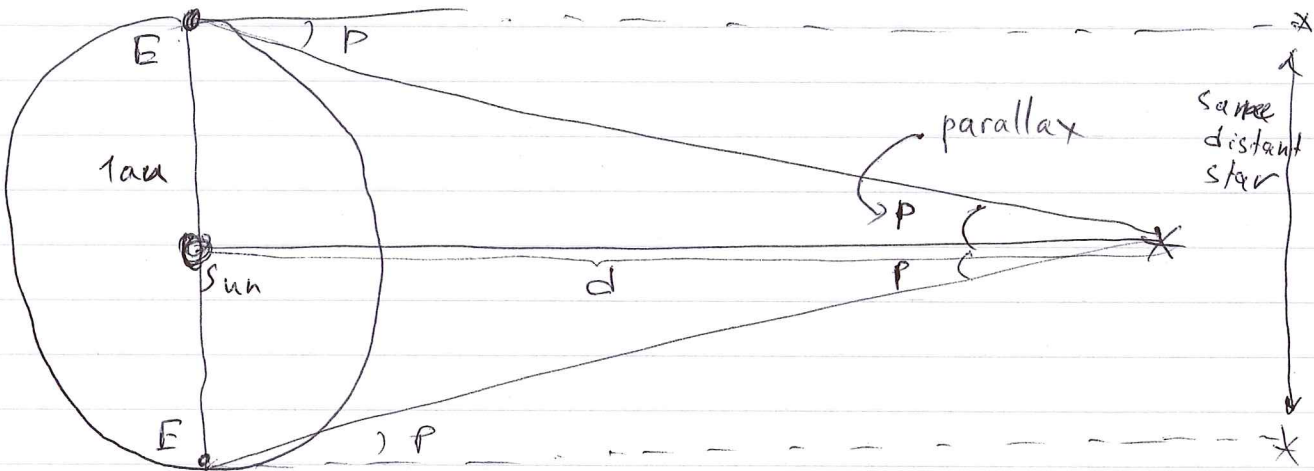


Lecture 3

(PO)

Definition of par sec
parallax second = pc

the largest base available to Earth based observer is the Diameter of Earth orbit.



$$1 \text{ au} \cdot d = \frac{1 \text{ au}}{\tan p} \approx \frac{1 \text{ au}}{p} \text{ (since } p \text{ is very small)}$$

Side note:
Closest star:
Proxima Centauri
has distance to us = 1.3 pc

$$\Rightarrow d = \frac{1 \text{ au}}{p}, \text{ since } p \text{ is small it is measured in arc second}$$

$$d = \frac{1 \text{ au}}{p ["] \cdot \frac{1^\circ}{3600"} \cdot \frac{\pi}{180}} = \frac{1 \text{ au}}{1/(2.06 \cdot 10^5) ["]}$$

$$d (p = 1") = \text{parsec} = \text{pc} = \frac{1.5 \cdot 10^{11} \cdot 2.06 \cdot 10^5}{1}$$

Generally distance
 $d = \frac{1 \text{ pc}}{p \text{ parallax}}$

$$\approx 3.08 \cdot 10^{16} \text{ m}$$

$$\approx 3.26 \text{ Lightyears} = 3.26 \text{ Ly}$$

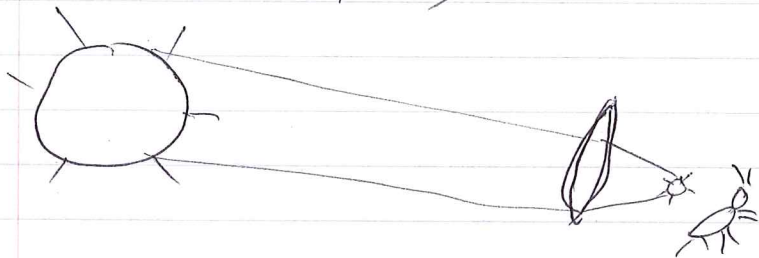
lecture 3 (P1)

Beyond naked eye observation.

