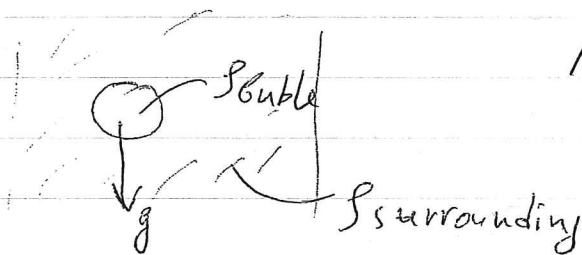


Lecture 24

gravitational mechanism of stellar pulsation

(PT)

g-mode - bubble slushing



Note that $p^{(b)} = p^{(s)}$ otherwise bubble would compress or expand until this condition is met

Buoyant force

$$f_{\text{net}} = (-\rho^{(b)}g + \rho^{(s)}g)V = (\rho^{(s)} - \rho^{(b)})gV$$

gravity pull down

once bubble displaced δr
we write Taylor expansion

$$\rho^{(s)} = \rho^{(s)}_{\text{initial}} + \frac{d\rho^{(s)}}{dr} \delta r$$

$$\rho^{(b)} = \rho^{(b)}_{\text{initial}} + \frac{d\rho^{(b)}}{dr} \delta r$$

assuming that $\rho^{(s)}_{\text{initial}} = \rho^{(b)}_{\text{initial}}$ i.e. they have the same starting point from which bubble emerge

$$f_{\text{net}} = g \left(\frac{d\rho^{(s)}}{dr} \delta r - \frac{d\rho^{(b)}}{dr} \delta r \right) V =$$

$$= \left(\frac{d\rho^{(s)}}{dr} \delta r - \frac{d\rho^{(b)}}{d\rho^{(b)}} \cdot \frac{d\rho^{(b)}}{dr} \delta r \right) gV$$

We require adiabatic process for bubble

$$PV^\gamma = \text{const}$$

since m of the bubble is a const,
i.e. we have mass conservation

$$P \left(\frac{V}{m} \right)^\gamma = \text{const}_2$$

$$d(P \rho^{-\gamma}) = dP \rho^{-\gamma} + (-\gamma) P \rho^{-\gamma-1} d\rho = 0$$

$$\frac{d\rho}{dP} = \frac{\rho^{-\gamma}}{\gamma \rho^{-\gamma-1} P} = \frac{\rho_i^{(\ell)}}{\gamma P_i^{(\ell)}}$$

recall that
this is
about bubble

now we can plug it ^{start cond.} to the fnet eq.

$$f_{\text{net}} = Vg \left(\frac{d\rho^{(s)}}{dr} - \frac{\rho_i^{(\ell)}}{\gamma P_i^{(\ell)}} \frac{dP^{(\ell)}}{dr} \right) dr$$

since $P^{(\ell)} = P^{(s)}$ to maintain the bubble
and $\rho_i^{(\ell)} = \rho_i^{(s)}$ we can drop
' ℓ ' and ' s ' subscript, keeping in mind
that it is actually ' s ' everywhere now


$$f_{\text{net}} = Vg \left(\rho \frac{d\rho}{dr} - \frac{\rho}{\gamma P} \frac{dP}{dr} \right) dr = Vg \left(\frac{1}{\rho} \frac{d\rho}{dr} - \frac{1}{\gamma P} \frac{dP}{dr} \right) \rho g dr$$

~~ρ~~ some coef. \xrightarrow{A}

$$f_{\text{net}} = Vg A \rho dr$$

$$A = \frac{1}{\rho} \frac{d\rho}{dr} - \frac{1}{\gamma P} \frac{dP}{dr}$$

if $A < 0$ we have a restoring force

think about  So our bubble will go ~~there~~ back to equilibrium and thus oscillate around this point.

Since $\rho V = m$ of the bubble

$$f_{\text{net}} = ma = mg A \delta r$$

$$a = \ddot{(\delta r)}$$

get again equation of harmonic oscillator

$$\omega_{\text{mode}} = \sqrt{-gA} = \sqrt{\left(\frac{1}{\rho} \frac{dP}{dr} - \frac{1}{\rho} \frac{d\rho}{dr}\right) g}$$

