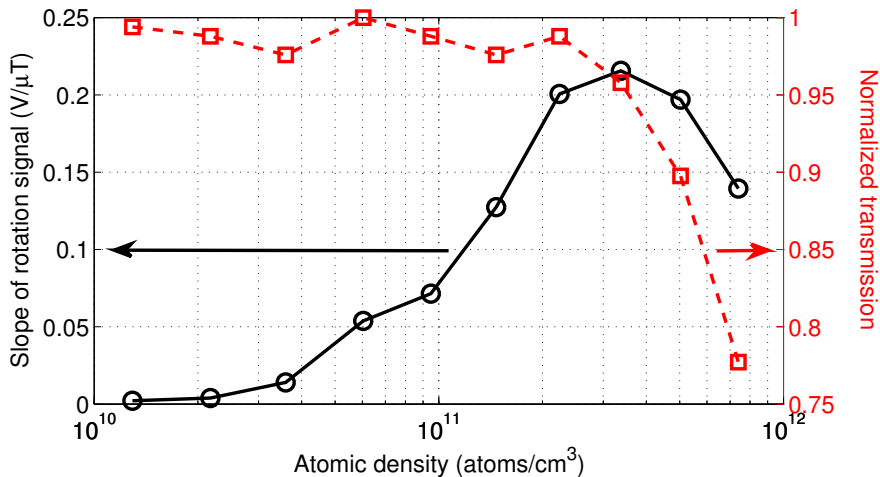
Naive model of rotation



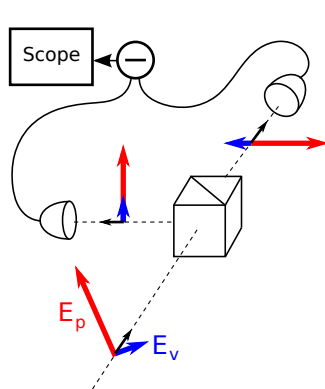
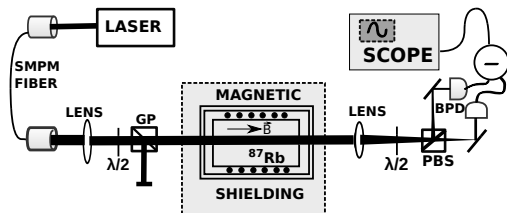
Experiment



Magnetometer response vs atomic density



Shot noise limit of the magnetometer

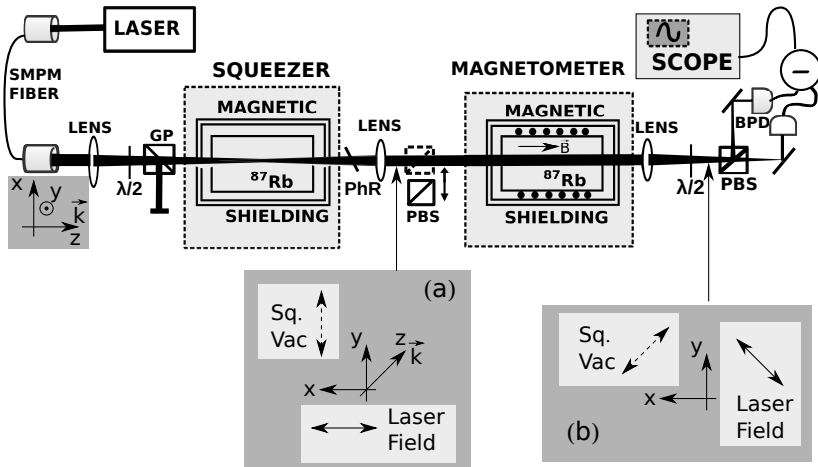


$$S = |E_p + E_v|^2 - |E_p - E_v|^2$$

$$S = 4E_p E_v$$

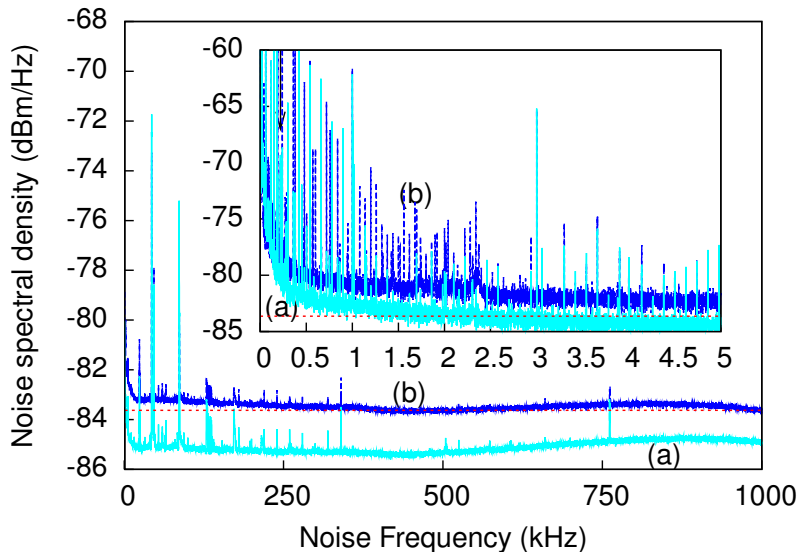
$$\langle \Delta S \rangle \sim E_p \langle \Delta E_v \rangle$$