Telescopes:

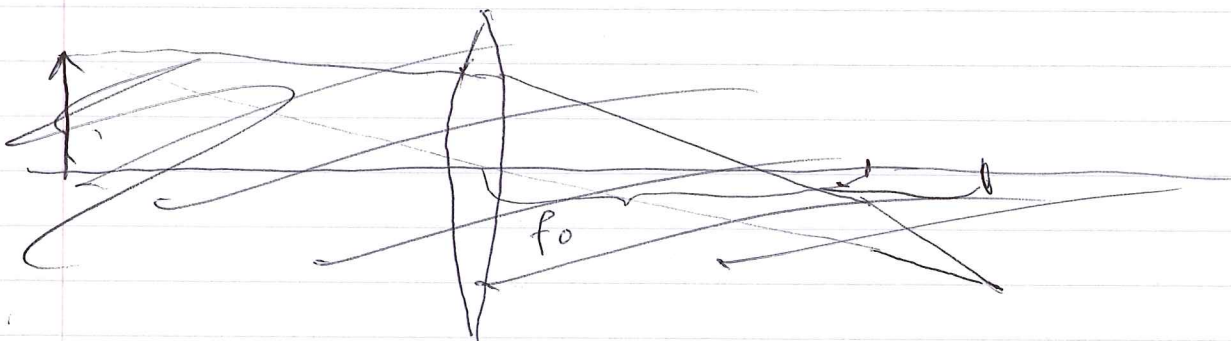
One lens is not enough though
it often refers to as a magnifying
glass.

Why it make very small image of
a distant objects. Recall burning ants
adventures

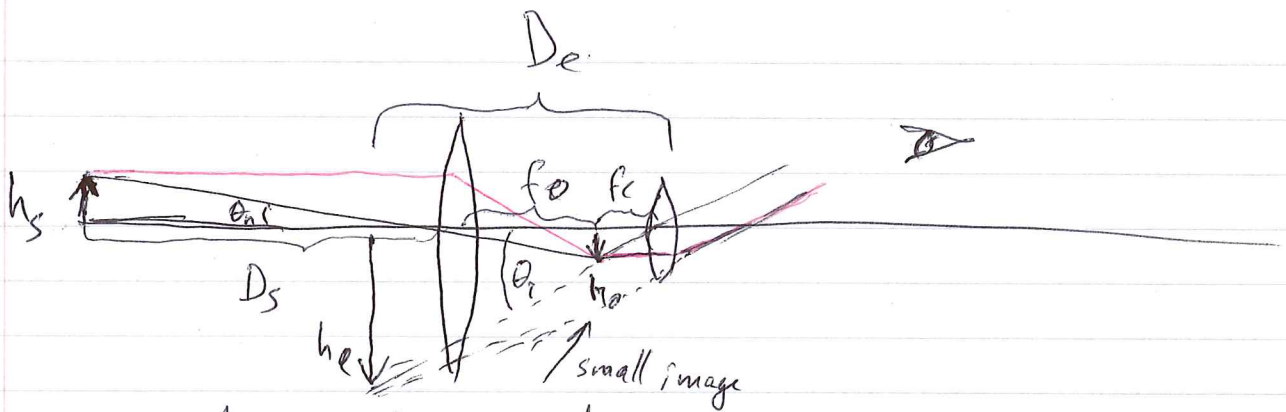


Side not.
Negative lenses
form "imaginary"
image so we need
2nd element

So we need two element system,
one form image, 2nd magnifies



(p2)



$$\frac{h_o}{h_s} = \frac{f_o}{D_s} \ll 1$$

$$\frac{h_e}{h_o} = \frac{D_e}{f_e} \gg 1$$

but by itself
it is not
very useful

Note though that

$\theta_{\text{naked eye}} = \frac{h_s}{D_s}$ - was old (unaided) angular size of the star

