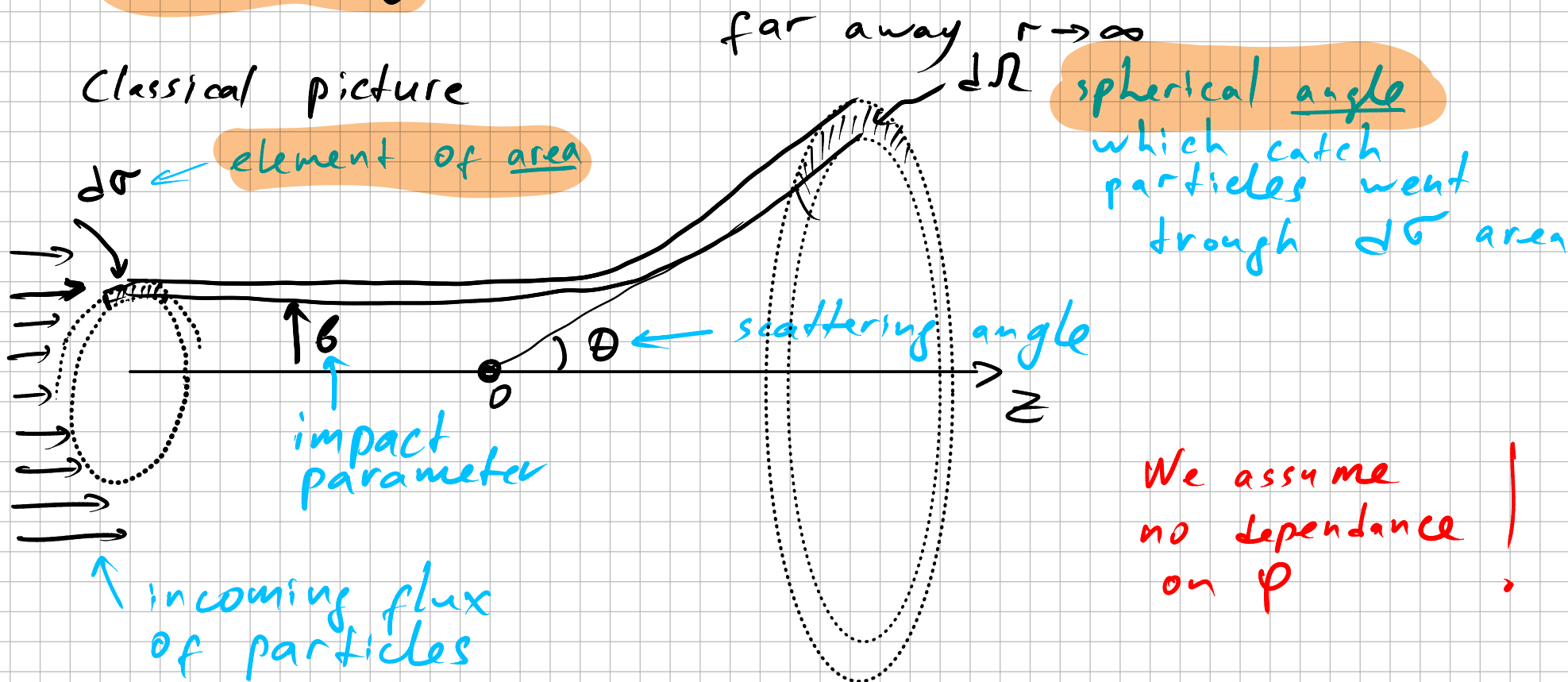


Scattering

Goal: find size and shape of the object



We assume no dependence on φ !

$d\sigma$ and $d\Omega$ seem to be linked, i.e. proportional

$$d\sigma = D(\theta) d\Omega$$

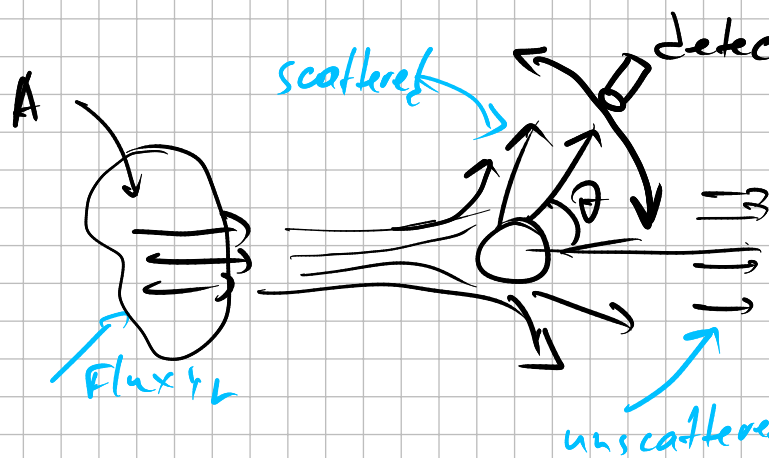
Canonical name: Differential cross-section often labeled as $\frac{d\sigma}{d\Omega}$

