

Relativistic Kinematics

Four-Vector Momentum: Generalization of $\vec{p} = m \frac{d\vec{x}}{dt}$

Let x^μ be the spacetime coordinates in some inertial frame.
parametrize the particle trajectory by $\vec{x}(t)$ or $\vec{x}(\tau)$

τ = proper time elapsed
along particle trajectory.

Recall $dt = \gamma d\tau$ ← "Moving clocks run slow"

Define $p^\mu = m \frac{dx^\mu}{d\tau}$ ← transforms like a 4-vector.
 $d\tau$ is Lorentz invariant.

$$p^0 = m \frac{dx^0}{d\tau} = mc\gamma$$

$$p^i = m \frac{dx^i}{d\tau} = m \vec{v} \gamma \quad \leftarrow v^i = \frac{dx^i}{dt}$$

In the nonrelativistic limit, $p^0 \approx mc \left(1 + \frac{1}{2} \frac{\vec{v}^2}{c^2} \right) = mc + \frac{1}{2} m \frac{\vec{v}^2}{c}$

mc is a constant independent of \vec{v} .

$\frac{1}{2} m \vec{v}^2$ is the nonrelativistic kinetic energy.

⇒ Identify $p^0 = E/c$, where $E =$ relativistic energy

we have $\boxed{E = mc^2 \gamma} = \underbrace{mc^2}_{\text{rest energy}} + \underbrace{mc^2(\gamma - 1)}_{\text{kinetic energy}}$

Spatial components of 4-momentum: $\vec{p} = m \frac{d\vec{x}}{dt} = m \frac{d\vec{x}}{dt} \gamma$

Nonrelativistic limit: $\vec{p} \approx m \frac{d\vec{x}}{dt} = m \vec{v}$. = nonrelativistic momentum.

Just as $ds^2 = dx^\mu dx^\nu \eta_{\mu\nu}$ is Lorentz invariant,

$$\begin{aligned} \text{so is } p^\mu p^\nu \eta_{\mu\nu} &= (p^0)^2 - \vec{p}^2 \\ &= (m\gamma c)^2 - (m\vec{v}\gamma)^2 \\ &= \frac{m^2 c^2}{1-v^2/c^2} (1-v^2/c^2) \\ &= m^2 c^2 \end{aligned}$$

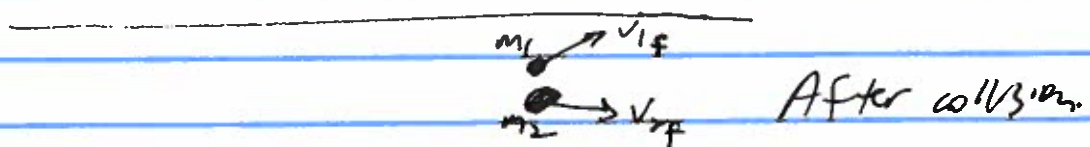
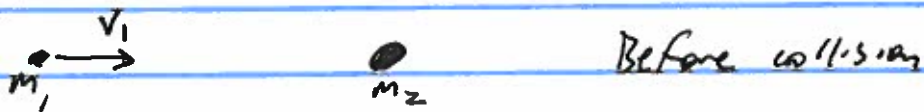
In other words, $E^2/c^2 = \vec{p}^2 + m^2 c^2$,

or
$$\boxed{E^2 = \vec{p}^2 c^2 + m^2 c^4}$$
 ← Relativistic relation between Energy and momentum.

Spacetime translation invariance → Conservation of (total) E and \vec{p} .

Kinematics of Collisions

Lab frame: one particle is at rest initially.



Initial total 4-momentum:

Suppose velocity of m_1 is in x -direction.

$$p^\mu = (m_1 c \gamma_1 + m_2 c, m_1 \gamma_1 v_1, 0, 0)$$

In the Center-of-Momentum frame, the total $\vec{p} = \vec{0}$.
(But the energy $p^0/c \neq 0$.)



$$\vec{p}' = m_1 \vec{v}'_1 \gamma'_1 + m_2 \vec{v}'_2 \gamma'_2 = \vec{0}$$

$$p'^\mu = (m_1 \gamma'_1 c + m_2 \gamma'_2 c, 0, 0, 0)$$

The boost $\vec{\beta}'$ to go from the lab to the COM frame:

$$\vec{\beta}' = -\vec{\beta}_2' \quad (\text{because } \vec{\beta}_2 = 0 \text{ — initial velocity of } m_2 \text{ in lab frame})$$

