

Two-level system evolution

1. Spin $\frac{1}{2}$ particle in $\vec{B} = \{0, 0, B_3\}$ magnetic field

$$\hat{H}|+z\rangle = \frac{1}{2}\hbar\omega_0 |+z\rangle \quad \hat{H}|-z\rangle = -\frac{1}{2}\hbar\omega_0 |-z\rangle$$

$$|+x\rangle(t) = \frac{1}{\sqrt{2}} (e^{-i\omega_0 t/2} |+z\rangle + e^{i\omega_0 t/2} |-z\rangle)$$

$$\langle S_z \rangle = 0 \quad \langle S_x \rangle(t) = \frac{\hbar}{2} \cos \omega_0 t \quad \langle S_y \rangle(t) = \frac{\hbar}{2} \sin \omega_0 t$$

2. Neutrino oscillations

three types of v:

three masses

$$\begin{matrix} v_e & v_u & v_t \\ m_1 & m_2 & m_3 \end{matrix}$$

do not
include
for
simplicity

Neutrino flavors don't match precisely
the masses!

mass \equiv energy eigenstates $|1,2\rangle$

$$E_{\text{kin}} = \sqrt{p^2 c^2 + m_{1,2}^2 c^4} \approx pc \left(1 + \frac{m_{1,2}^2 c^2}{2p^2}\right)$$

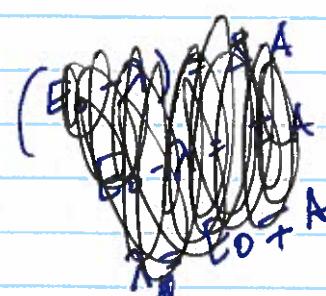
p-momentum, six-ed

$$\hat{H}|1,2\rangle = E_{1,2}|1,2\rangle \quad |1\rangle \rightarrow e^{-iE_1 t/\hbar} |1\rangle$$

$$|2\rangle \rightarrow e^{-iE_2 t/\hbar} |2\rangle$$

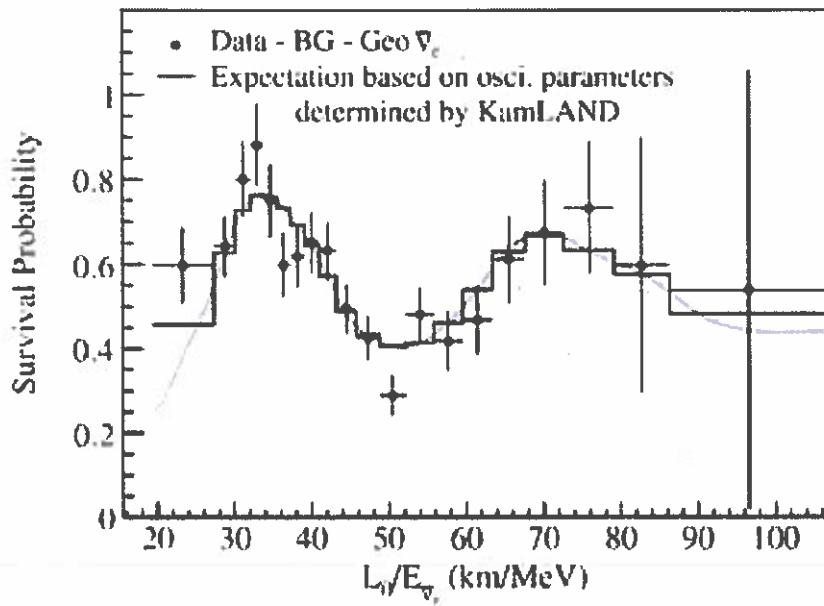
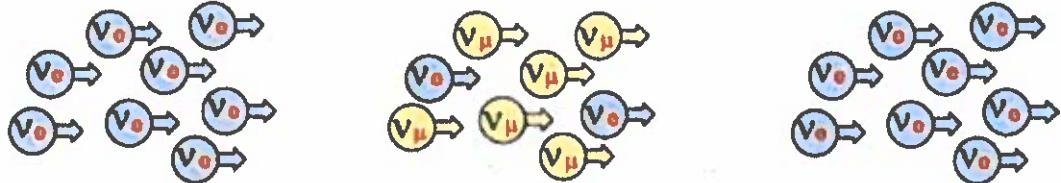
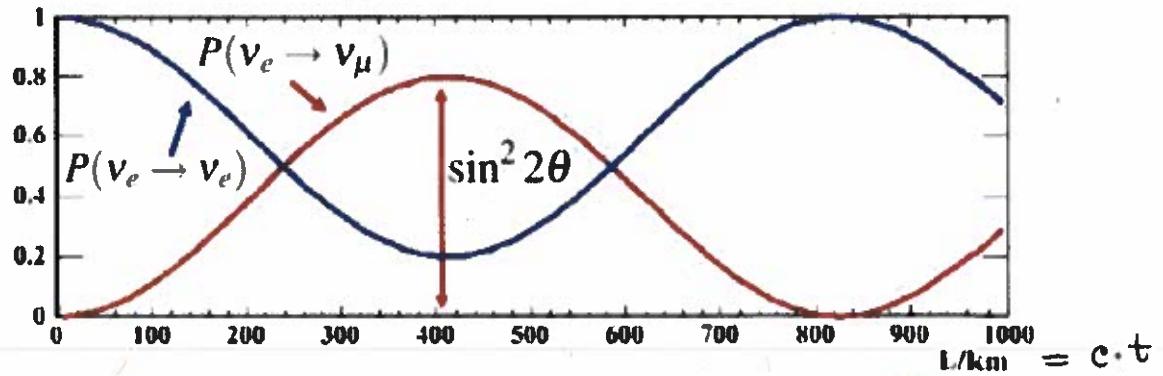
$$\begin{aligned} |\nu_e\rangle &= \cos\theta |1\rangle - \sin\theta |2\rangle \quad \Rightarrow \text{evolve in time!} \\ |\nu_\mu\rangle &= \sin\theta |1\rangle + \cos\theta |2\rangle \end{aligned}$$