Squeezed enhanced magnetometer setup

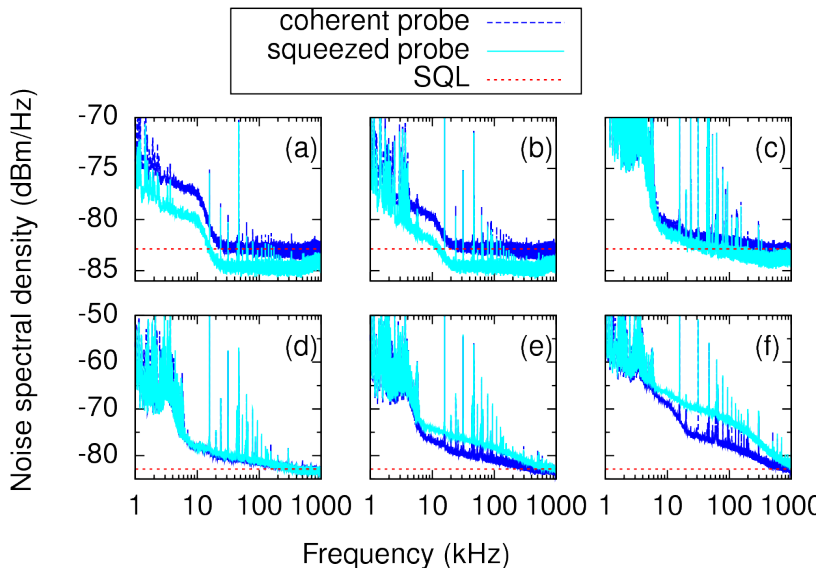


Note: Squeezed enhanced magnetometer was first demonstrated by Wolfgramm *et. al* Phys. Rev. Lett, **105**, 053601, 2010.

Magnetometer noise floor improvements

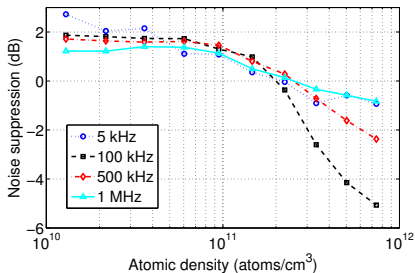


Magnetometer noise spectra

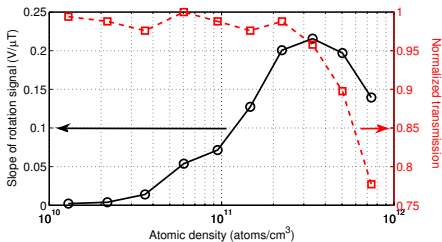


Noise suppression and response vs atomic density

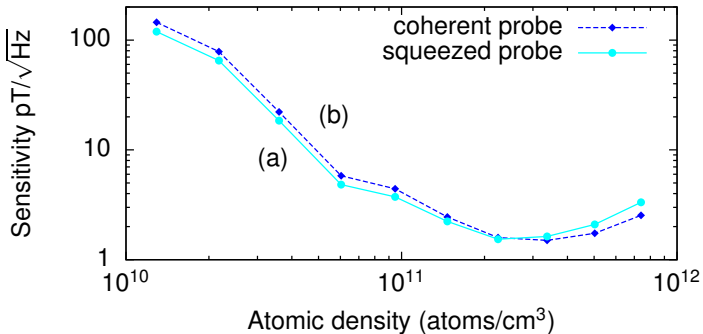
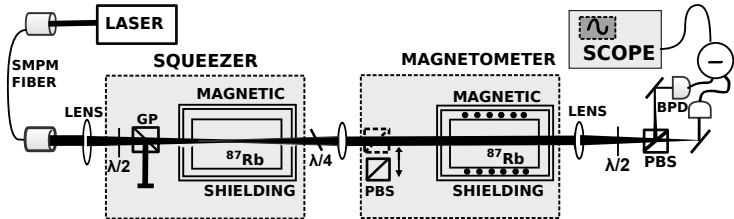
Noise suppression



Response



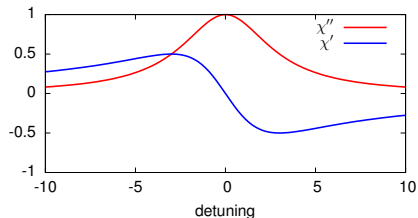
Magnetometer with squeezing enhancement



Light group velocity expression

$$\text{Group velocity } v_g = \frac{c}{\omega \frac{\partial n}{\partial \omega}}$$

Susceptibility

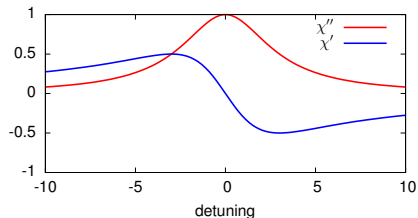


Rotation vs B field

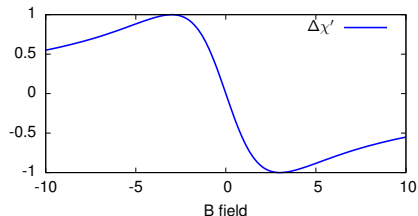
Light group velocity expression

$$\text{Group velocity } v_g = \frac{c}{\omega \frac{\partial n}{\partial \omega}}$$

