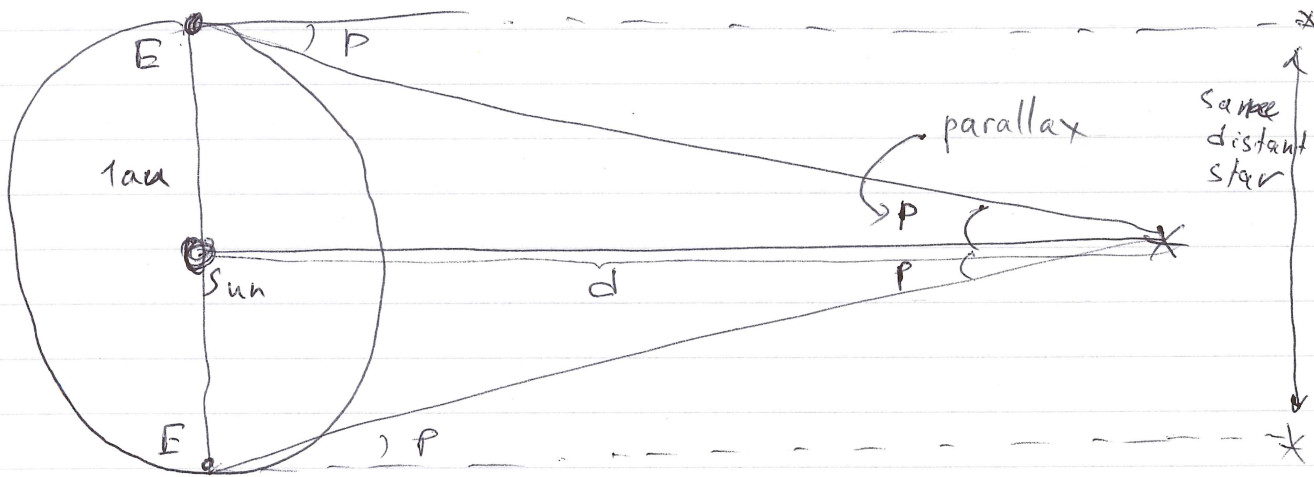


# Lecture 3 PO

Definition of par sec  
parallax second = pc

the largest base available to Earth based observer is the Diameter of Earth orbit.



$$1 \text{ au} \cdot d = \frac{1 \text{ au}}{\tan p} \quad / p \text{ is very small} /$$

Side note:  
Closest star:  
Proxima Centauri  
has distance to us = 1.3 pc

$$\Rightarrow d = \frac{1 \text{ au}}{p} \quad , \text{ since } p \text{ is small it is measured in arc second}$$

$$\Rightarrow d = \frac{1 \text{ au}}{p ["] \cdot \frac{1^\circ}{3600"} \cdot \frac{\pi}{180}} = \frac{1 \text{ au}}{1/(2.06 \cdot 10^5) ["]}$$

$$d (p = 1") = \boxed{\text{parsec} = \text{pc}} = \frac{1.5 \cdot 10^{11}}{1} \cdot 2.06 \cdot 10^5$$

Generally distance

$$d = \frac{1 \text{ pc}}{p \text{ parallax}}$$

$$\approx 3.08 \cdot 10^{16} \text{ m}$$

$$\approx 3.26 \text{ Lightyears} = 3.26 \text{ Ly}$$

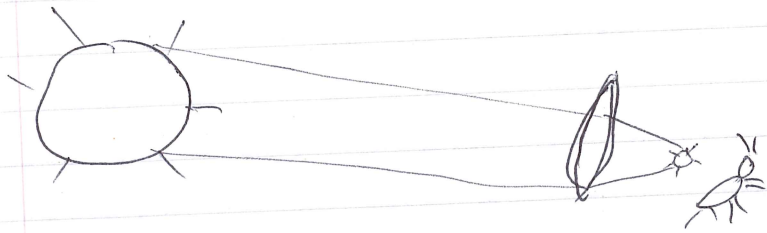
lecture 3 (pt)

Beyond naked eye observation.

