

Quantum Rabi Flopping

$$|\psi(t)\rangle = c_a(t) |a, n+1\rangle + c_b(t) |b, n\rangle$$

$$\text{let } c_b(t=0) = 1 \quad c_a(t=0) = 1$$

$$i\hbar \frac{\partial \psi}{\partial t} = \hat{H} \psi \quad \text{gives common phase factor}$$

$$i\hbar \begin{pmatrix} \dot{c}_b \\ \dot{c}_a \end{pmatrix} = \hbar \omega (n + \frac{1}{2}) \begin{pmatrix} c_b \\ c_a \end{pmatrix} + \frac{1}{2} \hbar \begin{pmatrix} 2g\sqrt{n+1} & \\ & 2g\sqrt{n+1} \end{pmatrix} \begin{pmatrix} c_b \\ c_a \end{pmatrix}$$

for simplicity

$$i\hbar \dot{c}_b = \hbar g \sqrt{n+1} c_a$$

$$i\hbar \dot{c}_a = \hbar g \sqrt{n+1} c_b$$

$$\ddot{c}_b + g^2 (n+1) c_b = 0$$

$$c_b = \cos(g\sqrt{n+1} t)$$

$$c_a = -i \sin(g\sqrt{n+1} t)$$

$$P_b = \cos^2(g\sqrt{n+1} t)$$

$$P_a = \sin^2(g\sqrt{n+1} t)$$

$$|\psi(t)\rangle = \cos(g\sqrt{n+1} t) |b, n\rangle - i \sin(g\sqrt{n+1} t) |a, n+1\rangle$$

Quantum Rabi fluctuations

notice that even if $n=0$ (vacuum) there still will be floppings with frequency $2g$ (vacuum Rabi flopping)

So, a Fock state with a fixed number of photons $|n\rangle$ behaves very similarly to a classical Rabi flopping.

What about a coherent state?

Initially

$$|\Psi_{\text{atom}}\rangle_0 = c_a |a\rangle + c_b |b\rangle$$

$$|\Psi_{\text{light}}\rangle_0 = \sum_{n=0}^{\infty} c_n |n\rangle \quad c_n = e^{-|d|^2/2} \frac{d^n}{\sqrt{n!}} \quad \text{for a coherent state}$$

$$|\Psi(t=0)\rangle = |\Psi_{\text{atom}}\rangle_0 |\Psi_{\text{light}}\rangle_0$$

As we discussed before, light-atom interaction couples states $|a, n+1\rangle$ and $|b, n\rangle$ for all n present.

$$|\Psi(t)\rangle = \sum_{n=0}^{\infty} \left\{ [c_b c_n \cos(gt\sqrt{n+1}) - i c_a c_{n+1} \sin(gt\sqrt{n+1})] |b\rangle + [-i c_b c_{n-1} \sin(gt\sqrt{n}) + c_a c_n \cos(gt\sqrt{n})] |a\rangle \right\} |n\rangle$$

For $c_b = 1$ (we start with an atom in the excited state)

$$|\Psi(t)\rangle = \sum_{n=0}^{\infty} \left\{ c_n \cos(gt\sqrt{n+1}) |b\rangle - i c_{n-1} \sin(gt\sqrt{n}) |a\rangle \right\} |n\rangle$$

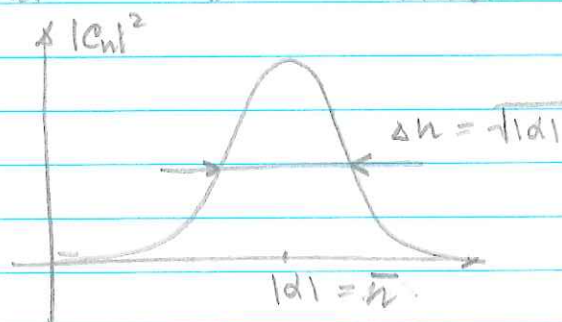
$$|\Psi_a(t)\rangle = \sum_{n=0}^{\infty} c_n \sin(gt\sqrt{n+1}) |n+1\rangle \quad \text{ground}$$

$$|\Psi_b(t)\rangle = \sum_{n=0}^{\infty} c_n \cos(gt\sqrt{n+1}) |n\rangle \quad \text{excited}$$

Average atomic inversion

$$\begin{aligned} \langle \Psi(t) | \hat{\sigma}_3 | \Psi(t) \rangle &= \langle \Psi(t) | |b\rangle \langle b| - |a\rangle \langle a| | \Psi(t) \rangle = \\ &= \langle \Psi_0 | \Psi_0 \rangle - \langle \Psi_a | \Psi_a \rangle = \sum_{n=0}^{\infty} |C_n|^2 \left[(\cos^2 gt \sqrt{n+1}) - \right. \\ &\quad \left. - \sin^2(gt \sqrt{n+1}) \right] = \sum_{n=0}^{\infty} |C_n|^2 \cos 2gt \sqrt{n+1} = \\ &= e^{-|d|^2} \sum_{n=0}^{\infty} \frac{|d|^{2n}}{n!} \cos(2gt \sqrt{n+1}) \end{aligned}$$

The output is a combination of many sine waves with somewhat different periods \rightarrow no clear Rabi fluctuations



Main contributing components lie between frequencies $2g\sqrt{n-\Delta n}$ and $2g\sqrt{n+\Delta n}$

Corresponding phase spread

$$\begin{aligned} 2gt_c (\sqrt{n+\Delta n} - \sqrt{n-\Delta n}) &\approx 2gt_c \sqrt{n} \left(\left(1 + \frac{\Delta n}{2n}\right) - \left(1 - \frac{\Delta n}{2n}\right) \right) \\ &\approx 2gt_c \frac{\Delta n}{\sqrt{n}} \approx 1 \quad \Rightarrow \quad gt_c \approx 1 \end{aligned}$$

$t_c \approx 1/g$ depends only on coupling strength

However we can also expect to see a revival of Rabi oscillations if

$$(g\sqrt{n+1} - g\sqrt{n}) t_R = 2\pi k \quad (k=0,1,2,\dots)$$

$$g\sqrt{n} \left(\frac{1}{2\sqrt{n}} - 1 \right) t_R = 2\pi \quad (k=1 \text{ for the first occurrence})$$

$$g \frac{1}{2\sqrt{n}} t_R = 2\pi$$

$$t_R = \frac{4\pi\sqrt{n}}{g}$$

The revival is never "complete" since the frequencies $\{g\sqrt{n}\}$ are not truly equidistant.

Why a coherent state is less "classical" than a number state?

Clear Rabi flopping require knowledge of precise intensity. Coherent state, as a minimum uncertainty state, has certain spread in its intensity distribution, that leads to the Rabi flopping diffusion.