Susceptibility



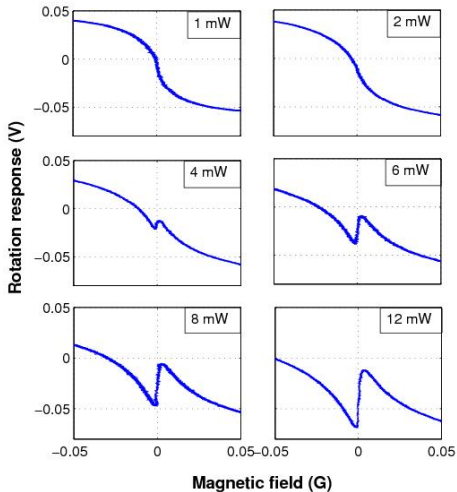
Rotation vs B field



Light group velocity estimate

$$\text{Group velocity } v_g = \frac{c}{\omega \frac{\partial n}{\partial \omega}}$$

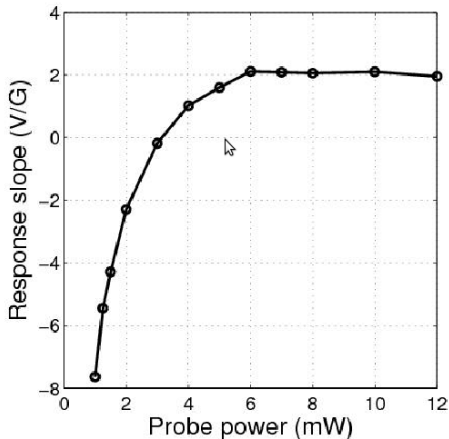
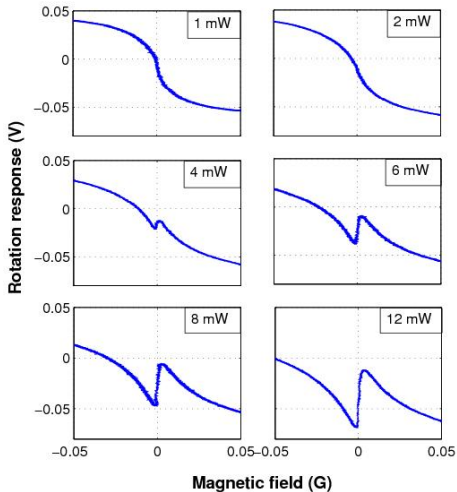
$$\text{Delay } \tau = \frac{L}{v_g} \sim \frac{\partial n}{\partial \omega} \sim \frac{\partial R}{\partial B}$$



Light group velocity estimate

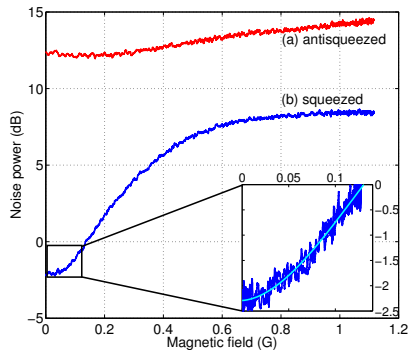
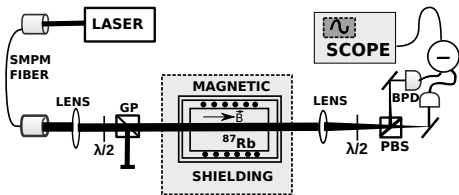
$$\text{Group velocity } v_g = \frac{c}{\omega \frac{\partial n}{\partial \omega}}$$

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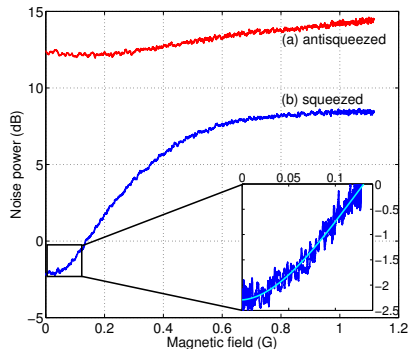
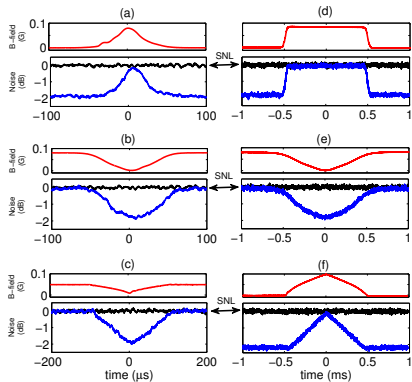
Squeezing vs magnetic field

Spectrum analyzer settings: Central frequency = 1 MHz, VBW = 3 MHz, RBW = 100 kHz