Telescopes:

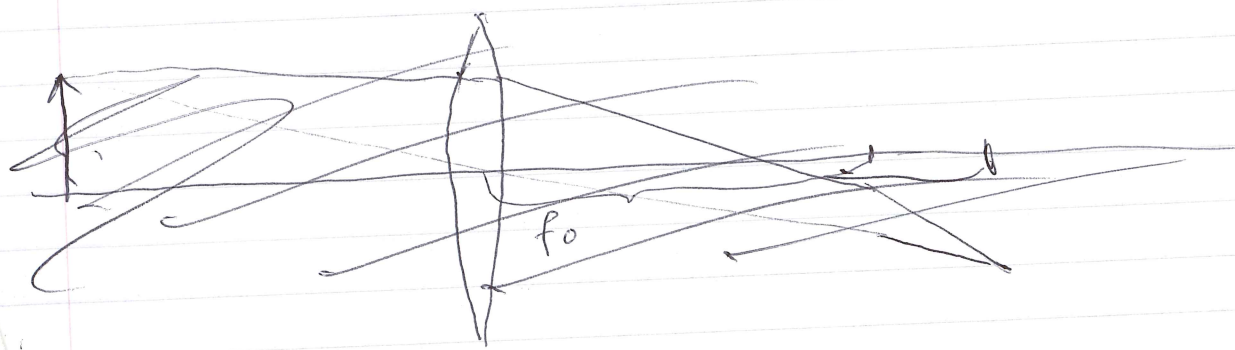
One lens is not enough though  
it often refers to as a magnifying  
glass.

Why? it makes very small image of  
a distant objects. Recall burning ants  
adventures

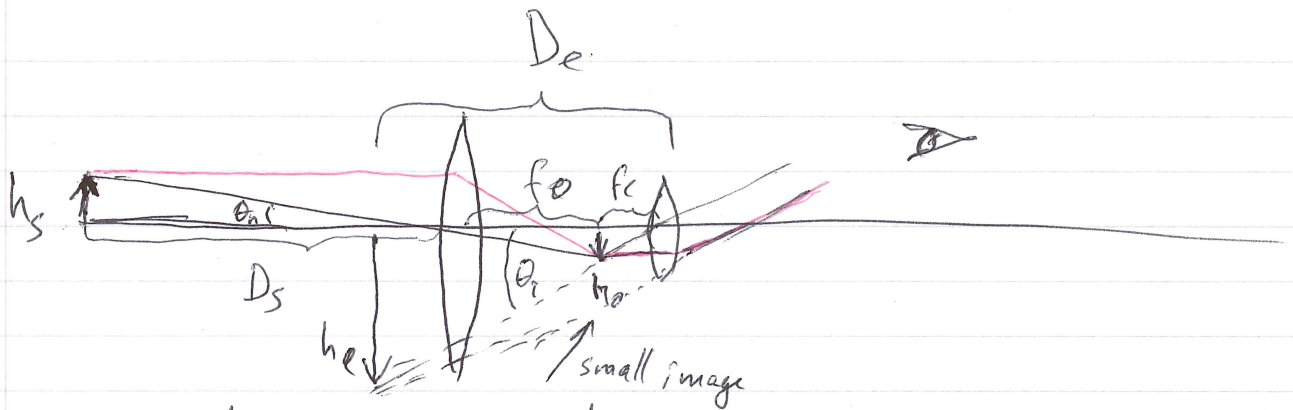


Side note.  
Negative lenses  
form "imaginary"  
image so we need  
2nd element

So we need two element system,  
one form image, 2nd magnifies



(p2)



$$\frac{h_o}{h_s} = \frac{f_o}{D_s} \ll 1$$

$$\frac{h_e}{h_o} = \frac{D_e}{f_e} \gg 1$$

but by itself  
it is not  
very useful

Note though that

$\theta_{\text{naked eye}} = \frac{h_s}{D_s}$  - was old (unaided) angular size of the star

