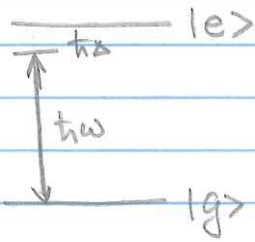


Near-resonant two-level system (Rabi model)



$$h\omega_0 = E_e - E_g$$

$$\Delta = \omega_0 - \omega$$

$$|\psi(t)\rangle = c_g(t) e^{-iE_g t/\hbar} |g\rangle + c_e(t) e^{-iE_e t/\hbar} |e\rangle$$

[recall $\dot{c}_e(t) = -\frac{i}{\hbar} \sum_k c_k(t) \langle e | \hat{H}_I | k \rangle e^{i\omega_{ek} t}$]

$$\langle g | \hat{H}_I | g \rangle = \langle e | \hat{H}_I | e \rangle = 0 \quad \langle e | \hat{H}_I | g \rangle = \rho_{eg} E_0 \cos \omega t$$

$$\dot{c}_g(t) = -\frac{i}{\hbar} c_e \rho_{eg}^* E_0 \cos \omega t e^{-i\omega_0 t}$$

$$\dot{c}_e(t) = -\frac{i}{\hbar} c_g \rho_{eg} E_0 \cos \omega t e^{i\omega_0 t}$$

Using RWA (keeping only terms $\sim \omega - \omega_0$)

$$\dot{c}_g(t) = -\frac{i}{2\hbar} \rho_{eg}^* E_0 e^{-i\Delta t} c_e(t)$$

$$\dot{c}_e(t) = -\frac{i}{2\hbar} \rho_{eg} E_0 e^{i\Delta t} c_g(t)$$

$$\frac{d}{dt} (\dot{c}_e e^{-i\Delta t}) = -\frac{i}{2\hbar} \rho_{eg} E_0 \dot{c}_g(t) = -\frac{|\rho_{eg} E_0|^2}{4\hbar^2} e^{-i\Delta t} c_e(t)$$

$$\ddot{c}_e e^{-i\Delta t} - i\Delta \dot{c}_e e^{-i\Delta t}$$

$$\ddot{c}_e - i\Delta \dot{c}_e + \frac{|\rho_{eg} E_0|^2}{4\hbar^2} c_e = 0 \quad \text{oscillator eqn}$$

The result: In using the RWA, the difference between the driving frequency and the natural frequency is more prominent!

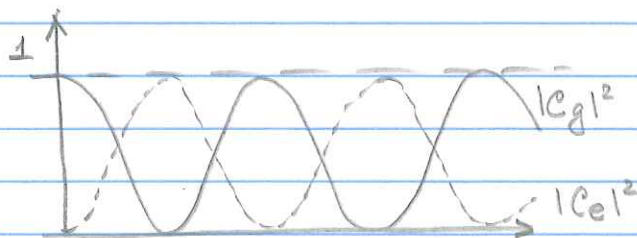
Rabi frequency $\Omega_R = \frac{\mu_{eg} E_0}{\hbar}$

$\Delta = 0$

$$\begin{cases} \ddot{c}_e + \frac{1}{4}\Omega_R^2 c_e = 0 \\ \ddot{c}_g + \frac{1}{4}\Omega_R^2 c_g = 0 \end{cases} \text{ oscillating}$$

assuming $c_g(t=0) = 1$, $c_e(t=0) = 0$

$$c_g(t) = \cos\left(\frac{1}{2}\Omega_R t\right) \quad c_e(t) = i \sin\left(\frac{1}{2}\Omega_R t\right)$$



$$\begin{aligned} P_e - P_g &= \cos^2\left(\frac{\Omega_R t}{2}\right) - \sin^2\left(\frac{\Omega_R t}{2}\right) \\ &= \cos \Omega_R t \end{aligned}$$

For $\Delta \neq 0$

$$c_e(t) = i \frac{\Omega_R}{\sqrt{\Delta^2 + \Omega_R^2}} e^{i\Delta t/2} \sin\left(\frac{1}{2}\sqrt{\Omega_R^2 + \Delta^2} \cdot t\right)$$

$$c_g(t) = e^{i\Delta t/2} \left\{ \cos\left(\frac{1}{2}\sqrt{\Omega_R^2 + \Delta^2} \cdot t\right) - \frac{i\Delta}{\sqrt{\Delta^2 + \Omega_R^2}} \sin\left(\frac{1}{2}\sqrt{\Omega_R^2 + \Delta^2} \cdot t\right) \right\}$$

Often $\tilde{\Omega}_R = \sqrt{\Omega_R^2 + \Delta^2}$ is called generalized Rabi frequency.

Notice that Rabi oscillations are fully coherent, since at any moment of time we can define atomic quantum state, i.e., for $\Delta = 0$

$$\psi(t) = \cos\left(\frac{1}{2}\Omega_R t\right) |g\rangle + i \sin\left(\frac{1}{2}\Omega_R t\right) |e\rangle$$

$$P_g = \cos^2\left(\frac{\Omega_R t}{2}\right) \quad P_e = \sin^2\left(\frac{\Omega_R t}{2}\right)$$

time-dependent population
(electron oscillates (flips) p/w these two states)