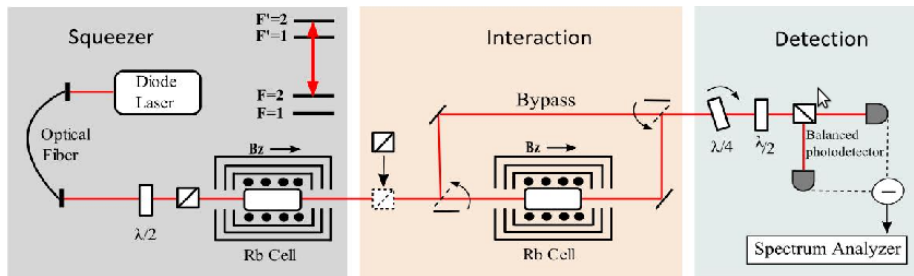


Squeezing vs magnetic field

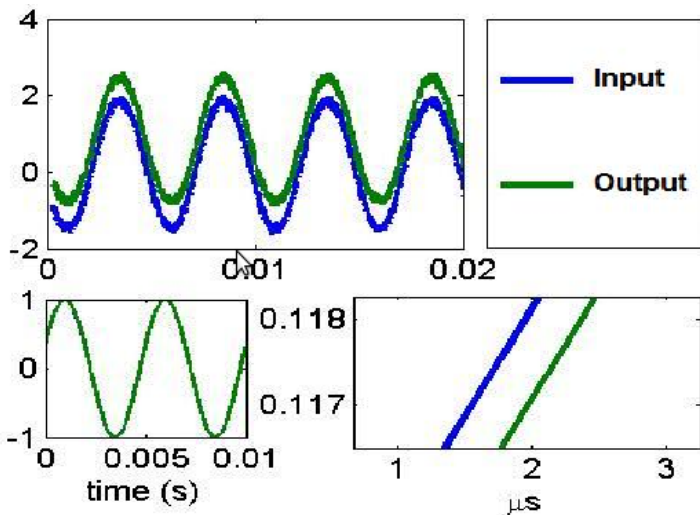
Spectrum analyzer settings: Central frequency = 1 MHz, VBW = 3 MHz, RBW = 100 kHz



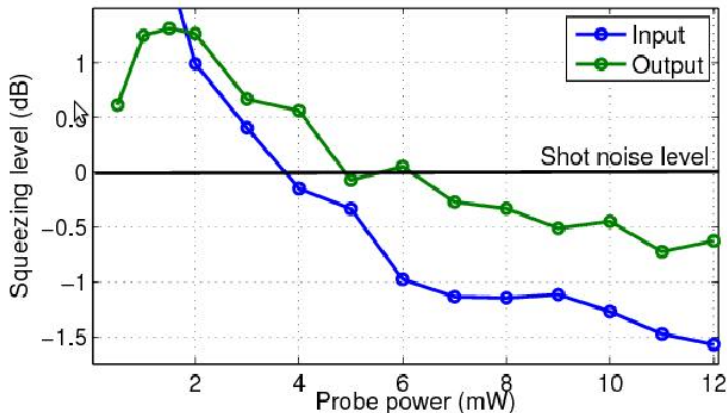
Time advancement setup



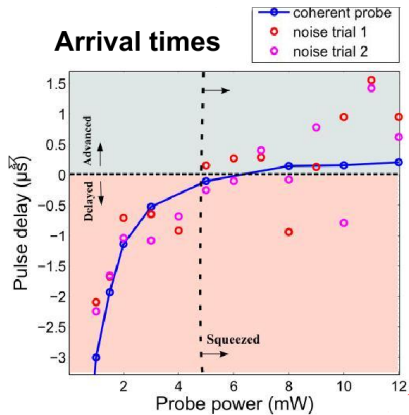
Squeezing modulation and time advancement



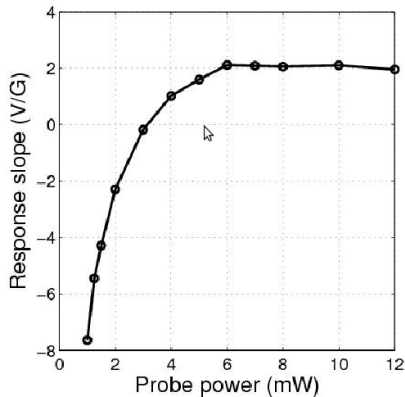
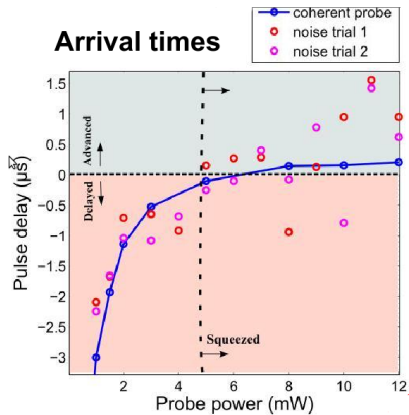
Squeezing after advancement cell



Advancement vs power

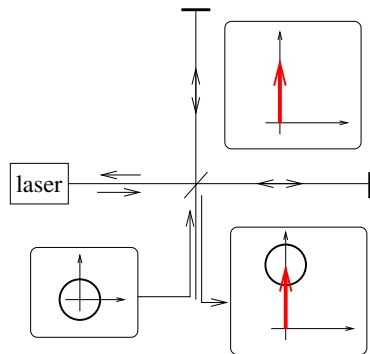


Advancement vs power

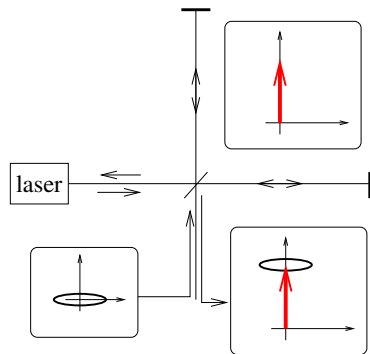


Quantum limited interferometers revisited

Vacuum input



Squeezed input



Limiting noise - Quantum Optical noise

Next generation of LIGO will be
quantum optical noise limited at almost all detection frequencies.