$\theta_{\text{telescope}} = \frac{h_e}{D_e}$  - is new angular size

So angular magnification

$$M = \frac{\theta_T}{\theta_n} = \frac{h_s}{D_s} \cdot \frac{D_e}{h_e} = \frac{h_o}{f_o} \cdot \frac{f_e}{h_o} = \frac{f_e}{f_o}$$

$$M = \frac{f_e}{f_o}$$

So we can boost  
angular resolution  
with a telescope

## A few practical limitations

$M$  - can be large but we need to hold lens assembly together  
 $\downarrow$   
 telescope

$$\text{So } f_e + f_o \approx 10 \text{ m}$$

$f_e \approx \lambda \approx 1 \mu\text{m}$  ← but this is hard to operate

$$\text{So max } M \approx \frac{10 \text{ m}}{1 \mu\text{m}} = 10^7.$$

So theoretically we can do the smallest angular resolution of

$$\theta_{\text{detect}} / M = 0.5' / 10^7$$

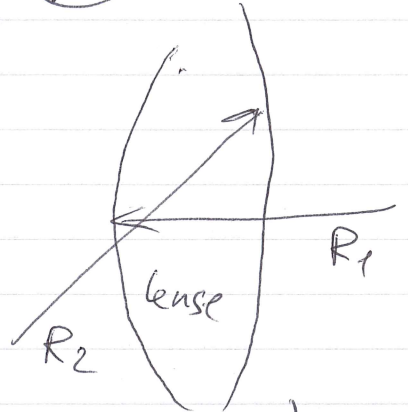
$\swarrow$   
 naked eye

sounds quite good but is it all to consider?

(p4)

## Chromatic aberrations

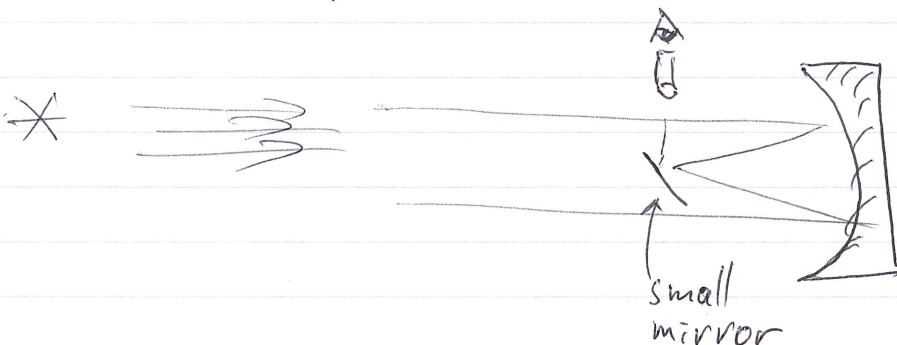
$$\frac{1}{f} = (n-1) \left( \frac{1}{R_1} - \frac{1}{R_2} \right)$$



Recall that  $n$  is a function of wavelength  $n = n(\lambda)$  otherwise we will see set of green, red, and blue images on top of each other.

Solutions: \* color filter  $\Rightarrow$  less light to detect

\* Reflector telescope suggested by Newton i.e. ~~re~~ replace lenses with mirror.



This technique was mastered by Herschel who built very large mirrors by literally polishing horse manure.

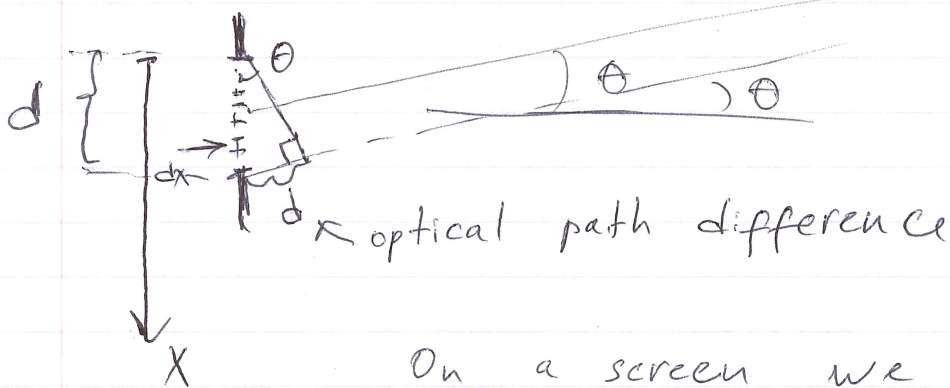
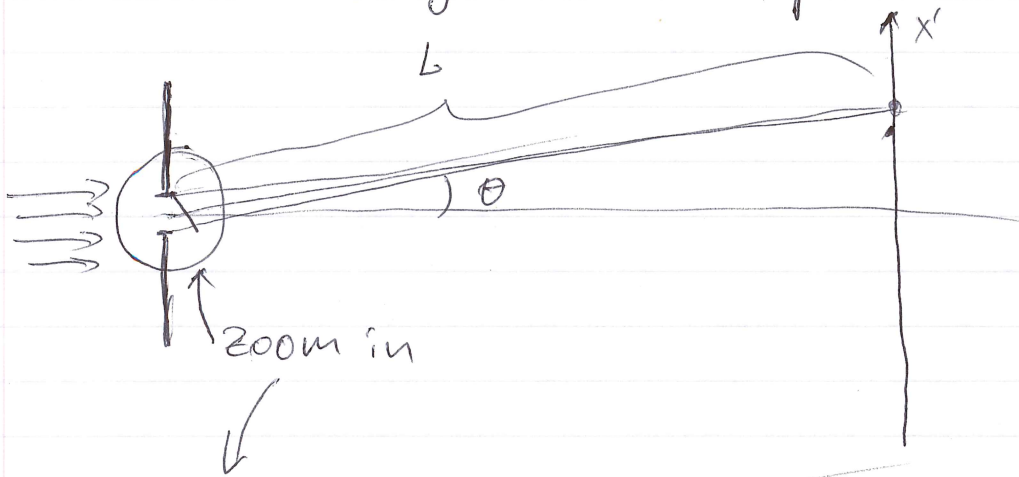
$\rightarrow$  1781 - find new planet Uranus