$$P(\nu_e \rightarrow \nu_e) = 1 - 8 \sin^2 2\theta - 8 \sin^2 \left(\frac{\Delta m^2 c^4}{4 E h c} \right); \quad \Delta m^2 = m_1^2 - m_2^2$$



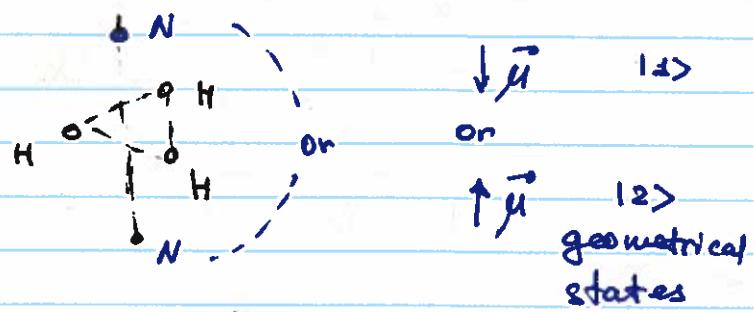
Neutrino oscillations

e.g. $\Delta m^2 = 0.003 \text{ eV}^2$, $\sin^2 2\theta = 0.8$, $E_\nu = 1 \text{ GeV}$



3. Ammonia molecules

Two orientations
 NH_3



Hamiltonian: $\hat{H} = \begin{pmatrix} E_0 & -A \\ -A & E_0 \end{pmatrix}$

$-A$ characterizes the energy required to "flip" the N -atom position

Eigenvalues of \hat{H} : $\det \begin{pmatrix} E_0 - \lambda & -A \\ -A & E_0 - \lambda \end{pmatrix} = 0$

$$(E_0 - \lambda)^2 - A^2 = 0 \quad \lambda_{\pm} = E_0 \pm A$$

Eigenstates $H | \pm \rangle = \lambda_{\pm} | \pm \rangle$

$$| \pm \rangle = \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \quad \lambda_+ = E_0 + A \quad \begin{pmatrix} -A & -A \\ -A & -A \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = 0$$

$$\lambda_- = E_0 - A \quad \begin{pmatrix} A & -A \\ -A & A \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = 0 \quad c_1 = -c_2$$

$$c_1 = c_2$$



$$E_+ = E_0 + A$$

$$| + \rangle = \frac{1}{\sqrt{2}} (| 12 \rangle + | 21 \rangle)$$

Symmetric superposition

$$E_- = E_0 - A$$

$$| - \rangle = \frac{1}{\sqrt{2}} (| 12 \rangle - | 21 \rangle)$$

anti-symmetric superposition

$$| 12 \rangle = \frac{1}{\sqrt{2}} | + \rangle + \frac{1}{\sqrt{2}} | - \rangle \rightarrow \frac{1}{\sqrt{2}} e^{-i(E_0+A)t/\hbar} | + \rangle + \frac{1}{\sqrt{2}} e^{-i(E_0-A)t/\hbar} | - \rangle$$