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shot noise

Uncertainty in number of photons

$$h \sim \sqrt{\frac{1}{P}} \quad (2)$$

Limiting noise - Quantum Optical noise

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shot noise

Uncertainty in number of photons

$$h \sim \sqrt{\frac{1}{P}} \quad (2)$$

radiation pressure noise

Photons impart momentum to mirrors

$$h \sim \sqrt{\frac{P}{M^2 f^4}} \quad (3)$$

Limiting noise - Quantum Optical noise

Next generation of LIGO will be **quantum optical noise limited** at almost all detection frequencies.

shot noise

Uncertainty in number of photons

$$h \sim \sqrt{\frac{1}{P}} \quad (2)$$

radiation pressure noise

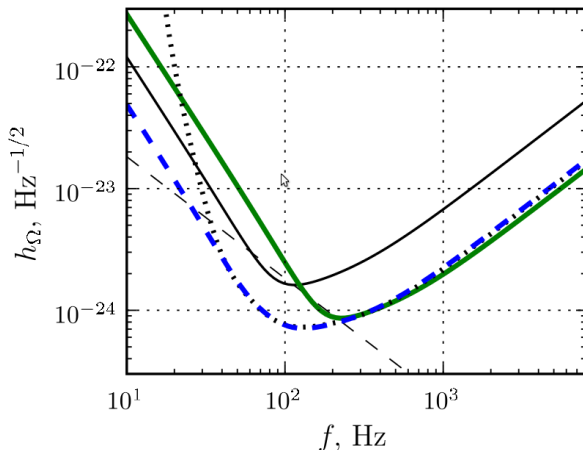
Photons impart momentum to mirrors

$$h \sim \sqrt{\frac{P}{M^2 f^4}} \quad (3)$$

There is no optimal light power to suit all detection frequency.
Optimal power depends on desired detection frequency.

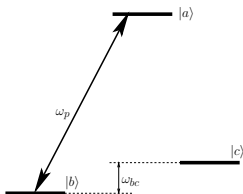
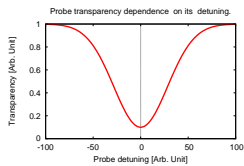
Interferometer sensitivity improvement with squeezing

Projected advanced LIGO sensitivity

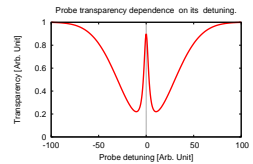
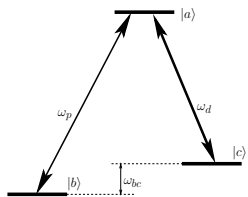
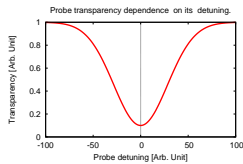


F. Ya. Khalili Phys. Rev. D 81, 122002 (2010)

EIT filter



EIT filter

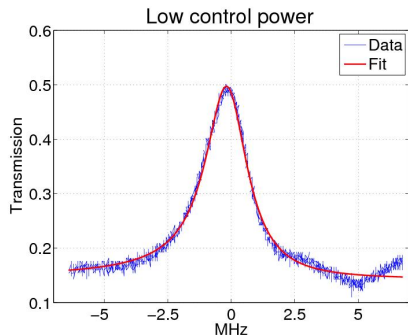
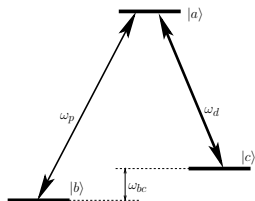


Squeezing and EIT filter

$$\begin{pmatrix} V_1^{out} \\ V_2^{out} \end{pmatrix} = \begin{pmatrix} A_+^2 & A_-^2 \\ A_-^2 & A_+^2 \end{pmatrix} \begin{pmatrix} V_1^{in} \\ V_2^{in} \end{pmatrix} + [1 - (A_+^2 + A_-^2)] \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\varphi_{\pm} = \frac{1}{2} (\Theta_+ \pm \Theta_-)$$

$$A_{\pm} = \frac{1}{2} (T_+ \pm T_-)$$

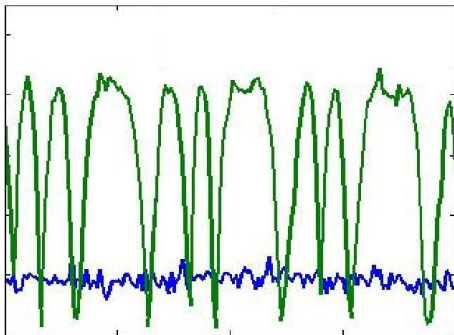
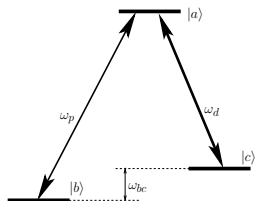


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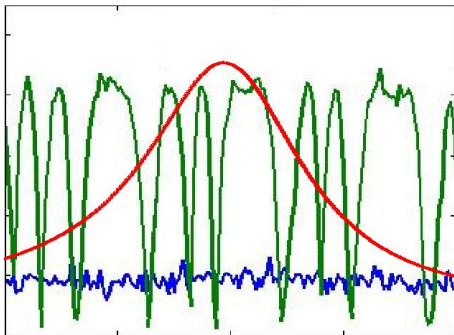
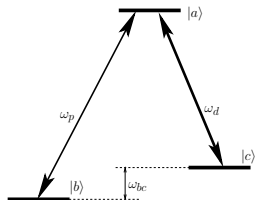


Squeezing and EIT filter

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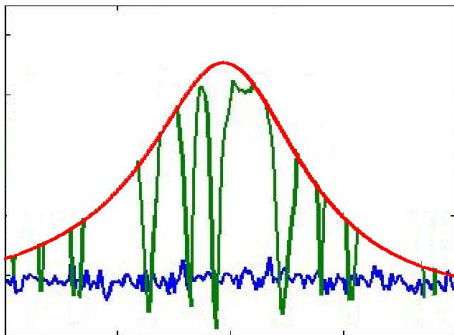
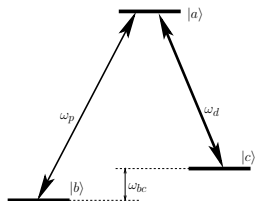