1803-1804 - binary star system (20 years of observations)

Same guy "map" of our galaxy as brightness over distance estimate.

A larger issue is diffraction!

Recall single slit experiment



On a screen we collect all light from the slit i.e. from every element  $dx$

Electromag (light) field will sum up on the screen

$$E_{\text{screen}} \propto \int_0^d e^{i \frac{2\pi(L + x \sin \theta)}{\lambda}} dx = \frac{E_0}{d} e^{i \frac{2\pi L}{\lambda}} \int_0^d e^{i \frac{2\pi x \sin \theta}{\lambda}} dx$$

proportionality

coef field per "pixel"

$\theta \ll 1$  but we will keep it

boring common factor

$$E_s = \frac{E'}{d} \int_0^d e^{i \frac{2\pi}{\lambda} \sin \theta x} dx =$$

$$= \frac{E'}{d} \frac{1}{i \frac{2\pi}{\lambda} \sin \theta} e^{i \frac{2\pi}{\lambda} \sin \theta x} \Big|_0^d =$$

$$= E' \frac{1}{i \frac{2\pi}{\lambda} \sin \theta} \cdot \left( e^{i \frac{2\pi}{\lambda} \sin \theta d} - 1 \right) \cdot \frac{e^{-\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}{e^{-\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}$$

$$= \frac{E' e^{-\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}{i \frac{2\pi}{\lambda} \sin \theta} \cdot \frac{e^{i \frac{\pi}{\lambda} \sin \theta d} - e^{-i \frac{\pi}{\lambda} \sin \theta d}}{2i} \cdot 2i$$

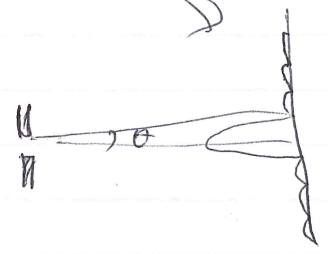
"  $\sin \left( \frac{\pi}{\lambda} \sin \theta d \right)$  "

$$= \frac{2E''}{2E} \cdot \frac{\sin \left( \frac{\pi}{\lambda} (\sin \theta) d \right)}{\frac{\pi}{\lambda} (\sin \theta) d}$$

const phase  $\varphi$

$$E_s = \frac{E \cdot e^{i\varphi}}{\frac{\pi}{\lambda} d (\sin \theta)} \sin \left( \frac{\pi}{\lambda} d \cdot (\sin \theta) \right)$$