$$\begin{aligned}
 |\pm\rangle(t) &= \frac{1}{\sqrt{2}} e^{-iE_0 t/\hbar} \left[e^{-iAt/\hbar} |+\rangle + e^{iAt/\hbar} |-\rangle \right] = \\
 &= \frac{1}{\sqrt{2}} e^{-iE_0 t/\hbar} \left[\left(\cos \frac{At}{\hbar} + i \sin \frac{At}{\hbar} \right) |+\rangle + \left(\cos \frac{At}{\hbar} + i \sin \frac{At}{\hbar} \right) |-\rangle \right] \\
 &= e^{-iE_0 t/\hbar} \left[\cos \frac{At}{\hbar} |\pm\rangle + \frac{i}{\hbar} \sin \frac{At}{\hbar} |\mp\rangle \right]
 \end{aligned}$$

If we placed a NH_3 molecule in one position, it is going to flip back and forth with frequency $\nu_{\text{NH}_3} = 2A/\hbar = 24 \text{ GHz}$

The basis for NH_3 maser operation

In reality, we can much more efficiently flip the system b/w two states using resonant electric or magnetic field oscillating on the frequency of this transition

Back to spins: strong constant B_0 along z and rf field $B_1 \cos \omega t$ along x

The off-diagonal term in the hamiltonian is proportional to B_1

$$\hat{H} = \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 \cos \omega t \\ \omega_1 \cos \omega t & -\omega_0 \end{pmatrix}$$

induced transitions,
the highest p rate
if $\omega = \omega_0$

$|+\rangle$ $E_+ = \hbar \omega_0 / 2$



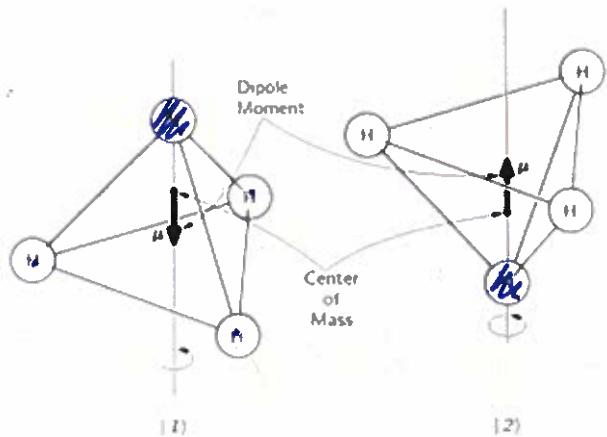
$|-\rangle$ $E_- = -\hbar \omega_0 / 2$



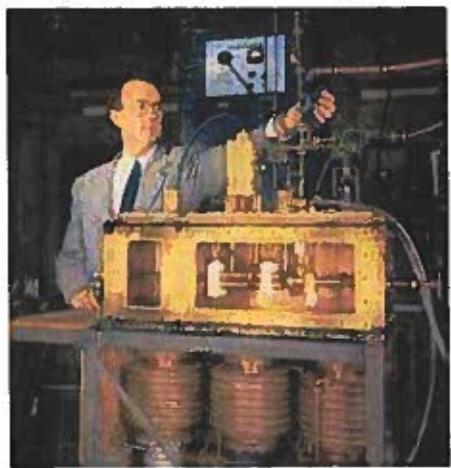
$$P_{-z}(t) = \cos^2 \frac{\omega_1 t}{2}$$

$$P_{+z}(t) = \sin^2 \frac{\omega_1 t}{2}$$

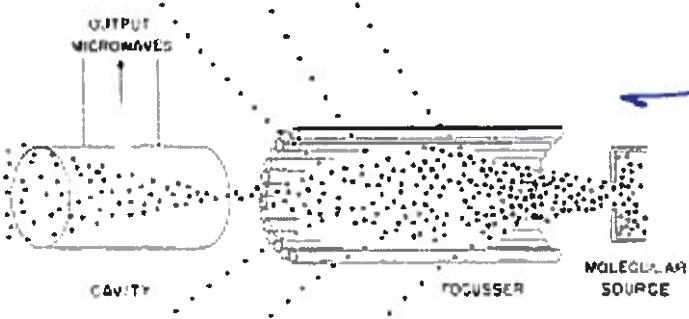
NH_3



Two equally possible
orientations



NH_3 maser



Microwave
Amplification via
Stimulated
Emission of
Radiation