

Molecules (at least simple ones)

We saw that there is no analytical solution for He (1 nucleus and 2 electrons), so there is little hope to have solutions for a general molecule (many electrons and at least two nucleuses).

So in general we would be concerned with a question: can the system be bound? i.e. is combined ground state energy lower than constituting parts?

If we somehow can predict the ground level Energy behavior vs R (nucleus separation)



minima at infinity \Rightarrow no bound state

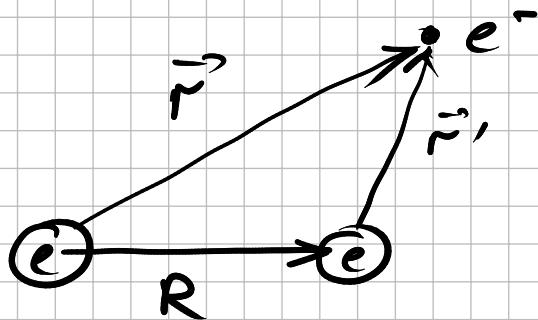
localized minima \Rightarrow bound state exist

Hydrogen molecule ion

this is the next "easiest" case

Born-Oppenheimer approximation

Nucleus are heavy \Rightarrow slow
so we will disregard their
relative motion with
respect to the center of mass.



$$\hat{H} = \frac{\vec{p}^2}{2me} - \frac{e^2}{r} - \frac{e^2}{r'} + \frac{e^2}{R}$$

nucleus repulsion
for any ψ

R is constant
in our model

$$\begin{aligned} \langle \frac{e^2}{R} \rangle &= \frac{e^2}{R} = \frac{e^2}{a_0} \frac{a_0}{R} \\ &= -\frac{2a_0}{R} \cdot E_g (z=1) \end{aligned}$$

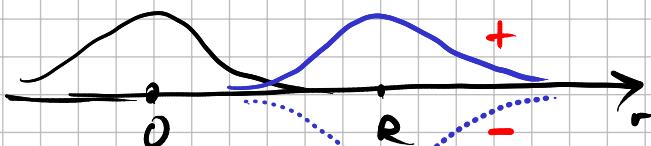
As usual, the art
is in choosing good
anzatz ψ .

We will use

$$\psi = C (\Psi_{1,0}(r) \pm \Psi_{1,0}(r')) , \quad \vec{r}' = \vec{r} - \vec{R}$$

Hydrogen ground
level eigenstate

Hydrogen ground energy



right now it is unclear which + or - is better. But if we think that two nuclei repel each other \Rightarrow we need electron in between, so '+' is more promising

Ψ in the form of $(\psi_{10}(r) + \psi_{10}(r'))$ is what quantum chemist call a linear combination of atomic orbitals (LCAO).

Our 1st task to normalize Ψ

$$\langle \Psi | \Psi \rangle = c^2 \left(\underbrace{\langle \psi_{10}(r) | \psi_{10}(r) \rangle}_{=1} + \underbrace{\langle \psi_{10}(r') | \psi_{10}(r') \rangle}_{=1} + \pm 2 \langle \psi_{10}(r) | \psi_{10}(r') \rangle \right)$$

$$= c^2 2 (1 \pm \underbrace{\langle \psi_{10}(r) | \psi_{10}(r') \rangle}_{\text{overlap integral}})$$

overlap integral which we label I

$$\text{Recall } \psi_{10}(r) = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}$$

Calculation of I is tedious (see Ch.8 Griffiths)

$$I = e^{-R/a} \left[1 + \left(\frac{R}{a}\right) + \frac{1}{3} \left(\frac{R}{a}\right)^2 \right]$$

$$c = \frac{1}{2(1 \pm I)}$$

Next we need the expectation value of truncated \hat{H}_T (no nucleus repulsion) $\frac{p^2}{2m} - \frac{e^2}{r} - \frac{e^2}{r'}$

Observe that

$$\begin{aligned}\hat{H}_T |\Psi_{10}(r)\rangle &= \left(\underbrace{\frac{p^2}{2m} - \frac{e^2}{r}}_E - \frac{e^2}{r'} \right) |\Psi_{10}(r)\rangle = \\ &\quad E \text{ of } M_{\text{hydrogen}} \\ &= E_g(z=1) |\Psi_{10}(r)\rangle - \frac{e^2}{r'} |\Psi_{10}(r)\rangle \quad \text{no prime} \\ \text{Similarly} \quad & \quad \text{ground energy of Hydrogen} \quad \text{prime} \\ \hat{H}_T |\Psi_{10}(r')\rangle &= E_g(z=1) |\Psi_{10}(r')\rangle - \frac{e^2}{r} |\Psi_{10}(r')\rangle\end{aligned}$$

$$\begin{aligned}\langle \psi | \hat{H}_T | \psi \rangle &= \langle \psi | \hat{H}_T | \underbrace{c(\Psi_{10}(r) \pm \Psi_{10}(r'))}_{\psi} \rangle = \\ &= E_g(z=1) - c^2 e^2 \left(\langle \Psi_{10}(r) \pm \Psi_{10}(r') | \left(\frac{1}{r} |\Psi_{10}(r)\rangle \pm \frac{1}{r'} |\Psi_{10}(r')\rangle \right) \right) \\ &= E_g(z=1) - c^2 e^2 \left[\langle \Psi_{10}(r) | \frac{1}{r} |\Psi_{10}(r)\rangle + \langle \Psi_{10}(r') | \frac{1}{r'} |\Psi_{10}(r')\rangle \right. \\ &\quad \left. \langle \Psi_{10}(r') | \frac{1}{r} |\Psi_{10}(r)\rangle \pm \langle \Psi_{10}(r') | \frac{1}{r'} |\Psi_{10}(r)\rangle \right] \quad \text{variable relabel}\end{aligned}$$

$$\langle H_r \rangle = E_g(z=1) - \frac{2C^2 e^2}{2a} \left[a \langle \Psi_{10}(r) | \frac{1}{r} | \Psi_{10}(r) \rangle \pm a \langle \Psi_{10}(r) | \frac{1}{r} | \Psi_{10}(r') \rangle \right]$$

D - direct integral

X - exchange integral

$$D = \frac{a}{R} - \left(1 + \frac{a}{R}\right) e^{-2R/a}$$

$$X = \left(1 + \frac{R}{a}\right) e^{-R/a}$$

$$\langle H \rangle = \langle H_r \rangle - \underbrace{\frac{2a}{R} E_g(z=1)}_{\text{nuclear repulsion}} = \left[\left(1 + 2 \frac{D \mp X}{1 \pm \lambda} \right) - \frac{2a}{R} \right] E_g(z=1)$$

notice that we can introduce unitless

parameter $\lambda = \frac{R}{a}$ and search for minimum

λ_{\min} for (+) case $\frac{R}{a} = 2.4 \Rightarrow R = 1.3 \text{ \AA}$

$$\langle H \rangle - E_g(z=1) = 1.8 \text{ eV}$$

https://quantummechanics.ucsd.edu/ph130a/130_notes/node398.html

