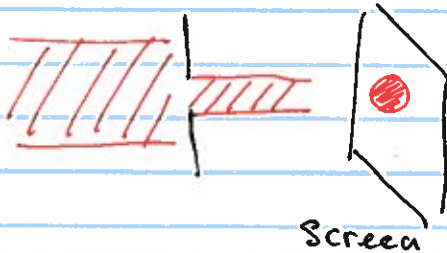


Diffraction

Diffraction is also an interference phenomena. We usually reserve this term to describe what happens to light when it passes through apertures or edges.

Naive expectation
(or light as particles)



is a spot matching
a size of the opening

Huygens principle
Very small opening



acts as a point
source

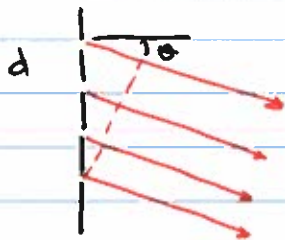


How to calculate the diffraction image?

Method 1: imaging the opening as
a combination of many slits

Method 2: exact calculations (with
complex waves)

Reminder: Interference pattern from N slits at distance d from each other



Each consecutive slit introduces delay of $d \sin \theta$ or phase delay $2\pi d \sin \theta / \lambda$

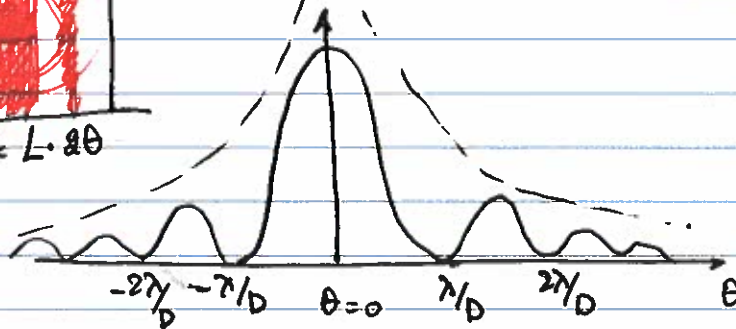
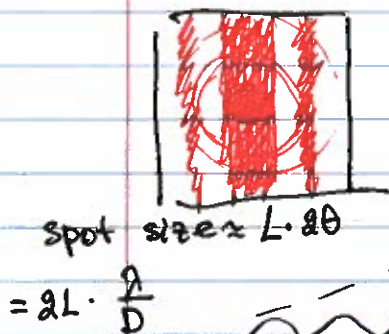
$$I_{N \text{ slits}} = I_2 \left(\frac{\sin \left(\frac{\pi N d \sin \theta}{\lambda} \right)}{\sin \left(\frac{\pi d \sin \theta}{\lambda} \right)} \right)^2 \quad I_2 - \text{intensity from one slit}$$

Method 1: ~~define~~ an opening of size $D = dN$
 $\Rightarrow d \rightarrow 0$ and $N \rightarrow \infty$, $I_2 = I_0 / N^2$

$$\sin \frac{\pi N d \sin \theta}{\lambda} \Rightarrow \sin \frac{\pi D \sin \theta}{\lambda} \quad \uparrow \text{total intensity}$$

$$I_{\text{opening}} = \frac{I_0}{N^2} \left[\frac{\sin \left(\frac{\pi D \sin \theta}{\lambda} \right)}{\frac{\pi d \sin \theta}{\lambda}} \right]^2$$

$$= I_0 \left[\frac{\sin \frac{\pi D \sin \theta}{\lambda}}{\frac{\pi D \sin \theta}{\lambda}} \right]$$



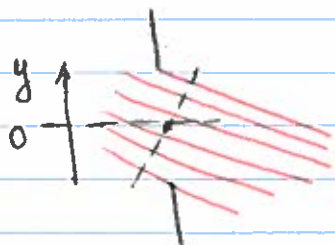
First diffraction minimum

$$\sin \frac{\pi D \sin \theta_{\min}}{\lambda} = 0$$

$$\frac{D \sin \theta_{\min}}{\lambda} = 1$$

$$\theta_{\min} \approx \sin \theta_{\min} = \frac{\lambda}{D}$$

for a 1D opening



For any height above the center of the slit y , the retardation (compare to the center) is $y \cdot \sin \theta$
 $-a/2 < y < a$ - slit. $k = \frac{2\pi}{\lambda}$