Velocity $\vec{\beta}'_1 c$ of mass m_1 in COM frame in terms of $\vec{\beta}'_1$ in lab frame and boost $\vec{\beta}'$ — velocity addition.

$$\beta'_1 = \frac{\beta_1 - \beta'}{1 - \beta_1 \beta'}$$

To find the boost velocity β' relating the two frames, we can use the Lorentz invariance of $p^\mu p_\mu$:

In the lab frame,

$$\begin{aligned} (p^0)^2 - \vec{p}^2 &= (m_1 c \gamma_1 + m_2 c)^2 - (m_1 v_1 \gamma_1)^2 = (m_1^2 + m_2^2 + 2m_1 m_2 \gamma_1) c^2 \\ &= (m_1^2 + m_2^2) c^2 + 2m_2 E_1 \\ &\quad E_1 = m_1 c^2 \gamma_1 \end{aligned}$$

In the COM frame:

$$(p'^0)^2 - (\vec{p}')^2 = (m_1 \gamma_1' c + m_2 \gamma_2' c)^2$$

$$= \left\{ -\frac{m_2 \gamma_2' \beta_2' c}{\beta_1'} + m_2 \gamma_2' c \right\}^2 \quad \text{using } \gamma_1' \beta_1' = -m_2 \gamma_2' \beta_2'$$

$$= m_2^2 (\gamma_2')^2 c^2 \left\{ 1 + \frac{\beta_1' (1 - \beta_1' \beta_2')}{\beta_1 - \beta_1'} \right\}^2 \quad \text{using } \beta_1' = \frac{\beta_1 - \beta_2'}{1 - \beta_1 \beta_2'}$$

and $\beta_2' = -\beta_1'$

$$= m_2^2 (\gamma_2')^2 c^2 \left\{ \frac{\beta_1' (1 - (\beta_1')^2)}{\beta_1 - \beta_1'} \right\}^2$$

$$= \boxed{\frac{m_2^2 c^2 \beta_1'^2 (1 - (\beta_1')^2)}{(\beta_1 - \beta_1')^2}} \quad \text{using } (\gamma_2')^2 = \frac{1}{1 - (\beta_1')^2}$$

This must be the same as $(p^0)^2 - \vec{p}^2$ in the lab frame.

Hence, the boost from the lab to COM frame is determined by

$$\boxed{\frac{m_2^2 c^2 \beta_1'^2 (1 - (\beta_1')^2)}{(\beta_1 - \beta_1')^2} = \left[m_1^2 + m_2^2 + \frac{2m_1 m_2}{\sqrt{1 - \beta_1^2}} \right] c^2}$$

Particle Production: Consider the nuclear reaction $p + n \rightarrow p + \bar{p} + n$

$p = \text{proton}$, $\bar{p} = \text{antiproton}$, $n = \text{neutron}$.

$m_p = 938 \text{ MeV}$, similar to m_n .

- treat all masses as equal.

Question: If the initial neutron is at rest in the lab frame, what is the minimum required kinetic energy of the proton in order to create the $p\bar{p}$ pair?

Insight: In lab frame calculate $p^\mu p_\mu$ γ_{rel} :
 $(p^0)^2 - (\vec{p})^2 = (m_p^2 + m_n^2) c^2 + 2 m_n \underbrace{m_p c^2 \gamma_p}_{E_p - \text{proton energy}}$

Write $E_p = m_p c^2 + T_p$
 \leftarrow proton kinetic energy

$$\Rightarrow (p^0)^2 - (\vec{p})^2 = (m_p + m_n)^2 c^2 + 2 m_n T_p$$

At threshold for the reaction (i.e. minimum T_p), in the COM frame all $\overset{\text{final}}{n}$ particles are at rest.

$$\Rightarrow (p^{0'})^2 - (\vec{p}')^2 = (3m_p + m_n)^2 c^2 \geq (p^0)^2 - (\vec{p})^2$$

by conservation of p^μ and Lorentz invariance of $(p^0)^2 - (\vec{p})^2$.

$$\Rightarrow T_p = \frac{[(3m_p + m_n)^2 - (m_p + m_n)^2] c^2}{2 m_n} = \frac{[8m_p^2 + 4m_p m_n] c^2}{2 m_n}$$

$$T_p \approx 6 m_p c^2 = 5.63 \text{ GeV}$$

\rightarrow This is 3x the difference in final and initial rest energies!