$\theta_{\text{telescope}} = \frac{h_e}{D_e}$ - is new angular size

So angular magnification

$$M = \frac{\theta_T}{\theta_n} = \frac{h_s}{D_s} \cdot \frac{D_e}{h_e} = \frac{h_o}{f_o} \cdot \frac{f_e}{h_o} = \frac{f_e}{f_o}$$

$$M = \frac{f_e}{f_o}$$

So we can boost
angular resolution
with a telescope

A few practical limitations

M - can be large but we need to hold lens assembly together



telescope

$$\text{So } f_e + f_o \approx 10 \text{ m}$$

$f_e \approx \lambda \approx 1 \mu\text{m}$ ← but this is hard to operate

$$\text{So max } M \approx \frac{10 \text{ m}}{1 \mu\text{m}} = 10^7.$$

So theoretically we can do the smallest angular resolution of

$$\theta_{\text{detect}} / M = 0.5' / 10^7$$

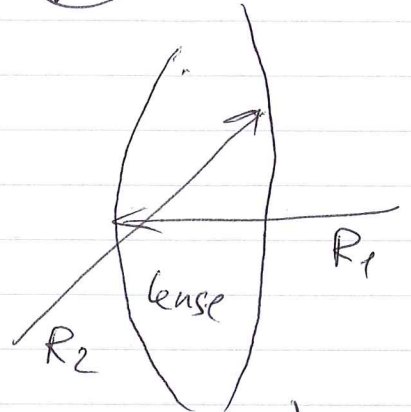
→ naked eye

sounds quite good but is it all to consider?

(p4)

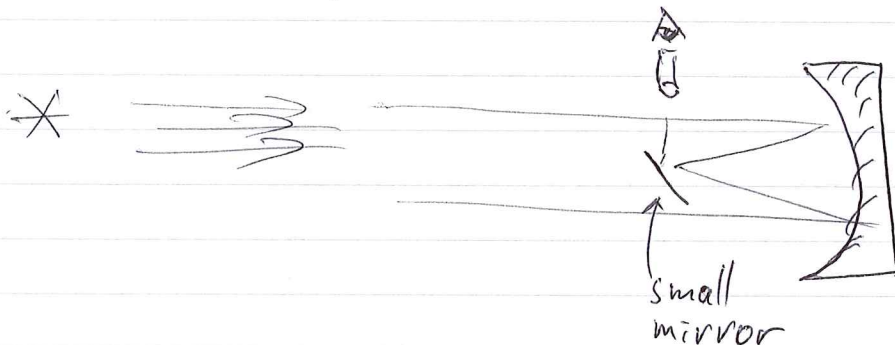
Chromatic aberrations

$$\frac{1}{f} = (n-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$



Recall that n is a function of wavelength $n = n(\lambda)$ otherwise we will see set of green, red, and blue images on top of each other.

- Solutions: *
- * color filter \Rightarrow less light to detect
 - * Reflector telescope suggested by Newton i.e. ~~replace~~ replace lenses with mirror.