$$d\sigma = \frac{d\sigma}{d\Omega} d\Omega$$

good luck to understand this form of equation

Experimentally we usually have access to

$$D(\theta) d\Omega$$



detector of fixed size $d\Omega$
moving along variable θ

of "clicks" in detector
proportional to

$$\text{Flux} \cdot d\Omega = \text{Flux} \cdot D(\theta) d\Omega = dN_{\text{detected Scattered}}$$

particle
unit area · unit time

$$\int d\sigma = \int D(\theta) d\Omega$$

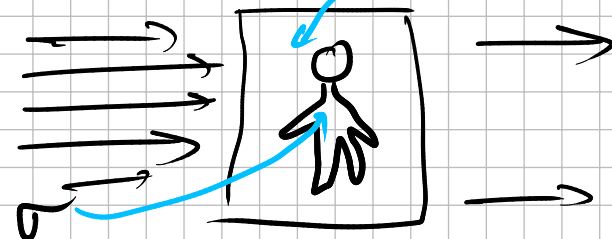
$$\sigma = \int D(\theta) d\Omega = \int \frac{dN_{\text{scattered}}}{\text{Flux}} = \frac{N_{\text{scattered}}}{\text{Flux}}$$

$$N_{\text{in}} = A \cdot \text{Flux}$$

↑ # particles
which we send
in

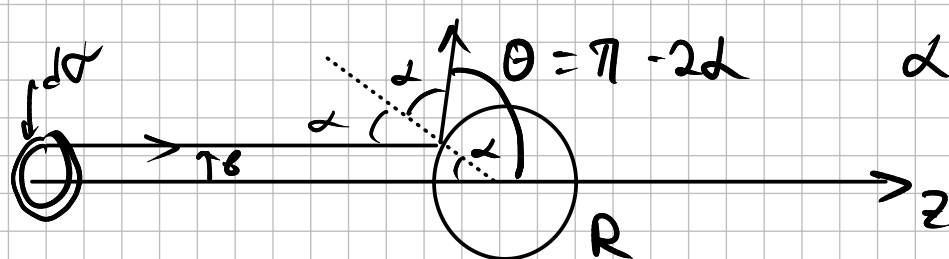
$$\frac{N_{\text{sc}}}{N_{\text{in}}} = \frac{\sigma}{A}$$

size of the "shadow"
missing particles
from $\theta = 0$

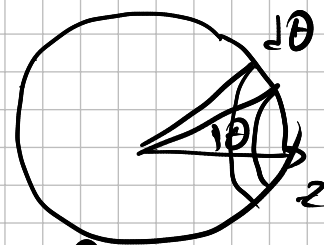


$D(\theta)$ carry information
about object and interaction
forces

Example: scattering on a sphere if $b > R \Rightarrow \theta = 0$



$$\alpha = \arcsin \frac{b}{R}$$



$$d\sigma = 2\pi b db = D(\theta) \cdot d\Omega = D(\theta) \cdot \underbrace{\frac{2\pi r \cdot r \sin \theta}{r^2}}_{d\Omega} d\theta$$

$$D(\theta) = \frac{d\sigma}{d\Omega} = \frac{2\pi b db}{2\pi \sin \theta d\theta} = \frac{b}{\sin \theta} \cdot \frac{db}{d\theta}$$

$$\left(\theta = \pi - 2\alpha = \pi - 2 \arcsin \frac{b}{R} \right)$$

$$d\theta = -2 \frac{1}{\sqrt{1 - \left(\frac{b}{R}\right)^2}} \cdot \frac{db}{R} = -2 \left| \frac{1}{\cos \alpha} \right| \frac{db}{R} = -2 \left| \frac{1}{\cos\left(\frac{\pi}{2} - \frac{\theta}{2}\right)} \right| \frac{db}{R}$$

$\uparrow \sin \alpha$
 $\uparrow \frac{\pi}{2} - \frac{\theta}{2}$

$$= -2 \left| \frac{1}{\sin \frac{\theta}{2}} \right| \frac{db}{R}$$

differentiate
LHS and RHS

$$\frac{db}{d\theta} = -\frac{R}{2} \left| \sin \frac{\theta}{2} \right| \Rightarrow$$

$$D(\theta) = \frac{b}{\sin \theta} \cdot \frac{db}{d\theta} = -\frac{b}{\sin \theta} \frac{R}{2} \left| \sin(\theta/2) \right| = -R^2 \frac{\overbrace{\cos(\theta/2) \sin(\theta/2)}^{\frac{1}{2} \sin \theta}}{2 \sin \theta}$$

$$D(\theta) = -\frac{R^2}{4}$$

Finally recall that $\int_0^\pi$

$$\sigma = \int D(\theta) d\Omega = - \int_0^\pi \frac{R^2}{4} 2\pi \sin \theta d\theta = \pi R^2 = \sigma$$

as expected
this "shadow" area
of a sphere

Important fact: Rutherford scattering of a charged particle
on a coulomb potential $\sim \frac{Qq}{r}$ gives $\sigma = \infty$

$$D(\theta) \sim \frac{1}{\sin^4(\theta/2)}$$

It was experimentally confirmed
with exact $D(\theta)$

Usually $\theta_{\text{scattering}}$ grows as impact parameter b decreases.

Think about $b \rightarrow \infty$ far away from collision $\theta \rightarrow 0$
 $b \rightarrow 0$ we guarantee to have a hit
 $\Rightarrow \theta$ grows