Squeezing and EIT filter

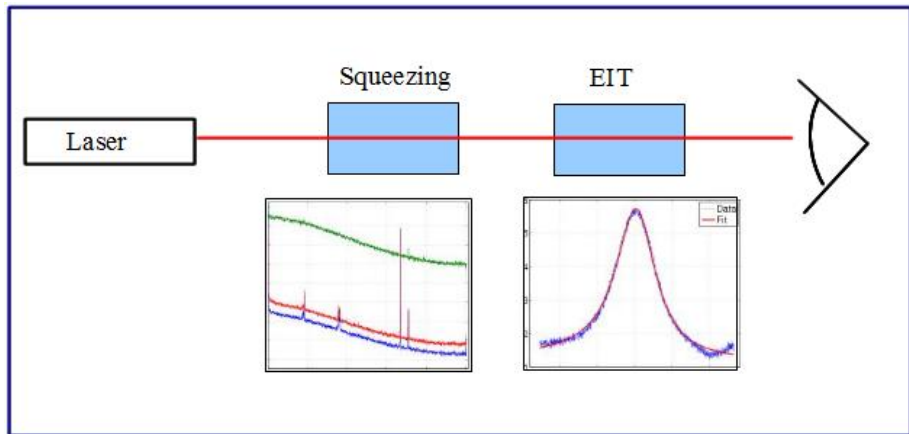
$$\begin{pmatrix} V_1^{out} \\ V_2^{out} \end{pmatrix} = \begin{pmatrix} A_+^2 & A_-^2 \\ A_-^2 & A_+^2 \end{pmatrix} \begin{pmatrix} V_1^{in} \\ V_2^{in} \end{pmatrix} + [1 - (A_+^2 + A_-^2)] \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\varphi_{\pm} = \frac{1}{2} (\Theta_+ \pm \Theta_-)$$

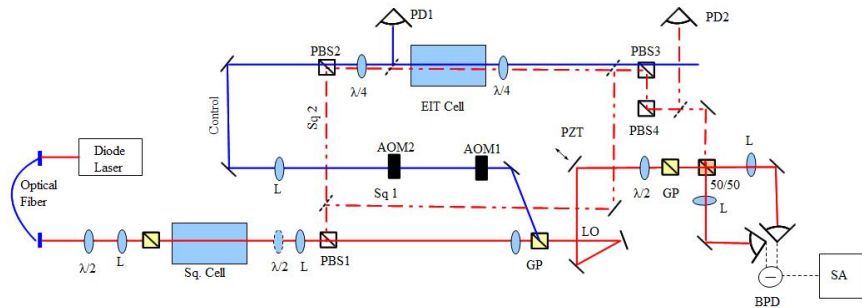
$$A_{\pm} = \frac{1}{2} (T_+ \pm T_-)$$



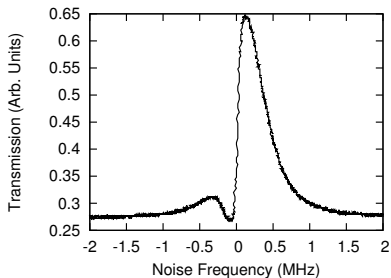
Squeezing and EIT filter setup



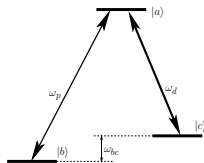
Squeezing and EIT filter setup



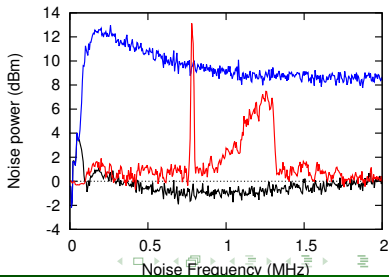
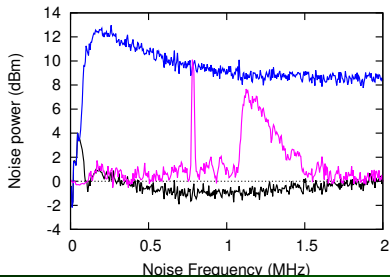
EIT filter and measurements without light



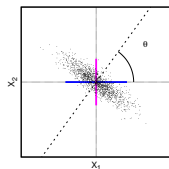
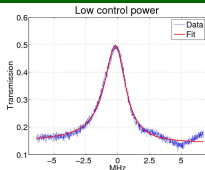
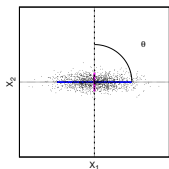
Coherent signal



Signal in the noise quadratures



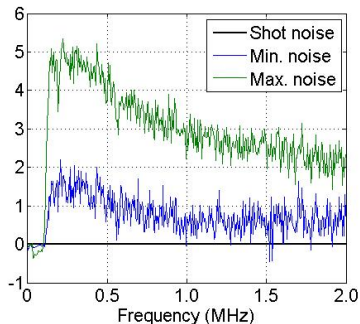
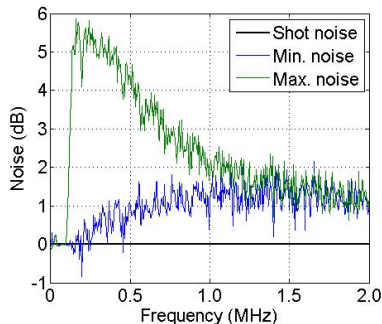
Squeezing angle rotation



$$\begin{pmatrix} V_{1out} \\ V_{2out} \end{pmatrix} = \begin{pmatrix} \cos^2 \varphi_+ & \sin^2 \varphi_+ \\ \sin^2 \varphi_+ & \cos^2 \varphi_+ \end{pmatrix} \begin{pmatrix} A_+^2 & A_-^2 \\ A_-^2 & A_+^2 \end{pmatrix} \begin{pmatrix} V_{1in} \\ V_{2in} \end{pmatrix} + [1 - (A_+^2 + A_-^2)] \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