The electric field in each point along the wavefront $dE(y) = E_0 \frac{dy}{a} e^{ik \cdot y \cdot \sin \theta}$

Total electric field

$$E_{\text{tot}} = \int_{-a/2}^{a/2} dE_y = \frac{E_0}{a} \int_{-a/2}^{a/2} dy e^{ik \sin \theta \cdot y} dy =$$

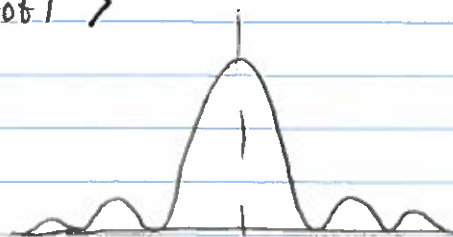
$$= \frac{E_0}{a} \frac{1}{ik \sin \theta} \left[e^{ik \sin \theta a/2} - e^{-ik \sin \theta a/2} \right] =$$

$$= \frac{E_0}{a} \frac{2}{k} \frac{\sin[k \sin \theta \cdot a/2]}{\sin \theta} = E_0 \frac{\sin \left[\frac{ka}{2} \sin \theta \right]}{\left[\frac{ka}{2} \sin \theta \right]}$$

$\left\{ \begin{aligned} \int e^{+dx} dx &= \\ &= e^{+dx} / +d \end{aligned} \right.$

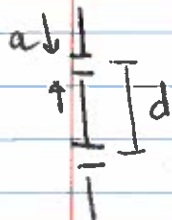
Power (or intensity) $P \propto \langle |E_{\text{tot}}|^2 \rangle$

$$P = P_0 \frac{\sin^2 \left[\frac{\pi a}{\lambda} \sin \theta \right]}{\left(\frac{\pi a}{\lambda} \sin \theta \right)^2}$$



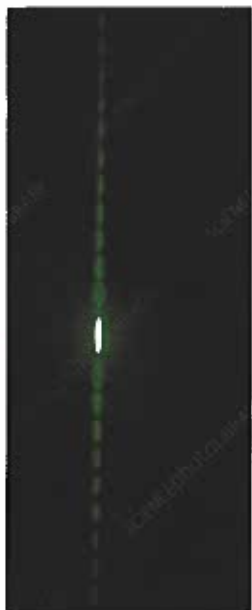
Two slits next to each other

$$P = P_0 \underbrace{\frac{\sin^2 \left[\frac{\pi a}{\lambda} \sin \theta \right]}{\left(\frac{\pi a}{\lambda} \sin \theta \right)^2}}_{\text{single-slit diffraction}} \cdot \underbrace{\cos^2 \left(\frac{\pi d \sin \theta}{\lambda} \right)}_{\text{two-slit interference}}$$