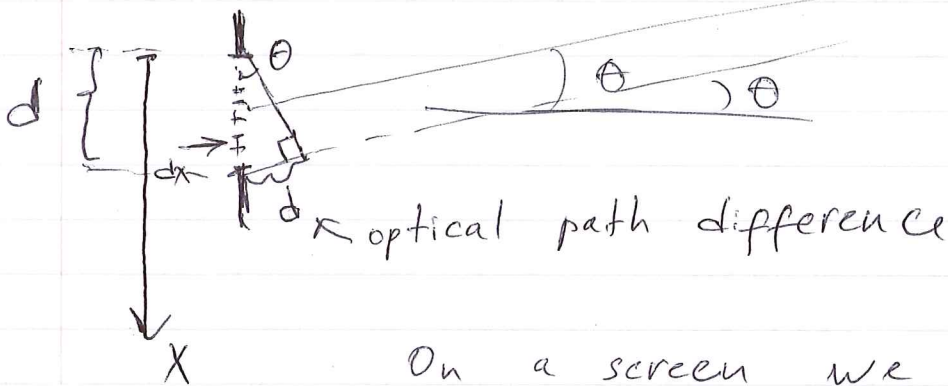
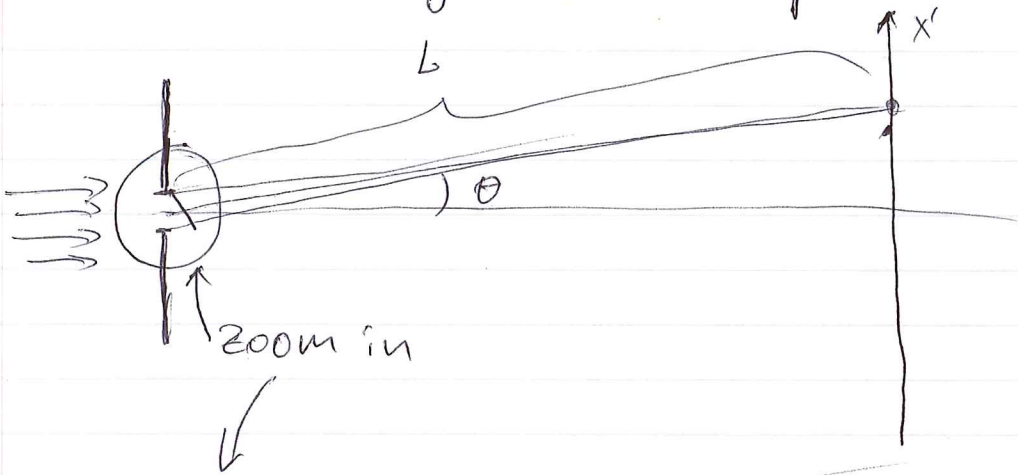
This technique was mastered by Herschel who built very large mirrors by literally polishing horse manure.

\rightarrow 1781 - find new planet Uranus

1803-1804 - binary star system (20 years of observations)
Same guy "map" of our galaxy as brightness over distance estimate.

A larger issue is diffraction!

Recall single slit experiment



On a screen we collect all light from the slit i.e. from every element dx

Electromag (light) field will sum up on the screen

$$E_{\text{screen}} \propto \int_0^d e^{i \frac{2\pi}{\lambda} (L + x \sin \theta)} dx = \frac{E'}{d} \int_0^d e^{i \frac{2\pi}{\lambda} x \sin \theta} dx$$

proportionality
coef field per "pixel"
= / $\theta \ll 1$ but we will keep it / = $\frac{E'}{d} \int_0^d e^{i \frac{2\pi}{\lambda} x \sin \theta} dx$

boring common factor

$$E_s = \frac{E'}{d} \int_0^d e^{i \frac{2\pi}{\lambda} \sin \theta x} dx =$$

$$= \frac{E'}{d} \frac{1}{i \frac{2\pi}{\lambda} \sin \theta} e^{i \frac{2\pi}{\lambda} \sin \theta x} \Big|_0^d =$$

$$= E' \frac{1}{i \frac{2\pi}{\lambda} \sin \theta} \cdot \left(e^{i \frac{2\pi}{\lambda} \sin \theta d} - 1 \right) \cdot \frac{e^{-\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}{e^{\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}$$

$$= \frac{E' e^{-\frac{1}{2} i \frac{2\pi}{\lambda} \sin \theta d}}{i \frac{2\pi}{\lambda} \sin \theta} \text{ "another boring factor"}$$

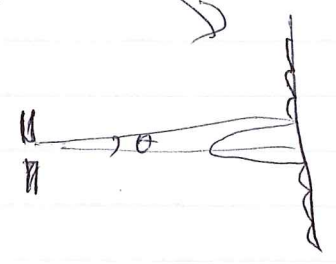
$$\frac{e^{i \frac{\pi}{\lambda} \sin \theta d} - e^{-i \frac{\pi}{\lambda} \sin \theta d}}{2i} \cdot 2i = \sin \left(\frac{\pi}{\lambda} \sin \theta d \right)$$

$$= \frac{2E' e^{i\varphi}}{2E} \cdot \frac{\sin \left(\frac{\pi}{\lambda} (\sin \theta) d \right)}{\frac{\pi}{\lambda} (\sin \theta) d}$$

const phase φ

$$E_s = \frac{E \cdot e^{i\varphi}}{\frac{\pi}{\lambda} d (\sin \theta)} \sin \left(\frac{\pi}{\lambda} d \cdot (\sin \theta) \right)$$