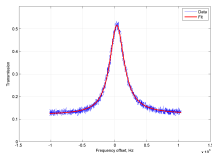
Locked at 300kHz

Locked at 1200kHz

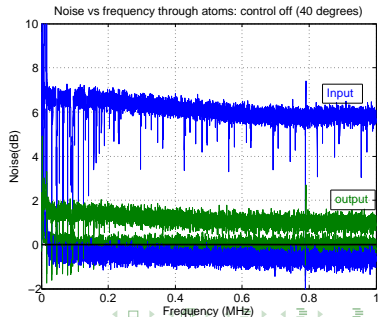
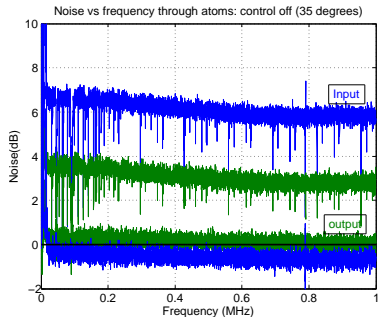
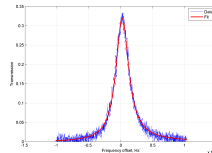


Narrower filter

$T=35^{\circ}\text{C}$, no control
transmission 42%

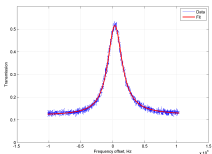


$T=40^{\circ}\text{C}$, no control
transmission 17%

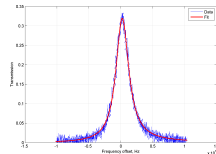


Narrower filter

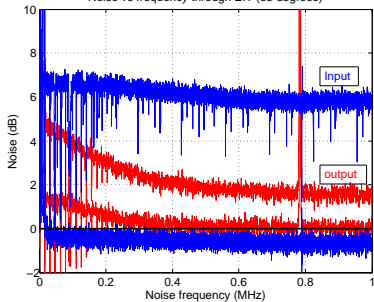
$T=35^{\circ}\text{C}$, no control
transmission 42%



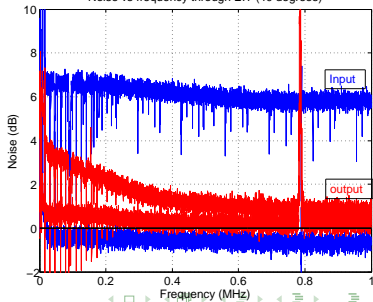
$T=40^{\circ}\text{C}$, no control
transmission 17%



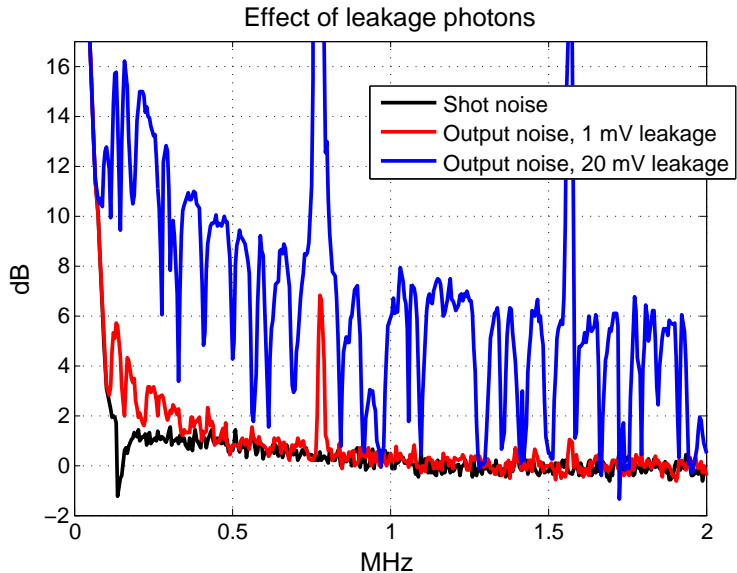
Noise vs frequency through EIT (35 degrees)



Noise vs frequency through EIT (40 degrees)

