

Homework 3

Physics 721-QFT
handed out 22 Sep. 2009

1. (a) For a complex scalar field

$$\phi(x) = \int \frac{d^3p}{(2\pi)^3 2\omega_p} \left(a_{\vec{p}} e^{-ipx} + b_{\vec{p}}^\dagger e^{ipx} \right)$$

find the charge operator that goes along with the current

$$j^\mu = : i\phi^\dagger \overleftrightarrow{\partial}^\mu \phi : = : i(\phi^\dagger \partial^\mu \phi - \partial^\mu \phi^\dagger \phi) :$$

in terms of the creation and annihilation operators ($a_{\vec{p}}^\dagger, b_{\vec{p}}^\dagger$ and $a_{\vec{p}}, b_{\vec{p}}$, respectively).

(b) Show that the two categories of single particle states, $|\vec{p}\rangle_a = a_{\vec{p}}^\dagger |0\rangle$ and $|\vec{p}\rangle_b = b_{\vec{p}}^\dagger |0\rangle$, are both eigenstates of the charge operator, with opposite eigenvalues.

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a) Can choose $t = 0$

$$Q = \int d^3x j^0(x) = : i \left(\int d^3x \phi^\dagger(x) \dot{\phi}(x) - h.c. \right) : \quad (1)$$

$$= i : \int (d^3p) \int (d^3p') \int d^3x \left(a_{\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}} + b_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} \right) (-i\omega_p) \left(a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} - b_{\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}} \right) - h.c. : \quad (2)$$

$$= \frac{1}{2} \left(\int (d^3p) \left[a_{\vec{p}}^\dagger a_{\vec{p}} - a_{-\vec{p}}^\dagger b_{\vec{p}}^\dagger + b_{-\vec{p}} a_{\vec{p}} - b_{\vec{p}}^\dagger b_{\vec{p}} \right] + h.c. \right) \quad (3)$$

$$\therefore Q = \int \frac{d^3p}{(2\pi)^3 2\omega_p} \left(a_{\vec{p}}^\dagger a_{\vec{p}} - b_{\vec{p}}^\dagger b_{\vec{p}} \right) \quad (4)$$

b)

$$Q|\vec{p}\rangle_a = \int \frac{d^3p'}{(2\pi)^3 2\omega_{p'}} a_{\vec{p}'}^\dagger a_{\vec{p}'} a_{\vec{p}}^\dagger |0\rangle = \int \frac{d^3p'}{(2\pi)^3 2\omega_{p'}} a_{\vec{p}'}^\dagger \left[(2\pi)^3 2\omega_p \delta^3(p' - p) + a_{\vec{p}}^\dagger a_{\vec{p}'} \right] |0\rangle = a_{\vec{p}}^\dagger |0\rangle \quad (5)$$

$$\therefore Q|\vec{p}\rangle_a = |\vec{p}\rangle_a \quad (6)$$

$Q|\vec{p}\rangle_b$ is the same except for the minus sign.

2. Show that the rapidity η , defined from the Lorentz parameter γ by $\cosh \eta = \gamma$, is additive for any two successive Lorentz transformations in the same direction.

(I.e., If we label a Lorentz transformation by the value of the rapidity, show that $\Lambda(\eta_2) \circ \Lambda(\eta_1) = \Lambda(\eta_1 + \eta_2)$.)

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 Choose axes so that the x -axis lines up with the direction of the Lorentz transformations.

Then

$$\Lambda(\eta) = \begin{pmatrix} \cosh \eta & \sinh \eta & 0 & 0 \\ \sinh \eta & \cosh \eta & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (7)$$

and

$$\Lambda(\eta_2)\Lambda(\eta_1) = \begin{pmatrix} \cosh \eta_2 & \sinh \eta_2 & 0 & 0 \\ \sinh \eta_2 & \cosh \eta_2 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cosh \eta_1 & \sinh \eta_1 & 0 & 0 \\ \sinh \eta_1 & \cosh \eta_1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (8)$$

$$= \begin{pmatrix} \cosh \eta_2 \cosh \eta_1 + \sinh \eta_2 \sinh \eta_1 & \cosh \eta_2 \sinh \eta_1 + \sinh \eta_2 \cosh \eta_1 & 0 & 0 \\ \cosh \eta_2 \sinh \eta_1 + \sinh \eta_2 \cosh \eta_1 & \cosh \eta_2 \cosh \eta_1 + \sinh \eta_2 \sinh \eta_1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (9)$$

$$= \begin{pmatrix} \cosh(\eta_1 + \eta_2) & \sinh(\eta_1 + \eta_2) & 0 & 0 \\ \sinh(\eta_1 + \eta_2) & \cosh(\eta_1 + \eta_2) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \quad (10)$$

$$= \Lambda(\eta_1 + \eta_2) \quad (11)$$