Minima condition $E_s = 0$

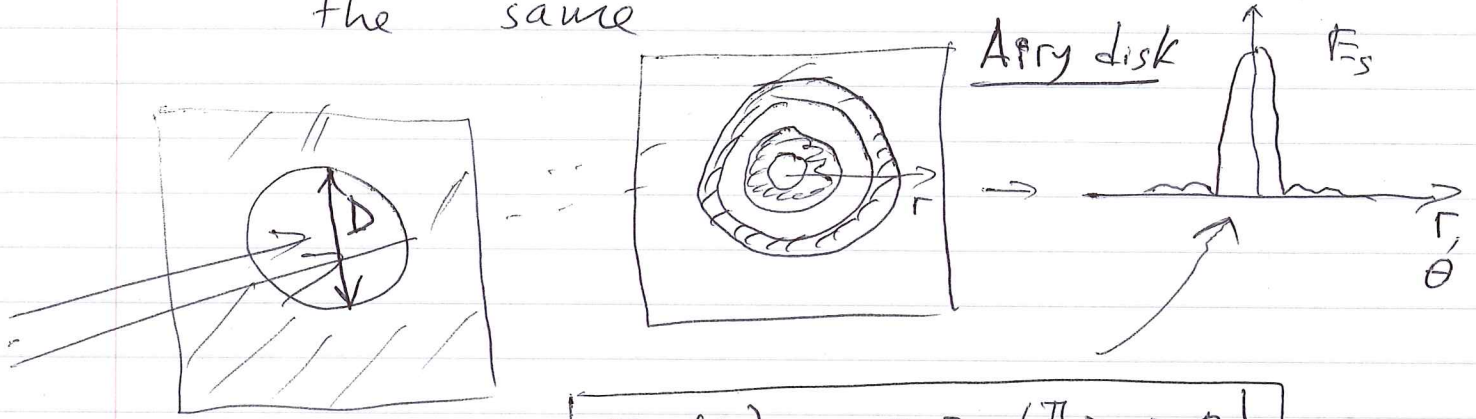


$$\frac{\pi}{\lambda} d \cdot \sin \theta = m \pi$$

$$\sin \theta = m \frac{\lambda}{d}$$

main maxima angle $\theta_m = \frac{\lambda}{d}$

Well round apertures are almost the same



$$E_s(\theta) = \frac{2 J_1\left(\frac{\pi}{\lambda} D \sin \theta\right)}{\frac{\pi}{\lambda} D \sin \theta}$$

Bessel function

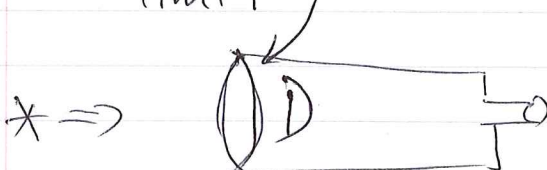
Now if we have two objects close to each other, they form two overlapping Airy disks, which in cross section looks like.



$$\theta_{min} = 1.22 \frac{\lambda}{D}$$

Rayleigh criterion

So size of the objective is the actual limit



$D \sim 10\text{m}$
 $\lambda \sim 1\mu\text{m}$

~~so it is nearly matches~~

$$\theta_m = 1.2 \frac{\lambda}{D} = 1.2 \cdot \frac{10^{-6}\text{m}}{10\text{m}} \approx 10^{-7}\text{rad}$$

~~Physical building construction limit $\theta_m \approx 10^{-5}\text{rad}$~~

Recall that aided eye limited by maximum M - achivable by construction

$$\theta_m = \theta_{eye}/M = \frac{2\text{mm}}{2\text{cm}} / \frac{10\text{m}}{1\text{m}} = 10^{-9} / 10^7 = 10^{-16} \text{ rad which}$$

is not achivable due to diffraction

with $\theta_m \approx \frac{\lambda}{D} = \frac{10^{-6}\text{m}}{10\text{m}} \approx 10^{-7} \text{ rad}$

$$1 \text{ rad} = 2 \cdot 10^5 \text{ arc sec}$$