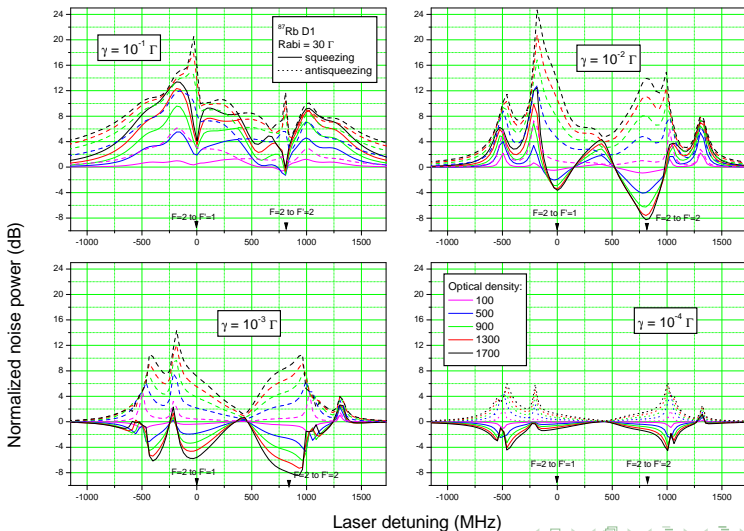


Excess noise and leakage



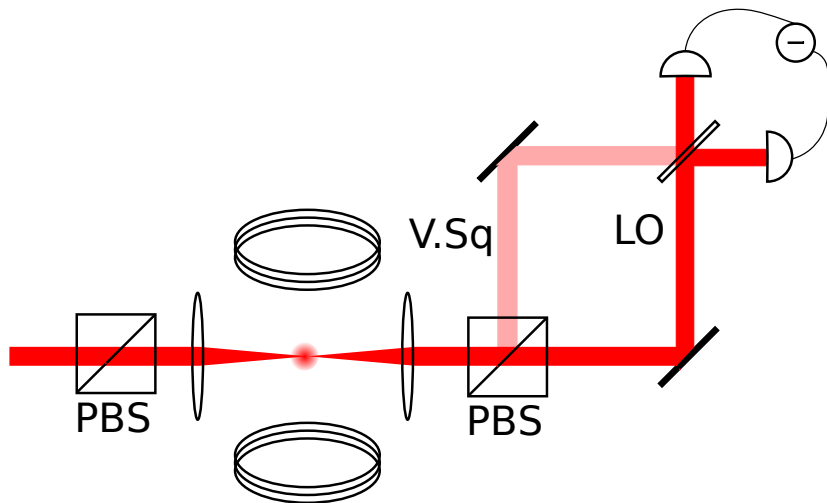
Theoretical prediction for MOT squeezing with ^{87}Rb

$F_g = 2 \rightarrow F_e = 1, 2$ high optical density is very important

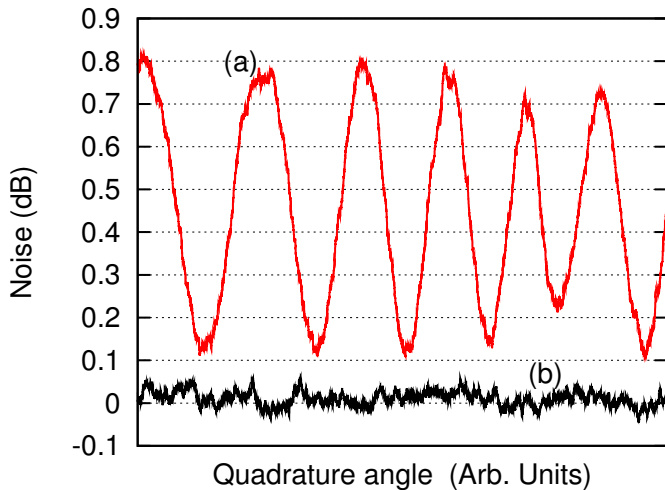


MOT squeezer

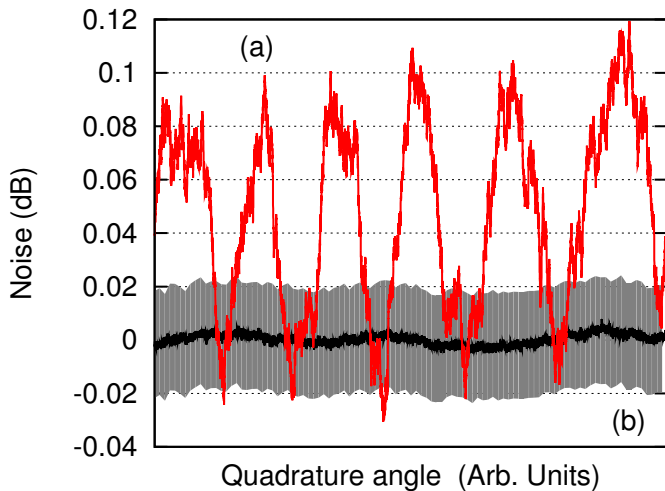
Cloud size = 1 mm, $T = 200 \mu\text{K}$, $N = 7 \times 10^9 \text{ 1/cm}^3$,
OD = 2, beam size = 0.1 mm, 10^5 interacting atoms



Noise contrast in MOT with ^{87}Rb $F_g = 2 \rightarrow F_e = 1$



Squeezing in MOT with ^{87}Rb $F_g = 2 \rightarrow F_e = 1$



People

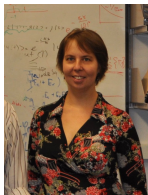
Travis Horrom and Gleb Romanov



Robinjeet Singh, LSU



Irina Novikova



Jonathan P. Dowling, LSU



Summary

- We demonstrate fully atomic squeezed enhanced magnetometer
- Magnetometer noise floor lowered in the range from several kHz to several MHz
- Demonstrated sensitivity as low as $1 \text{ pT}/\sqrt{\text{Hz}}$ in our particular setup
- First demonstration of superluminal squeezing propagation with $v_g = c/2000$ or time advancement of $0.5 \mu\text{S}$

For more details:

- T. Horrom, et al. “Quantum Enhanced Magnetometer with Low Frequency Squeezing”, **PRA**, 86, 023803, (2012).
- T. Horrom, et al. “All-atomic generation and noise-quadrature filtering of squeezed vacuum in hot Rb vapor”, arXiv:1204.3967.

Support from