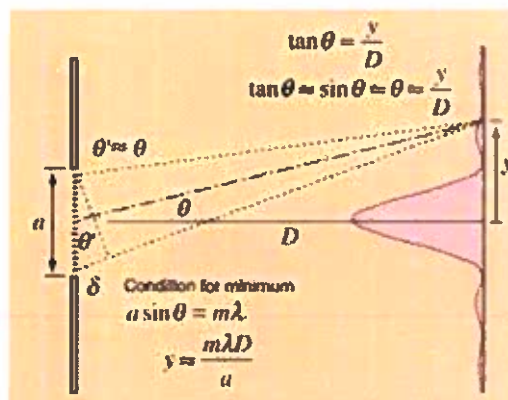


single-slit diffraction

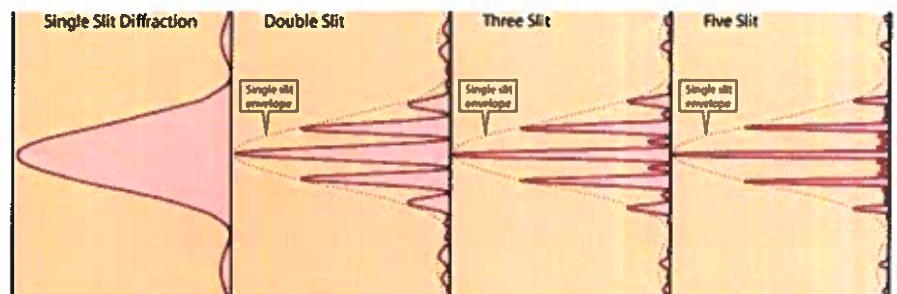
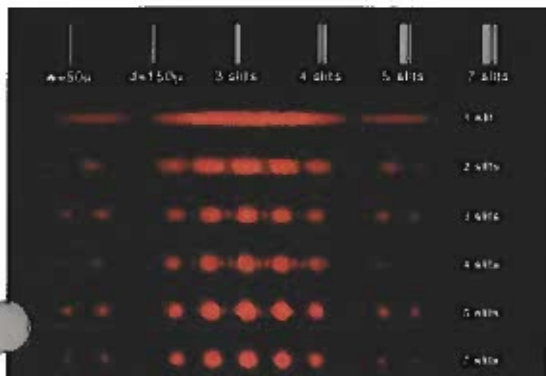
two-slit interference



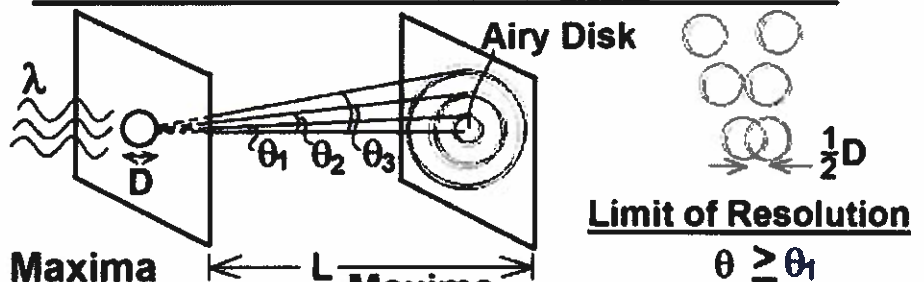
FQAAAAAdAAAAABAE



<https://www.google.com/url?sa=i&url=https%3A%2F%2Fsciencedemonstrations.fas.harvard.edu%2Fpresentations%2Fpoissons-spot&psig=AOvVaw3uY6WVfQUEBNp6IfC0AaMd&ust=1714228354551000&source=images&cd=vfe&opi=89978449&ved=0CBIQjRxqFwoTCPDjmbeM4IU>