Minima condition  $E_s = 0$

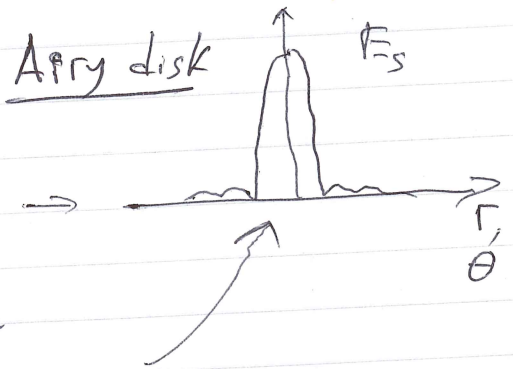
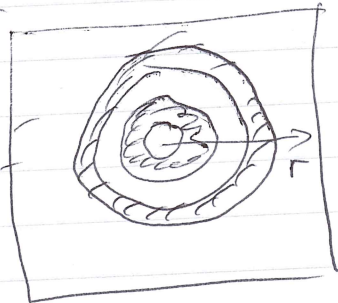
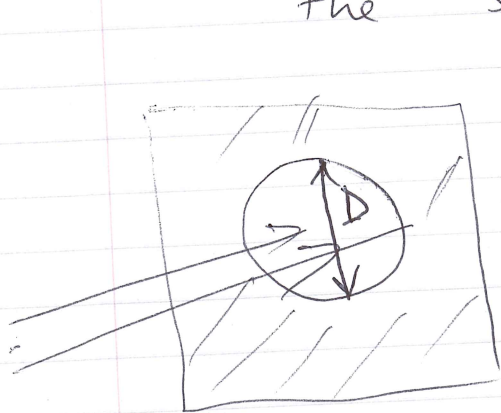


$$\frac{\pi}{\lambda} d \cdot \sin \theta = m \pi$$

$$\sin \theta = m \frac{\lambda}{d}$$

main maxima angle  $\theta_m = \frac{\lambda}{d}$

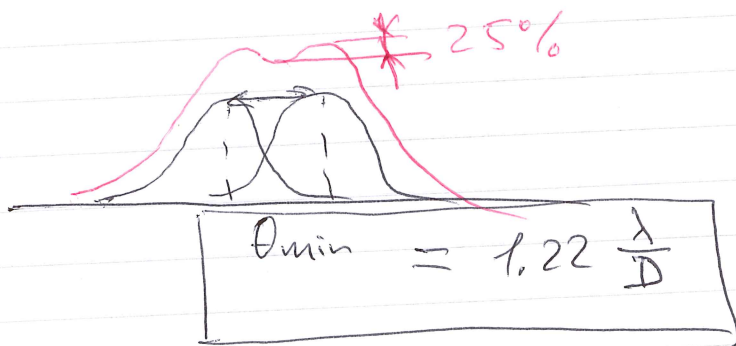
Well round apertures are almost the same



$$E_s(\theta) = \frac{2 J_1\left(\frac{\pi}{\lambda} D \sin \theta\right)}{\frac{\pi}{\lambda} D \sin \theta}$$

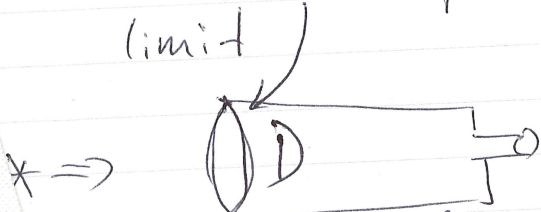
Bessel function

Now if we have two objects close to each other, they form two overlapping Airy disks, which in cross section looks like.



Rayleigh criterion

So size of the objective is the actual limit



$D \sim 10 \text{ m}$  ~~so it is nearly matches~~

$\lambda \sim 1 \mu\text{m}$

$$\theta_m = 1.2 \frac{\lambda}{D} = 1.2 \cdot \frac{10^{-6} \text{ m}}{10 \text{ m}} \approx 10^{-7} \text{ rad}$$

~~physical building~~  
~~ie. construction limit on~~



Recall that aided eye limited by  
~~to~~ maximum  $M$  - achievable by construction

$$\theta_m = \theta_{\text{eye}}/M = \frac{2\text{mm}}{2\text{cm}} / \frac{10\text{m}}{1\text{m}} =$$

$$= 10^{-4} / 10^7 = 10^{-11} \text{ rad which}$$

is not achievable due to diffraction

$$\text{with } \theta_m \approx \frac{\lambda}{D} = \frac{10^{-6}\text{m}}{10\text{m}} \approx 10^{-7} \text{ rad}$$

$$1 \text{ rad} = 2 \cdot 10^5 \text{ arc sec}$$