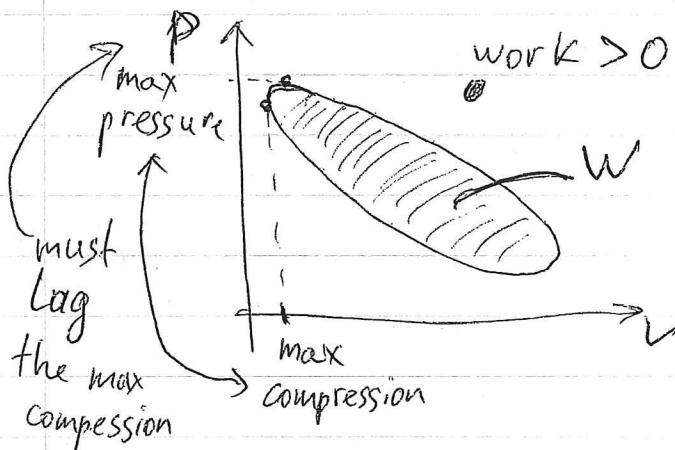
$$\pi = \frac{2\pi}{\omega}$$

Note that if $A > 0$ then there is no restoring force. Once the bubble starts to move up it will keep going to do it \Rightarrow ~~convection~~
convection condition

Oscillation Driving mechanism

It is not enough to have conditions for oscillations, ~~some~~ ^{source} mechanism must drive this oscillation since otherwise oscillation energy will dissipate and oscillation will cease.

On top of it whatever the source of energy (fusion) it must be applied right to do a positive work on a slice of a star.



if we go \curvearrowright
and negative if we
do \curvearrowleft direction

$$W = \oint P dV$$

Opacity effects, κ and γ - mechanism

Generally $\kappa \sim \frac{1}{T^{3.5}}$ \uparrow absorption
coef.

We need κ - to be increased with T
can be achieved in partially ionization
zone if ~~increase of~~ T injection of
heat goes not to kinetic energy of atoms
(and thus increase of T) but to
increased ionization and thus $\uparrow \kappa$

This will introduce delay:
Max pressure will happen after maximum
compression

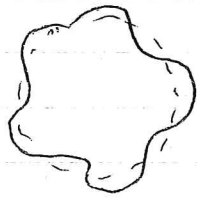


Additionally if some region is heated
less during compression (i.e. T is smaller
than ~~the~~ surrounding)
then heat additionally will go to such
region: γ - mechanism

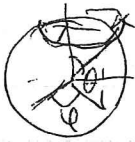
Partial ionization ^{zones} thus acts like a
pistons, and their location in star
defines what mode will oscillate if any

This in turn depends on star temperatures
see fig. 14-14

p-modes, going back to pressure governed mode. We discussed motion of the spherical layer as a whole, but we can also observe 'ripples' on the surface or in dipper part of the star



this described by spherical harmonic function $Y_l^m(\theta, \varphi)$

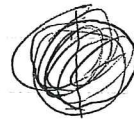


$$l = 0, 1, 2, 3 \dots$$

$$m = -l, \dots, 0, \dots, l \quad \text{overall } 2l+1 \text{ values}$$

$$Y_0^0(\theta, \varphi) = K_0^0 \text{ (const)}$$

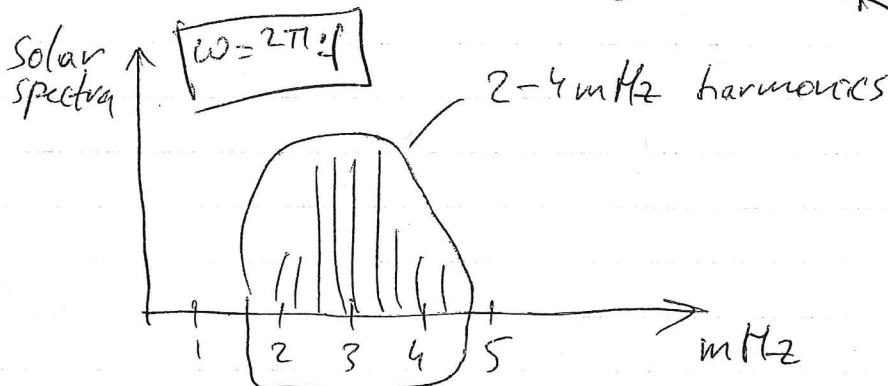
$$Y_1^0(\theta, \varphi) = K_1^0 \cos \theta$$



$$Y_l^{\pm 1}(\theta, \varphi) = K_l^{\pm 1} \sin \theta e^{i\varphi}$$

wavelength of oscillation $\lambda = \frac{2\pi R}{\sqrt{l(l+1)}}$

$$f_l = \frac{v_{\text{sound}}}{\lambda} = \sqrt{\frac{\gamma P}{\rho}} \frac{\sqrt{l(l+1)}}{2\pi R}$$



distance from center not the star radius

$$3 \text{ mHz} \approx 300 \text{ sec} = 5 \text{ min oscillation}$$