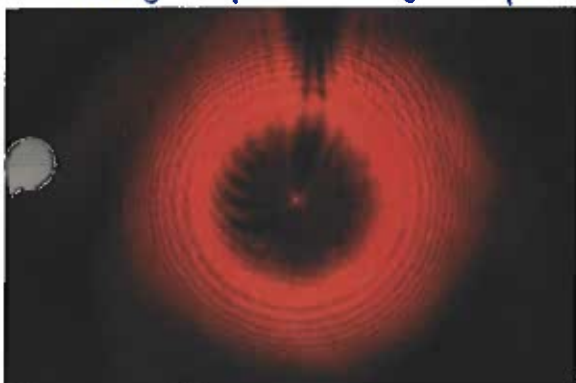


Diffraction Pattern of a Circular Aperture: Intro

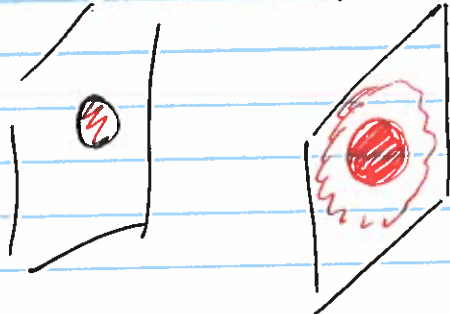


$$\begin{aligned}
 \sin \theta_1 &= 1.22 \frac{\lambda}{D} & \sin \theta_{\text{Max1}} &= 1.63 \frac{\lambda}{D} \\
 \sin \theta_2 &= 2.23 \frac{\lambda}{D} & \sin \theta_{\text{Max2}} &= 2.68 \frac{\lambda}{D} \\
 \sin \theta_3 &= 3.24 \frac{\lambda}{D} & \sin \theta_{\text{Max3}} &= 3.70 \frac{\lambda}{D}
 \end{aligned}$$

Arago spot - bright spot in the center of a circular target "shadow"



2D circular aperture



$$I(\theta) = I_0 \left(\frac{J_1\left(\frac{\pi D \sin \theta}{\lambda}\right)}{\left(\frac{\pi D \sin \theta}{\lambda}\right)} \right)^2$$

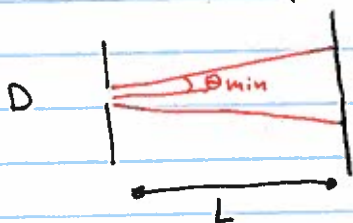
$J_1(x)$ - Bessel function of the first kind

$J_1(x) = 0$ for $x = 3.83$

$$\frac{\pi D \sin \theta_{\min}}{\lambda} = 3.83$$

$$\theta_{\min} \approx 1.22 \frac{\lambda}{D}$$

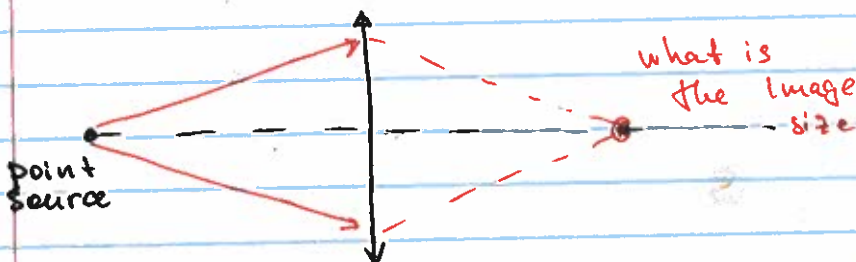
Optical resolution - how well one can resolve optical features



beam spot size $\approx L \theta_{\min}$

The smaller D is, the larger is $\theta_{\min} \rightarrow$ larger spot on the screen.

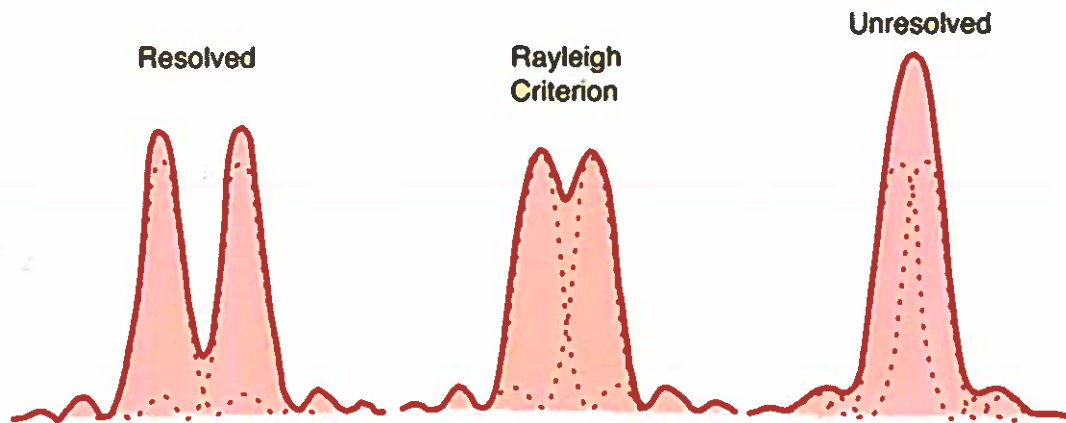
Same equation will be valid for a lens



We cannot recreate the image to be a point source exactly, since we cannot collect all the light emitted by the source!

$$\text{min spot} \propto f \cdot \frac{\lambda}{D}$$

larger lens $\rightarrow$ better sharpness



$$\sin \theta_R = \frac{\lambda}{d}$$

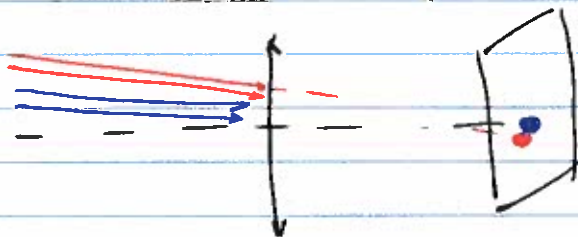
Single slit

$$\sin \theta_R = 1.22 \frac{\lambda}{d}$$

Circular aperture

Why the telescopes have to be large?

To resolve better!



Rayleigh criteria:

Two objects
can be resolved
if their angular
separation is
larger than $\theta_{\min}$

How well our eyes can do?

Assuming a ~~pupil~~ pupil size $D \approx 5\text{mm}$
and average visible light $\lambda \approx 500\text{nm} =$
 $= 5 \cdot 10^{-7}\text{m}$

$$\theta_{\min} = 1.22 \frac{\lambda}{D} = 1.22 \cdot \frac{5 \cdot 10^{-7}\text{m}}{5 \cdot 10^{-3}\text{m}} \approx 1.22 \cdot 10^{-4} \text{ rad}$$

~~angular minutes~~

$$= 180^\circ/\text{rad} \cdot 60 \frac{\text{min}}{\text{deg}} \cdot 1.22 \cdot 10^{-4} \text{ rad} \approx 1' \quad (\text{angular minute})$$

How far you can be to resolve $\Delta x = 1\text{-mm}$ size feature

$$L_{\min} = \frac{\Delta x}{\theta_{\min}} = \frac{10^{-3}}{1.22 \cdot 10^{-4}} \approx 8\text{m}$$

(if they are bright enough!)

How to increase the resolution?

- Use shorter wavelength
moving from CD (808nm) to DVD (405nm)
allowed to halve the size of features
and quadruple the information capacity
- X-ray diffraction (XRD) - $\lambda_{\text{x-ray}} = 0.03 - 3\text{nm}$
- Modern chip technology uses 10nm wavelength
(extreme UV light)

In 2030 they plan to reach a few
angstrom (10^{-10}m)

- To resolve even smaller features → use heavier particles

Anything can display wave properties

$$\lambda = \frac{h}{p}$$

De-Broglie wavelength

$$h = 10^{-34} \text{ J.s}$$

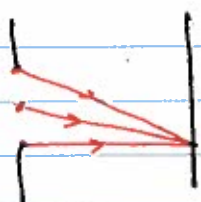
Planck's constant

p - particle momentum

Electron diffraction is a common tool for material science

Neutrons are also used for material science to probe deeper into material (since they don't have electric charge)

- Near-field imaging - collecting and analyzing light at distances close to the feature size



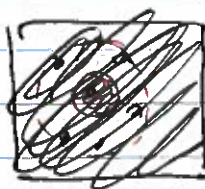
Near-field images contain much more information and can provide much better resolution → but hard to collect

(Check out Dr. Muntaz Qazilbash research)

- Nonlinear sub-diffraction imaging



normal
laser
beam



structured



only
unsaturated
source
is
visible

Large, diffraction-limited laser spot excites many objects, making it impossible to resolve the individual contributions

"blinding" light
STED imaging (Nobel prize 2014)
Vortex (donut) beam "blinds" all objects except for one sitting in the center, where its intensity is zero